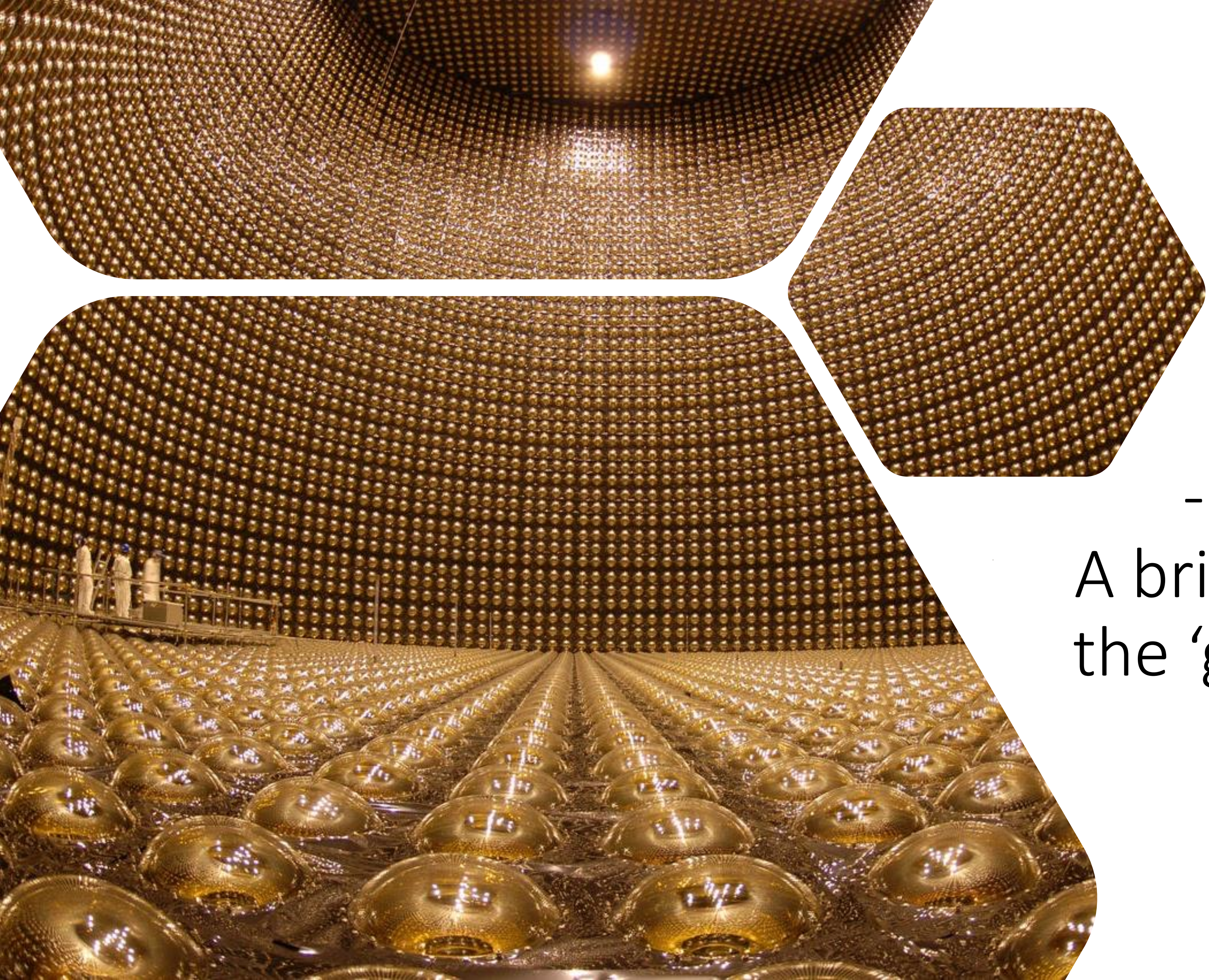
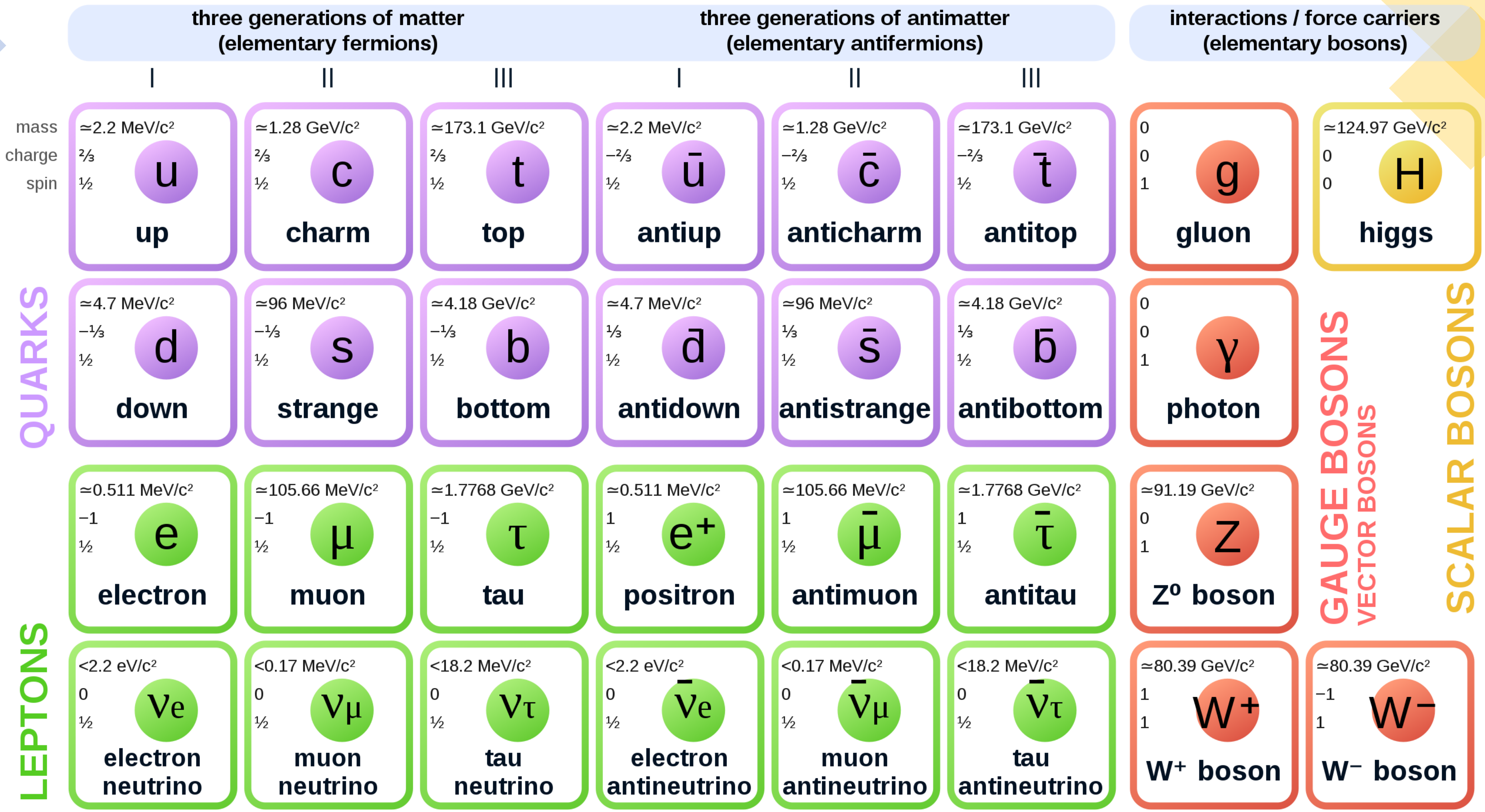


-Neutrinos-
A brief overview of the ‘ghost particle’.

By Jaiden Parlone

Standard Model of Elementary Particles



	three generations of matter (elementary fermions)			three generations of antimatter (elementary antifermions)			interactions / force carriers (elementary bosons)	
	I	II	III	I	II	III		
mass charge spin	≃2.2 MeV/c^2 ⅔ ½ u up	≃1.28 GeV/c^2 ⅔ ½ c charm	≃173.1 GeV/c^2 ⅔ ½ t top	≃2.2 MeV/c^2 −⅔ ½ $\bar{u}$ antiup	≃1.28 GeV/c^2 −⅔ ½ $\bar{c}$ anticharm	≃173.1 GeV/c^2 −⅔ ½ $\bar{t}$ antitop	0 0 1 g gluon	≃124.97 GeV/c^2 0 0 H higgs
QUARKS	≃4.7 MeV/c^2 −⅓ ½ d down	≃96 MeV/c^2 −⅓ ½ s strange	≃4.18 GeV/c^2 −⅓ ½ b bottom	≃4.7 MeV/c^2 ⅓ ½ $\bar{d}$ antidown	≃96 MeV/c^2 ⅓ ½ $\bar{s}$ antistrange	≃4.18 GeV/c^2 ⅓ ½ $\bar{b}$ antibottom	0 0 1 γ photon	GAUGE BOSONS VECTOR BOSONS SCALAR BOSONS
	≃0.511 MeV/c^2 −1 ½ e electron	≃105.66 MeV/c^2 −1 ½ μ muon	≃1.7768 GeV/c^2 −1 ½ τ tau	≃0.511 MeV/c^2 1 ½ e^+ positron	≃105.66 MeV/c^2 1 ½ $\bar{\mu}$ antimuon	≃1.7768 GeV/c^2 1 ½ $\bar{\tau}$ antitau	≃91.19 GeV/c^2 0 1 Z Z^0 boson	
LEPTONS	<2.2 eV/c^2 0 ½ ν_e electron neutrino	<0.17 MeV/c^2 0 ½ ν_μ muon neutrino	<18.2 MeV/c^2 0 ½ ν_τ tau neutrino	<2.2 eV/c^2 0 ½ $\bar{\nu}_e$ electron antineutrino	<0.17 MeV/c^2 0 ½ $\bar{\nu}_\mu$ muon antineutrino	<18.2 MeV/c^2 0 ½ $\bar{\nu}_\tau$ tau antineutrino	≃80.39 GeV/c^2 1 1 W^+ W^+ boson	≃80.39 GeV/c^2 −1 1 W^- W^- boson

1 proton : 1 electron : 1 neutron : 1 billion neutrinos

Abundance 2nd only to photons!

Bananas emit around 1,000,000 neutrinos per day, mostly from Potassium-40 beta decay.

You emit around 300,000,000 per day as well.

Both are these are nothing compared to the sun’s output, which leads to around 100,000,000,000,000 (100 trillion) passing through your body every second.

However…

There is only about a 1 in 4 chance of a neutrino interacting within your body within your lifetime.

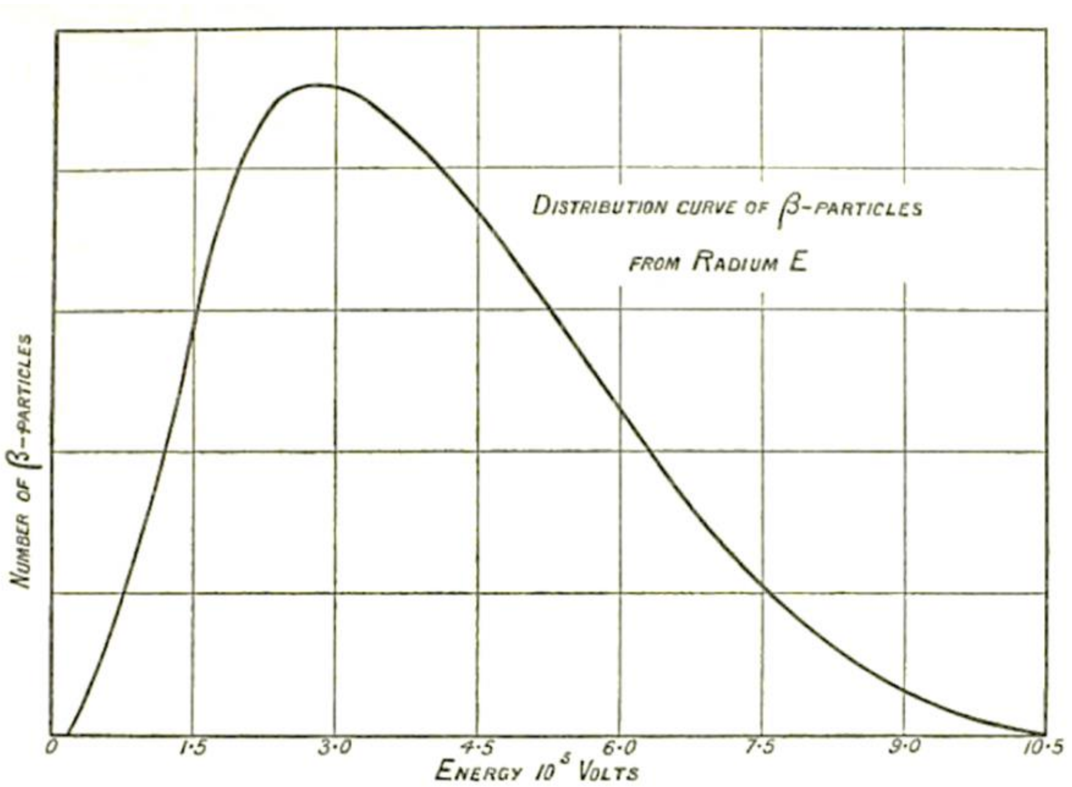
And if you were holding a banana, it would take about 2 billion years before a neutrino from it interacted with you.

In the early 20th Century, beta decay was thought to be a two body process, described by:

$$n \rightarrow p + e^{-}$$

In 1914 James Chadwick discovered that the energy spectrum of electrons emitted during beta decay of radium was continuous as opposed to discrete, apparently in violation of conservation of energy.

This type of decay would have also violated later laws of angular momentum and lepton number conservation.



Note that this diagram is actually from a confirmation paper from 1927; Doi: 10.1098/rspa.1927.0168.

In 1930 Pauli first postulated a neutral, spin-$\frac{1}{2}$ particle with negligible mass that was also produced in this process. This he named ‘neutron’ but it was later renamed to ‘neutrino’ (the Italian equivalent of "little neutral one") after Chadwick’s discovery of the atomic neutron. This particle would later be fully realised as the electron antineutrino.

$$n \rightarrow p + e^- + \bar{\nu}_e$$

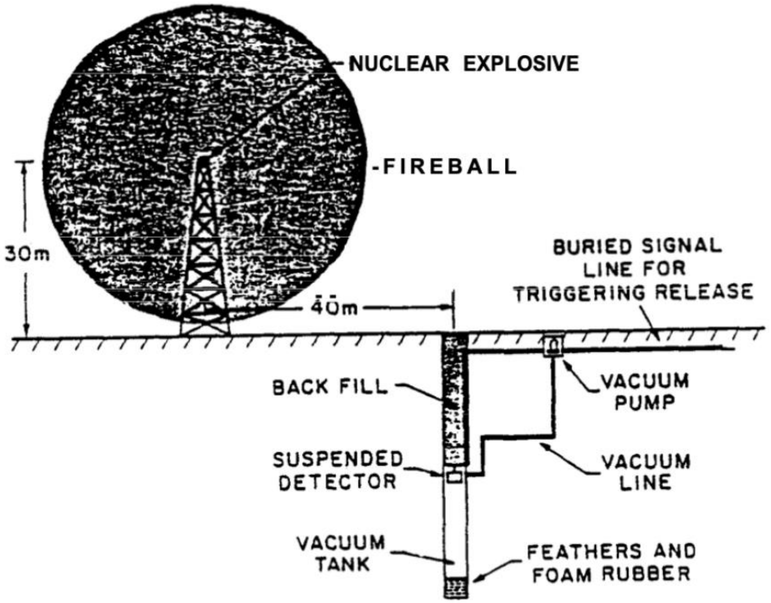
However, due to this particle’s properties Pauli famously stated:

In 1956 the team working on the Cowan-Reines neutrino experiment published results utilised antineutrinos created in a nuclear ~~blast~~ reactor to induce a process known as ‘inverse beta decay’:

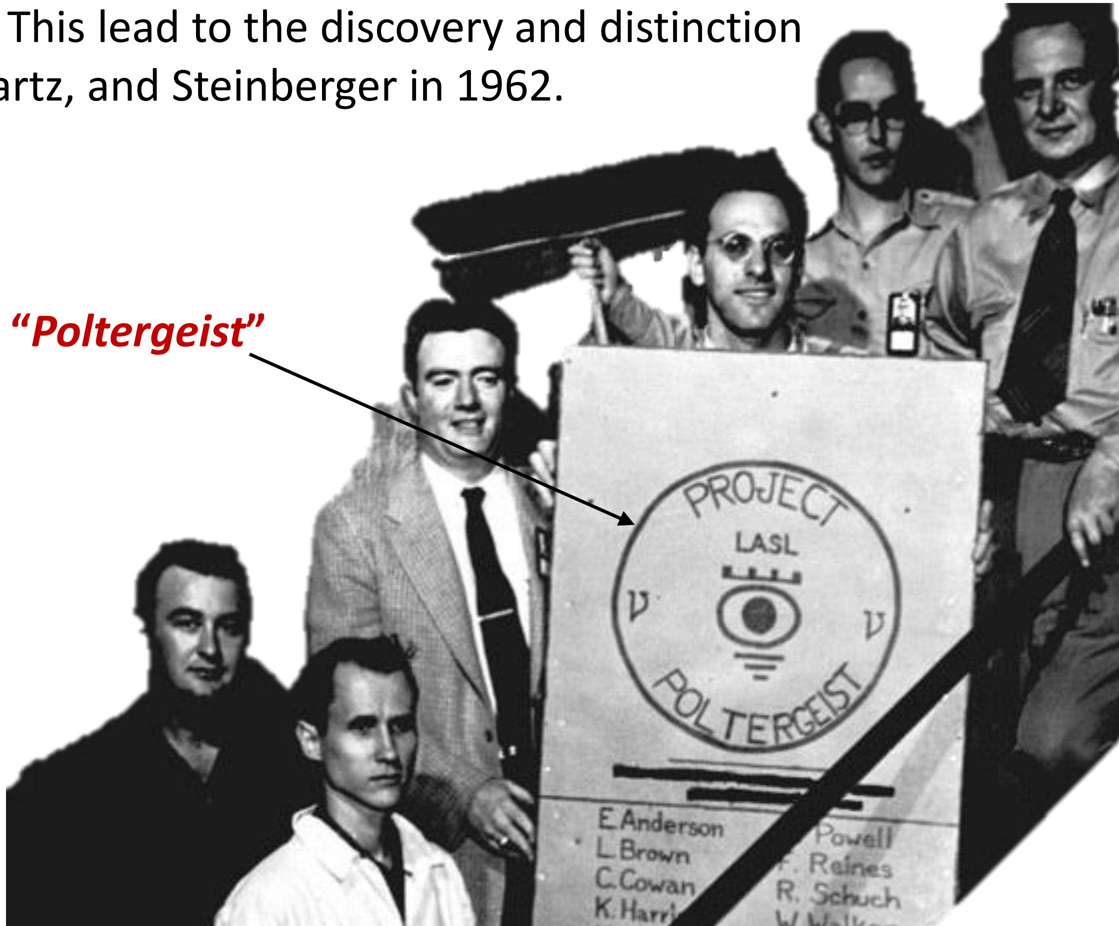
$$\bar{\nu}_e + p \rightarrow n + e^-$$

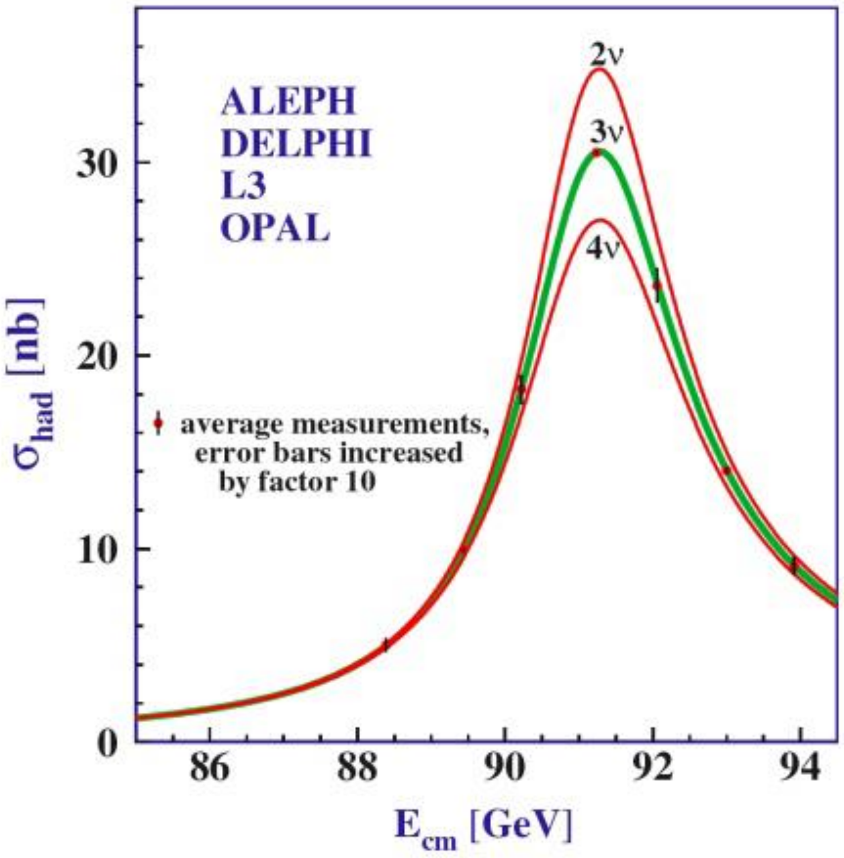
This effort was awarded a Nobel Prize in 1995, almost 45 years later.

However, open questions still remained. These neutrinos had only been observed to produce electrons in interactions, whereas there was now another known lepton, the muon. This lead to the discovery and distinction of the muon neutrino, ν_μ, by Lederman, Schwartz, and Steinberger in 1962.



Project “*Poltergeist*”





Studies of Z^0 boson decay allows for the number of light active neutrinos to be determined. This is done through the decay mode;

$$Z^0 \rightarrow \nu_l + \bar{\nu}_l$$

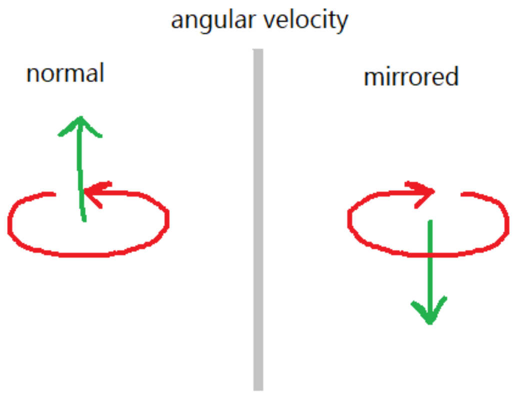
Where each neutrino flavour contributes to the decay width.

From this, the results from LEP determined to the number of light, active neutrino species to be: $N_\nu = 2.984 \pm 0.008$

This was achieved before the DONUT collaboration discovered the tau neutrino, ν_τ, in July 2000.

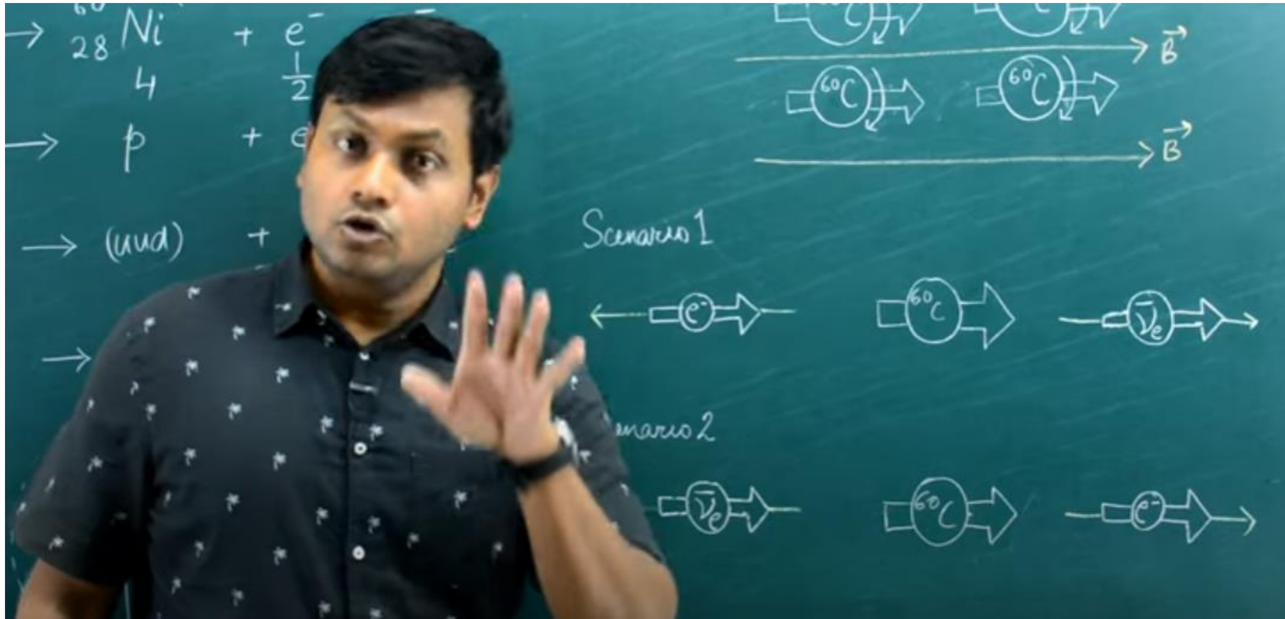
Time to get interactive

- Curl your hands!
- The direction of axial vector/ pseudo-vector follows a curling right hand.
  - The curl is the rotation (or similar in other examples), and in classical examples is formed of the combination (cross product) of two perpendicular polar (regular) vectors.
  - The thumb is the direction of the vector.
  - Where a regular vector would not be mirrored, pseudovectors are.

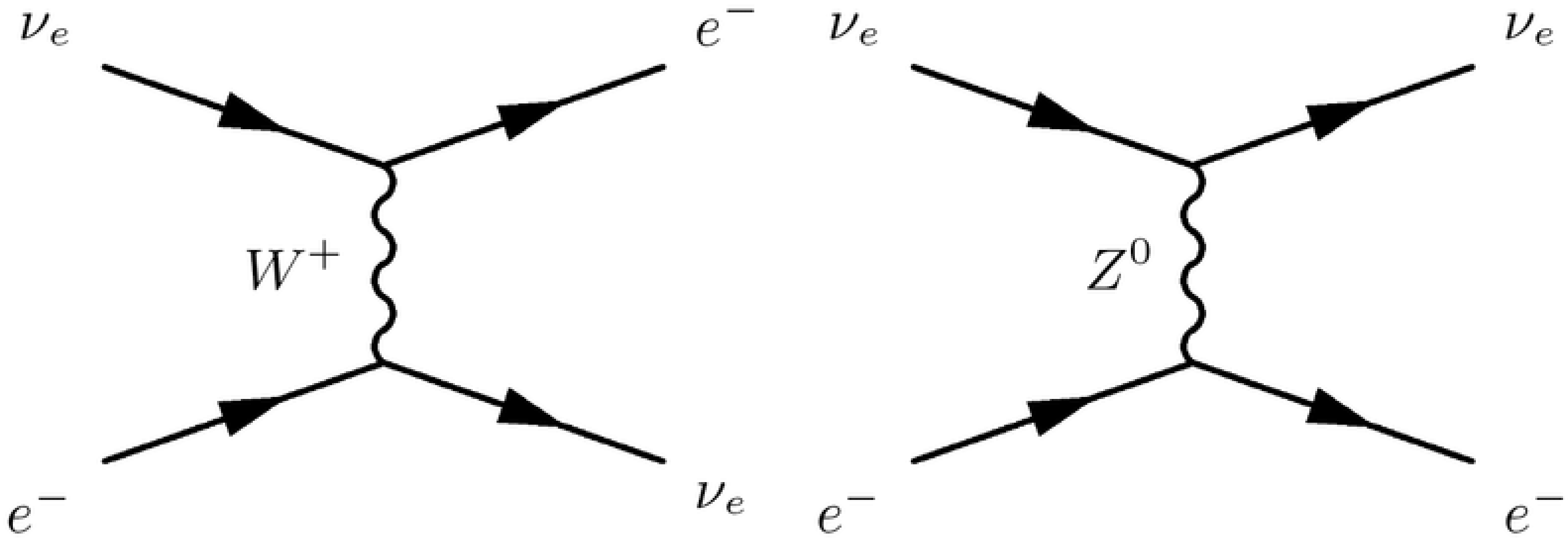


- We can also use this curling hand trick to visualise the ‘handedness’ of a particle, aligning or unaligning momentum with spin.

https://www.youtube.com/watch?v=v44jEXN4sSY



Because they have no colour charge, or electric charge, neutrinos interact only through the Weak Force.

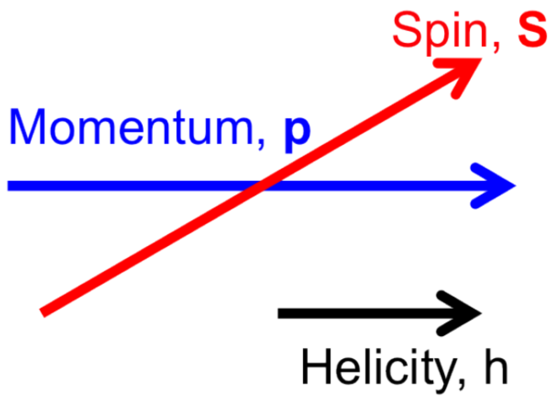


$W^\pm$mediated interactions are the only known flavour changing interaction, and are referred to as the 'Charged Current' in many sources.

The fact that the $W^\pm$ and Z^0 are massive are a result of symmetry breaking in the Higgs Field, the derivation of which tends to be covered in the 3rd or 4th year 'Gauge Theories' course.

Whilst this next bit is contained within the electroweak theory, it's good to describe its origins:

Helicity is defined as the projection of the spin on to the direction of the particle’s momentum.



$$h = \frac{\sigma \cdot \mathbf{p}}{|p|}$$

Right helicity = spin aligned with direction of momentum. $\vec{\sigma} \uparrow\uparrow \vec{p}$

Left helicity = spin anti-aligned with direction of momentum. $\vec{\sigma} \uparrow\downarrow \vec{p}$

Chirality is a bit more abstract, and is best thought about in terms of transformations (i.e. a reflection of a chiral object is not the same), but in the relativistic limit, helicity and chirality become equivalent.

When a parity operation (flip coordinates) is performed on helicity:

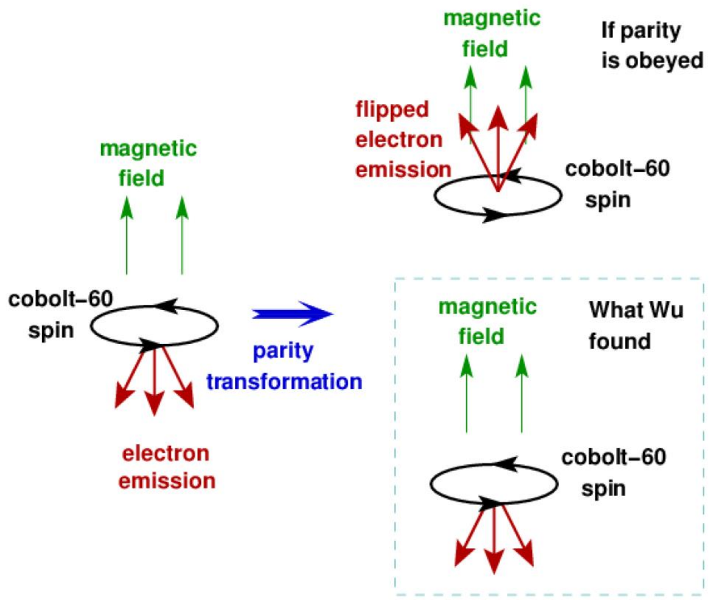
Momentum: Polar vector, changes sign: **p** → -**p**
Spin: Axial vector, $\sigma \rightarrow \sigma$

Therefore helicity should always change under parity transformation… right?

C.S. Wu checked this was correct for Weak interactions and in 1957 found that it simply wasn’t. Parity was violated, apparently maximally. This understandably traumatised many physicists.

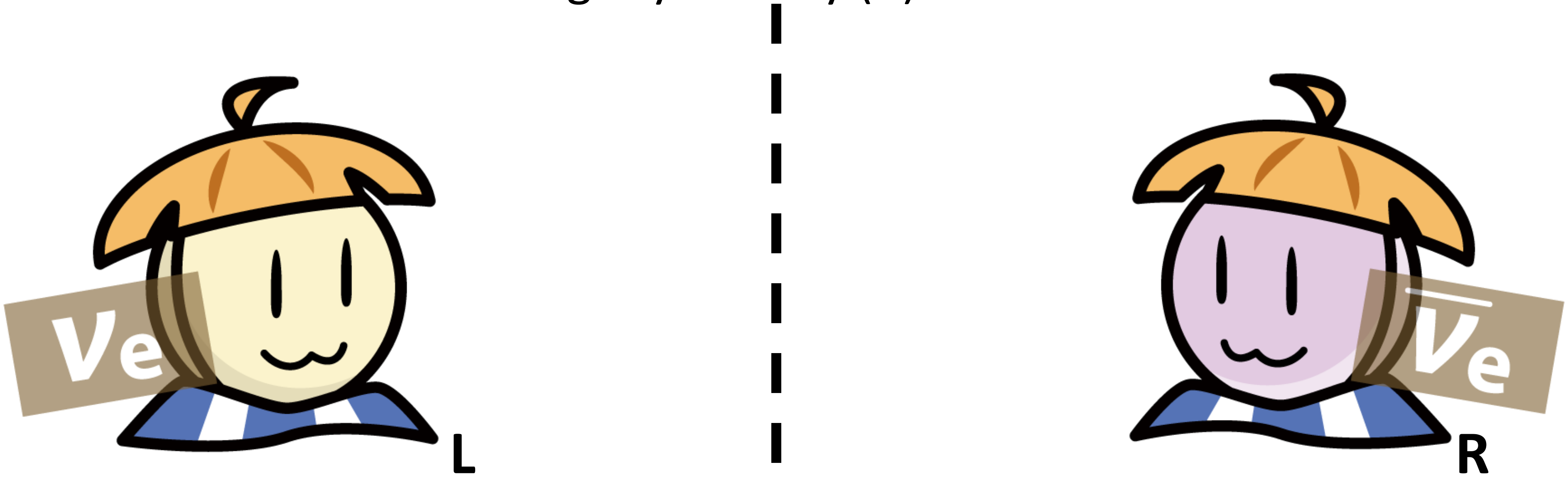
$$^{60}_{27}\mathrm{Co} \rightarrow {}^{60}_{28}\mathrm{Ni}^* + e^- + \bar{\nu}_e$$
$$^{60}_{28}\mathrm{Ni}^* \rightarrow {}^{60}_{28}\mathrm{Ni} + 2\gamma$$

- Apply magnetic field
- Cool apparatus to 0.01 K
- High proportion of ^{60}Co spins are aligned with the magnetic field
- Examine distribution of emitted electrons and photons.



Lee and Yangs’ contributions to the theory of Parity violation were rewarded with a Nobel prize in that same year. Though Wu’s role was mentioned, she unfortunately wasn’t officially honoured until 1978 with the Wolf Prize.

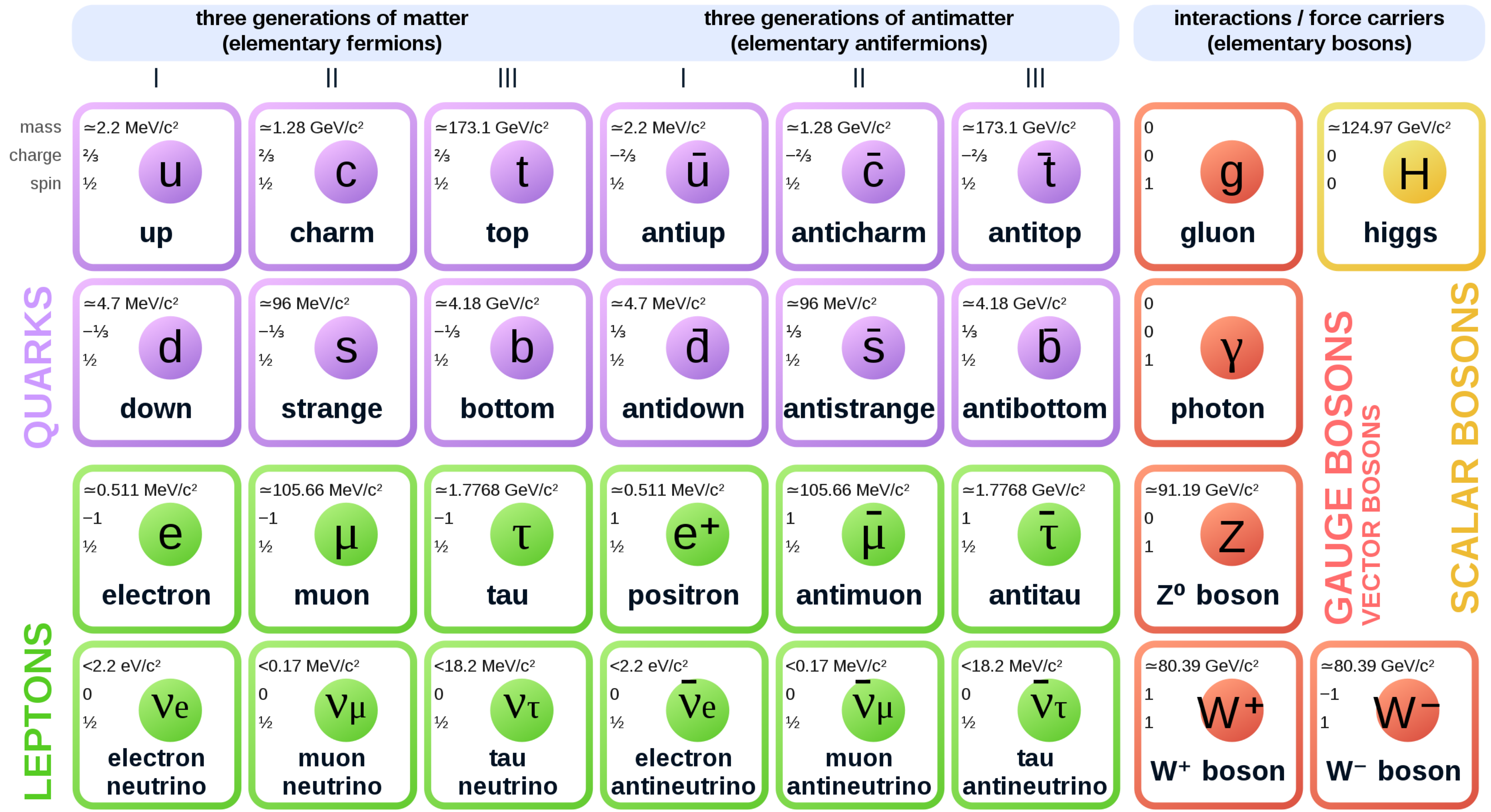
The charged weak bosons only couple to left handed particles, and right handed antiparticles. As neutrinos can only interact weakly, only left-handed neutrinos and right-handed antineutrinos exist, but their opposite chiral partners do not (in the standard model). This also means that Charge symmetry (C) is broken.



Charge-Parity Symmetry

A note is that charge conjugation reverses all internal quantum numbers as well as reversing electric charge.

Standard Model of Elementary Particles



	three generations of matter (elementary fermions)			three generations of antimatter (elementary antifermions)			interactions / force carriers (elementary bosons)	
	I	II	III	I	II	III		
mass charge spin	≃2.2 MeV/c^2 ⅔ ½ u **up**	≃1.28 GeV/c^2 ⅔ ½ c **charm**	≃173.1 GeV/c^2 ⅔ ½ t **top**	≃2.2 MeV/c^2 −⅔ ½ ū **antiup**	≃1.28 GeV/c^2 −⅔ ½ c̄ **anticharm**	≃173.1 GeV/c^2 −⅔ ½ t̄ **antitop**	0 0 1 g **gluon**	≃124.97 GeV/c^2 0 0 H **higgs**
QUARKS	≃4.7 MeV/c^2 −⅓ ½ d **down**	≃96 MeV/c^2 −⅓ ½ s **strange**	≃4.18 GeV/c^2 −⅓ ½ b **bottom**	≃4.7 MeV/c^2 ⅓ ½ d̄ **antidown**	≃96 MeV/c^2 ⅓ ½ s̄ **antistrange**	≃4.18 GeV/c^2 ⅓ ½ b̄ **antibottom**	0 0 1 γ **photon**	GAUGE BOSONS VECTOR BOSONS SCALAR BOSONS
	≃0.511 MeV/c^2 −1 ½ e **electron**	≃105.66 MeV/c^2 −1 ½ μ **muon**	≃1.7768 GeV/c^2 −1 ½ τ **tau**	≃0.511 MeV/c^2 1 ½ e^+ **positron**	≃105.66 MeV/c^2 1 ½ $\bar{\mu}$ **antimuon**	≃1.7768 GeV/c^2 1 ½ $\bar{\tau}$ **antitau**	≃91.19 GeV/c^2 0 1 Z **Z^0 boson**	
LEPTONS	<2.2 eV/c^2 0 ½ ν_e **electron neutrino**	<0.17 MeV/c^2 0 ½ ν_μ **muon neutrino**	<18.2 MeV/c^2 0 ½ ν_τ **tau neutrino**	<2.2 eV/c^2 0 ½ $\bar{\nu}_e$ **electron antineutrino**	<0.17 MeV/c^2 0 ½ $\bar{\nu}_\mu$ **muon antineutrino**	<18.2 MeV/c^2 0 ½ $\bar{\nu}_\tau$ **tau antineutrino**	≃80.39 GeV/c^2 1 1 W^+ **W^+ boson**	≃80.39 GeV/c^2 −1 1 W^- **W^- boson**

With this revelation of broken symmetries, the unified electroweak theory is fully within the standard model. So all is well! Neutrinos are massless, neutral, spin-$\frac{1}{2}$, chiral elementary fermions. Unless….

Solar neutrino problem

In the late 1960s Raymond Davis, a chemist by trade, knew that electron neutrinos interacting with chlorine would produce a radioactive isotope of argon:

$$\nu_e + Cl^{37} \rightarrow Ar^{37} + e^-$$

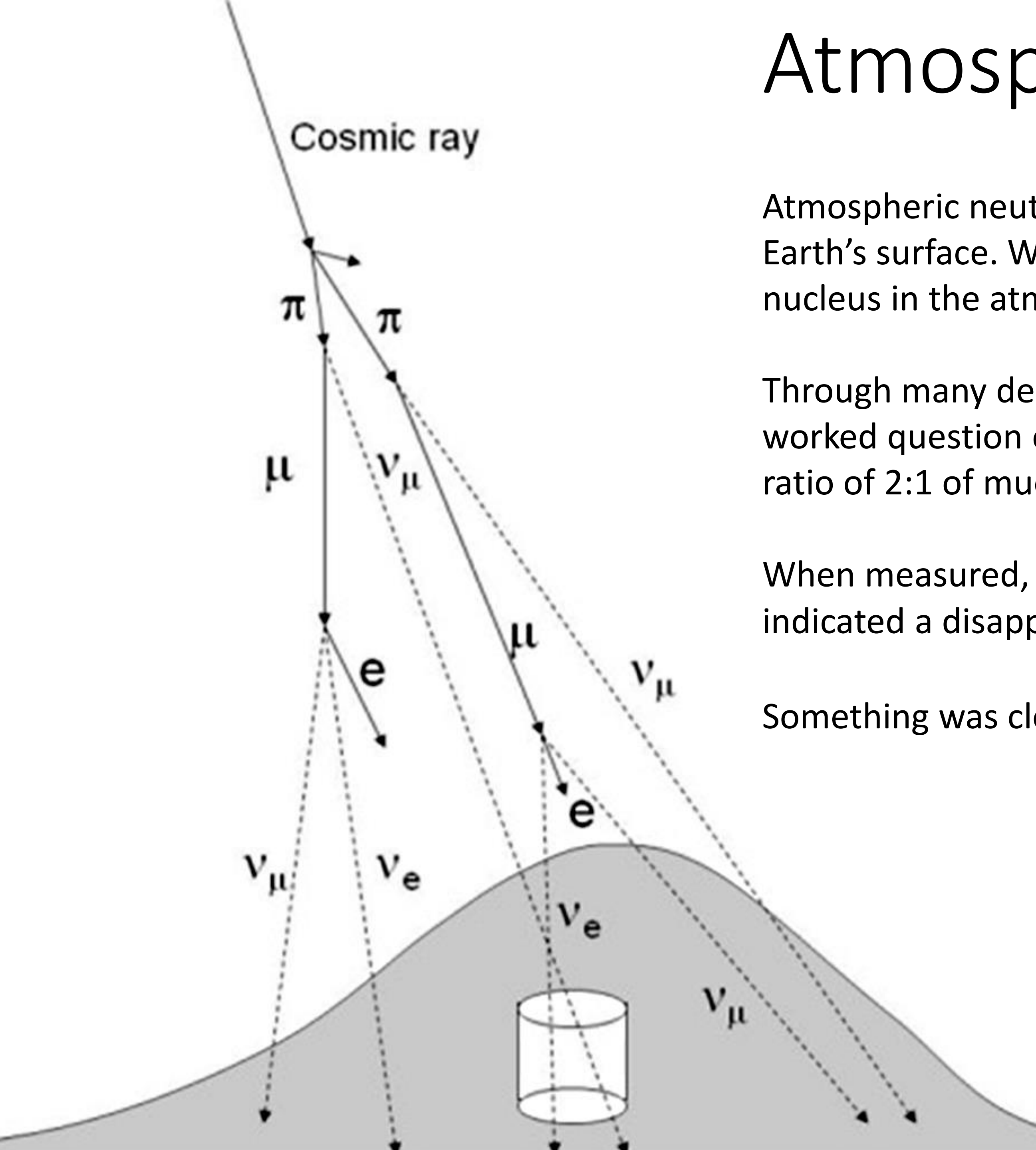
The idea was to fill a 380 cubic meter tank with perchloroethylene (dry-cleaning fluid), place it 1,478 meters underground to shield from cosmic rays, and count solar neutrinos (produced by fusion chains) by extracting and measuring the resulting argon.

With a week's operation, the experiment could expect to create 10 atoms of argon, compared to the 9×10^{30} chlorine atoms present in the tank.

However Davis did not find the expected 10 atoms, but instead 3. Many thought this to be experimental error, or an inaccurate prediction, but many later experiments confirmed this result.

This came to be known as the 'solar neutrino problem'.

Atmospheric neutrino problem



Atmospheric neutrinos are typically produced about 15 kilometres above the Earth's surface. When a cosmic ray (usually a proton) interact with an atomic nucleus in the atmosphere, it forms a cascade of particles.

Through many decays of the unstable products (here be a great relativity worked question on the lifetime of muons), neutrinos are produced with a ratio of 2:1 of muon:electron flavour.

When measured, this ratio was different to the prediction. Many experiments indicated a disappearance of ν_μ and an appearance of ν_e.

Something was clearly going on.

A possible solution?

Inspired by Kaon oscillations proposed by Gell-Mann and Pais, Bruno Pontecorvo proposed a possibility that another neutral particle, neutrinos, might oscillate between two states.

The proposals he made were;

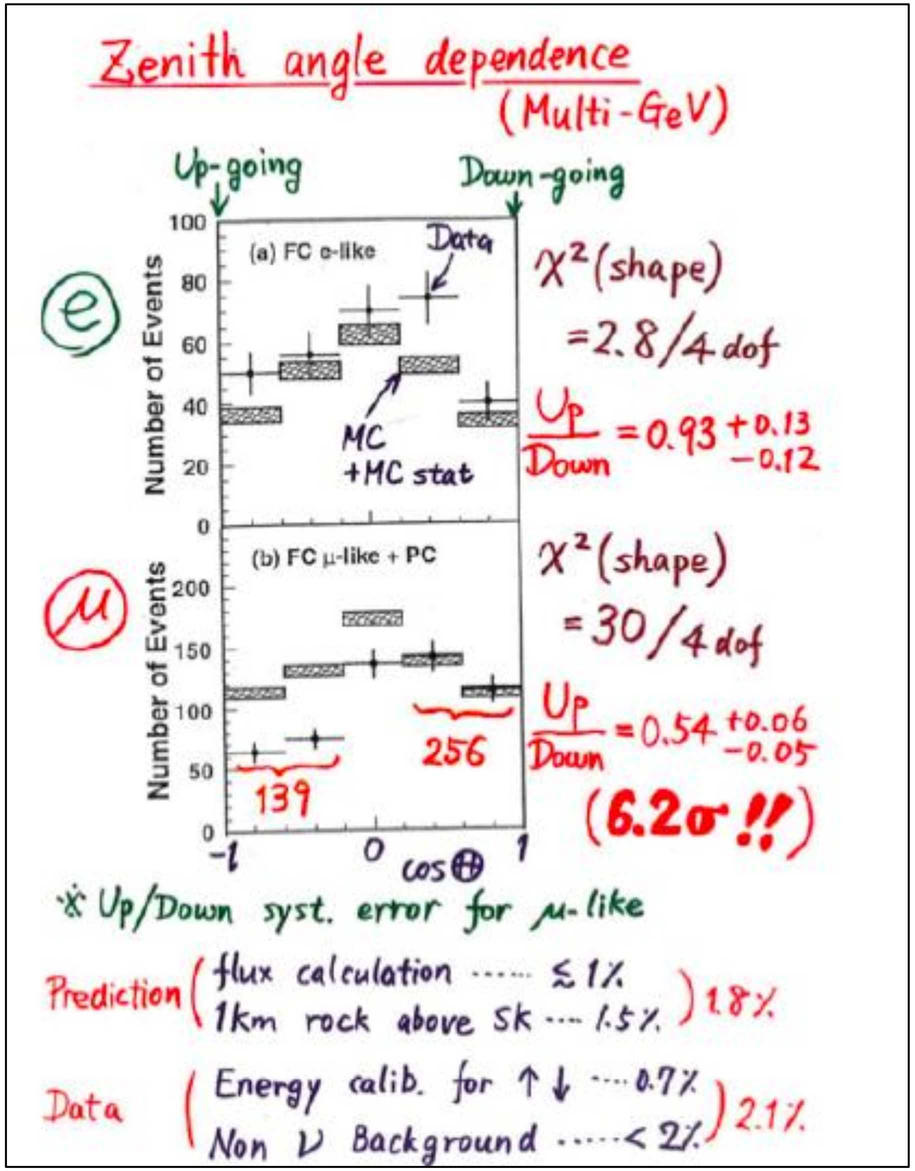
- $\nu_L \leftrightarrow \nu_R$ oscillations which were a bit of a shot in the dark.
- $\nu_L \leftrightarrow \bar{\nu}_L$ and $\nu_R \leftrightarrow \bar{\nu}_R$, which is a sterile oscillation.
- $\nu_e \leftrightarrow \nu_\mu$, allowing a neutrino to violate lepton flavour conservation, changing between the two lepton flavours known at the time.

These proposals can all be found in: Soviet Physics JETP. 26.984-988.

A complete confirmation

In 1998, the Super-Kamiokande water Cherenkov detector achieved the first measurement of atmospheric neutrino flavour disappearance, with the energy dependant pattern expected of oscillation phenomena.

On the right is shown the dependence on zenith angle (i.e. the angle the neutrino entered the detector from with from below being $\cos\Theta = -1$ and above $\cos\Theta = 1$) and therefore the distance travelled by the neutrino (from atmosphere above or through Earth as well).



Zenith angle dependence (Multi-GeV)

χ^2(shape) = 2.8/4 dof

$\frac{\text{Up}}{\text{Down}} = 0.93^{+0.13}_{-0.12}$

χ^2(shape) = 30/4 dof

$\frac{\text{Up}}{\text{Down}} = 0.54^{+0.06}_{-0.05}$

(6.2σ!!)

※ Up/Down syst. error for μ-like

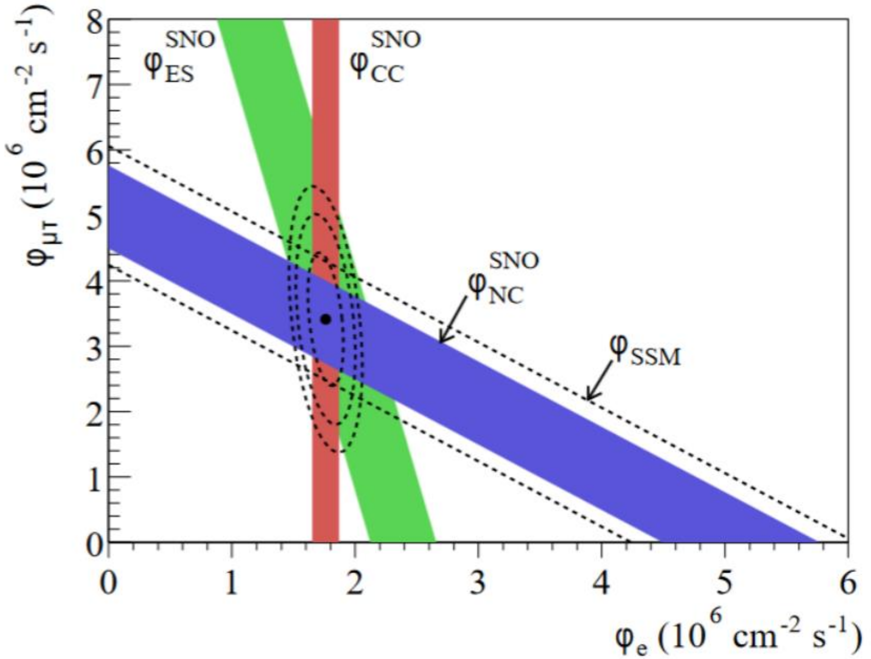
Prediction (flux calculation …… ≲ 1% ; 1km rock above SK ···· 1.5%) 1.8%

Data (Energy calib. for ↑↓ ···· 0.7% ; Non ν Background ·····< 2%) 2.1%

SNO place like home.

In 2002 the Sudbury Neutrino Observatory, a kiloton heavy water Cherenkov detector announced that it had measured solar neutrinos using three different interaction channels.

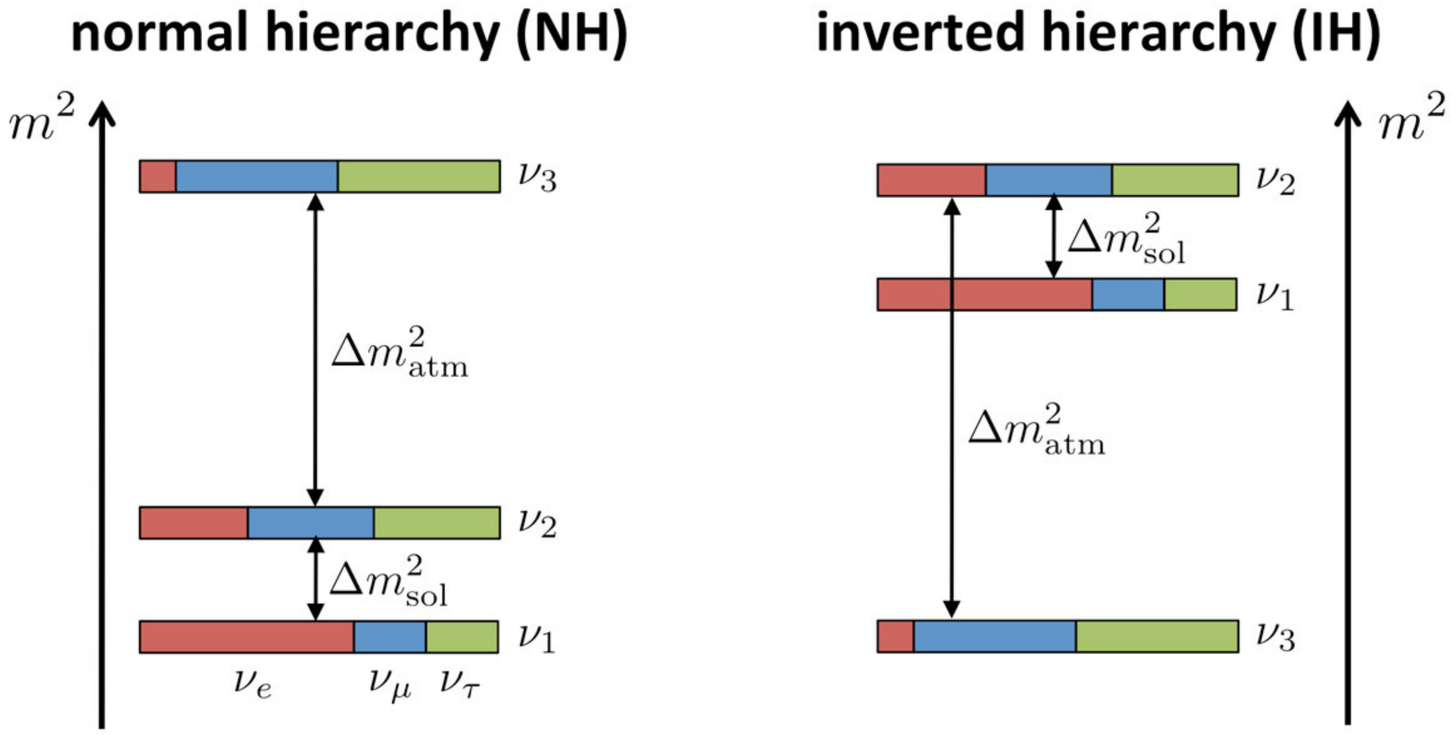
- The first, a charged current interaction, was sensitive to only ν_e. This saw the deficit Davis had at the Homestake experiment.
- The second, a neutral current interaction, was sensitive to all ν. This saw neutrino flux matching that predicted by the solar model.
- The third, an Elastic Scattering interaction, primarily sensitive to ν_e but had sensitivity to other flavours. This mostly gave directional info.



With these detections, the community was sure that neutrino flavour oscillations were occurring! However this realisation came with a troubling conclusion. Neutrinos had mass.

Well, kinda. What they actually required in theory was a mass difference between three, set, well defined states. And then each set neutrino flavour must be a transformation from these states, each flavour ending up as a mix of the three masses. This transformation is given by a rotation (contained in a matrix) from the mass states to the flavour states (or visa versa).

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$



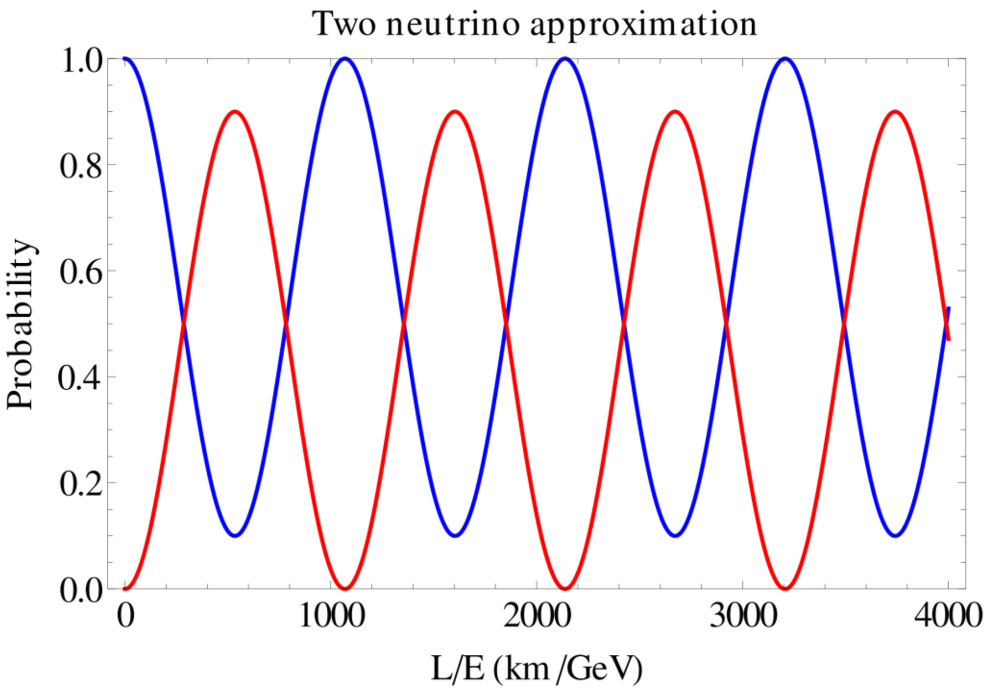
Warning,
maths
incoming.

If we start off in a two flavour world, this rotation matrix becomes familiar;

$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

And the survival probability of a flavour can be derived to be;

$$P(\nu_e \to \nu_e) = 1 - \sin^2 2\theta \sin^2 \left[\frac{(m_2^2 - m_1^2)x}{2E}\right]$$
$$= 1 - \sin^2 2\theta \sin^2 \left[\frac{1.27\Delta m^2 L}{E}\right]$$



From here we can see a couple of things;

- If the mass difference is 0, no oscillations can occur.
- Theta, named the mixing angle, defines how different the flavour states are from mass states. This determines amplitude of oscillation and is maximal at $\frac{\pi}{4}$.
- There are two free parameters that are either set for us by nature, or we can freely control;
  - ❖ Length, L, is the distance a neutrino will propagate.
  - ❖ Energy, E, is the energy of that neutrino.

$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{\mathrm{CP}}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{\mathrm{CP}}} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

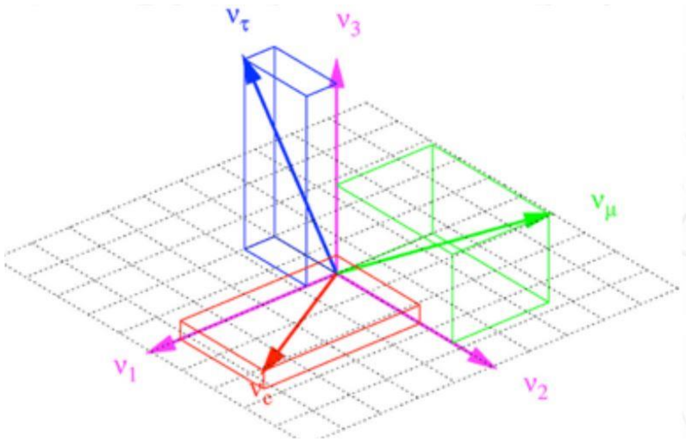
$$= \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{\mathrm{CP}}} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta_{\mathrm{CP}}} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta_{\mathrm{CP}}} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta_{\mathrm{CP}}} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta_{\mathrm{CP}}} & c_{23}c_{13} \end{bmatrix}.$$

$$P_{\alpha\to\beta} = \delta_{\alpha\beta} - 4\sum_{i>j} \mathrm{Re}\left(U^*_{\alpha i}U_{\beta i}U_{\alpha j}U^*_{\beta j}\right)\sin^2\left(\frac{\Delta m^2_{ij}L}{4E}\right) + 2\sum_{i>j} \mathrm{Im}\left(U^*_{\alpha i}U_{\beta i}U_{\alpha j}U^*_{\beta j}\right)\sin\left(\frac{\Delta m^2_{ij}L}{2E}\right),$$

With 3 flavour oscillations, things get more complex (no pun intended), but remain familiar;

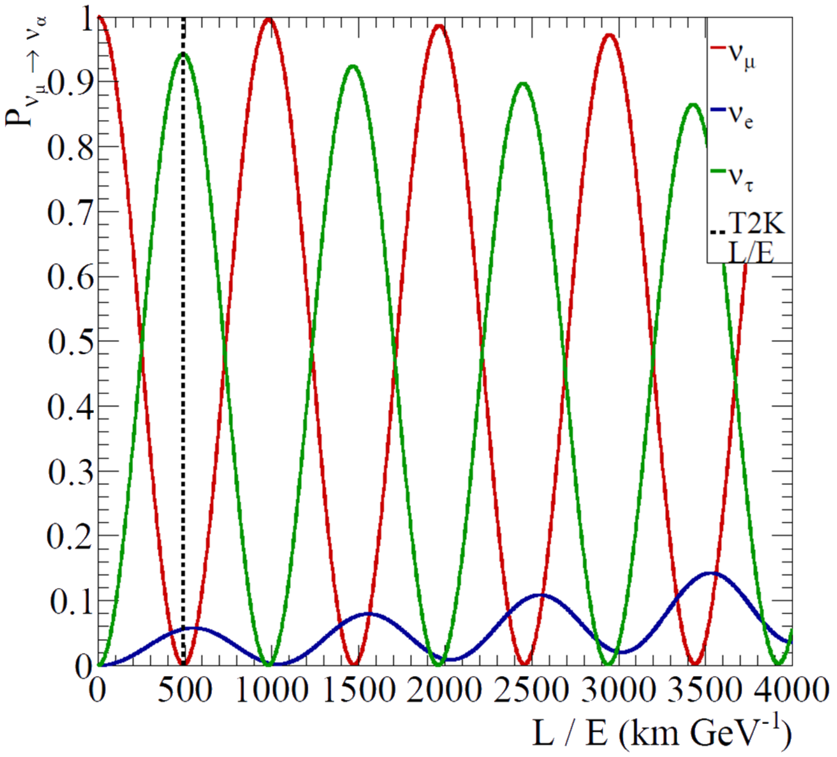
- There are now 2 mass differences, the large and the small; Δm^2_{32} & Δm^2_{21}. Note that because of the squared mass terms, the signs are not detectable without extra effects.
  - ❖ $\Delta m^2_{21} = 7.53 \times 10^{-5} eV^2$
  - ❖ $\Delta m^2_{32} \approx \Delta m^2_{13} = 2.45 \times 10^{-3} eV^2$
- There are 3 mixing angles; θ_{12}, θ_{23}, and θ_{13}. These values (shown below in terms of $\sin^2(\theta)$) are large, especially θ_{23} which is nearly maximal.
  - ❖ $\sin^2(\theta_{12}) = 0.307$
  - ❖ $\sin^2(\theta_{23}) = 0.547$
  - ❖ $\sin^2(\theta_{13}) = 0.0220$

The value for θ_{13} was found to be surprisingly large (though still smaller than the others) and opened up a whole new possibility; measuring the CP-violating phase factor, δ_{CP}.



Parameter values taken from PDG interactive listings, shown without error
https://pdglive.lbl.gov/Particle.action?node=S067&init=0

Assuming 3 flavour PMNS mixing in a pure ν_μ beam with a fixed baseline, L.

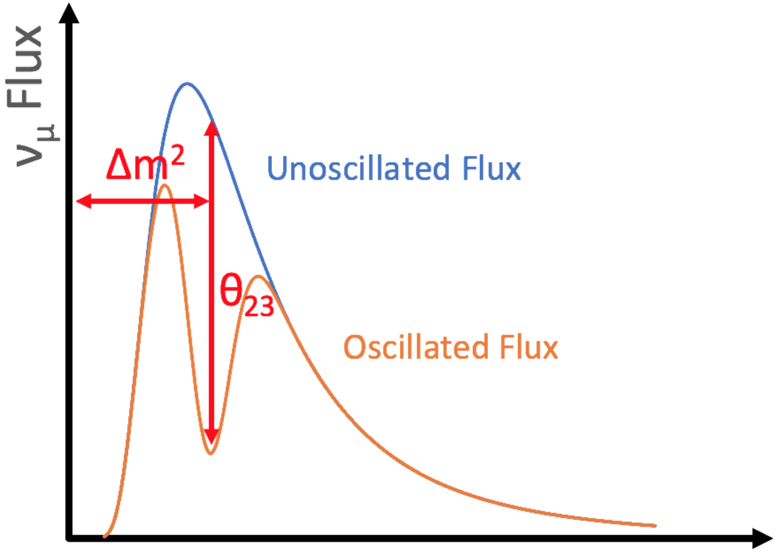


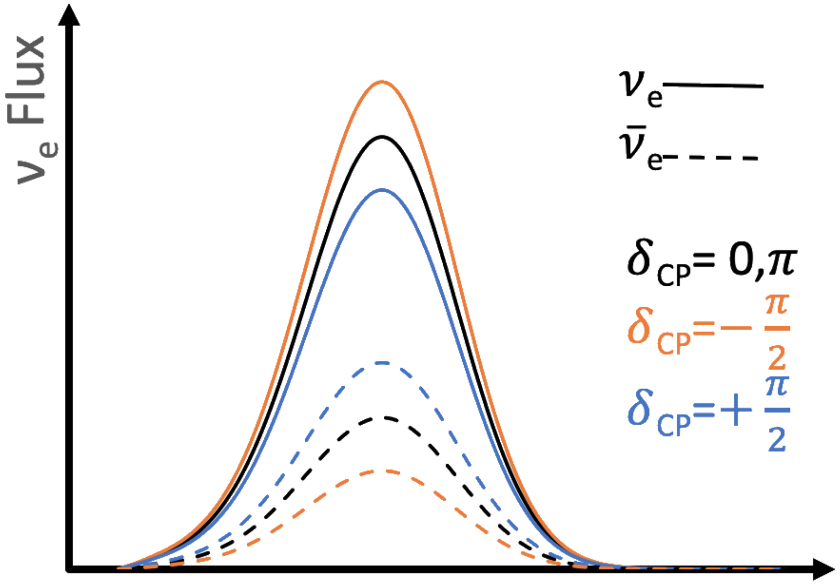
$$P(\nu_\mu \rightarrow \nu_\mu) \simeq 1 - \cos^4(\theta_{13}) \sin^2(2\theta_{23}) \sin^2\left(1.27 \Delta m^2_{32} \frac{L}{E_\nu}\right)$$

$$P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2(2\theta_{13}) \sin^2(\theta_{23}) \sin^2\left(1.27 \Delta m^2_{32} \frac{L}{E_\nu}\right)$$

$$\mp 1.27 \Delta m^2_{32} \frac{L}{E_\nu} 8 J_{CP} \sin^2\left(1.27 \Delta m^2_{32} \frac{L}{E_\nu}\right)$$

- θ_{23}
- Δm^2_{32} or $(|\Delta m^2_{31}|)$
- θ_{13}
- δ_{CP}





The international journal of science / 16 April 2020

nature

THE MIRROR CRACK'D

An indication of matter–antimatter symmetry violation in neutrinos

Coronavirus The models driving the global response to the pandemic

Hot source Remnants of primordial nitrogen in Earth's mantle

Origin of a species Revised age for Broken Hill skull adds twist to human evolution

CP violation would mean: $P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$

The magnitude of CP effect is given by the Jarlskog Invariant:

$$J_{CP}^{PMNS} = \frac{1}{8}\, sin2\theta_{12}\; sin2\theta_{13}\; sin2\theta_{23}\; cos\theta_{13}\; sin\delta_{CP}$$

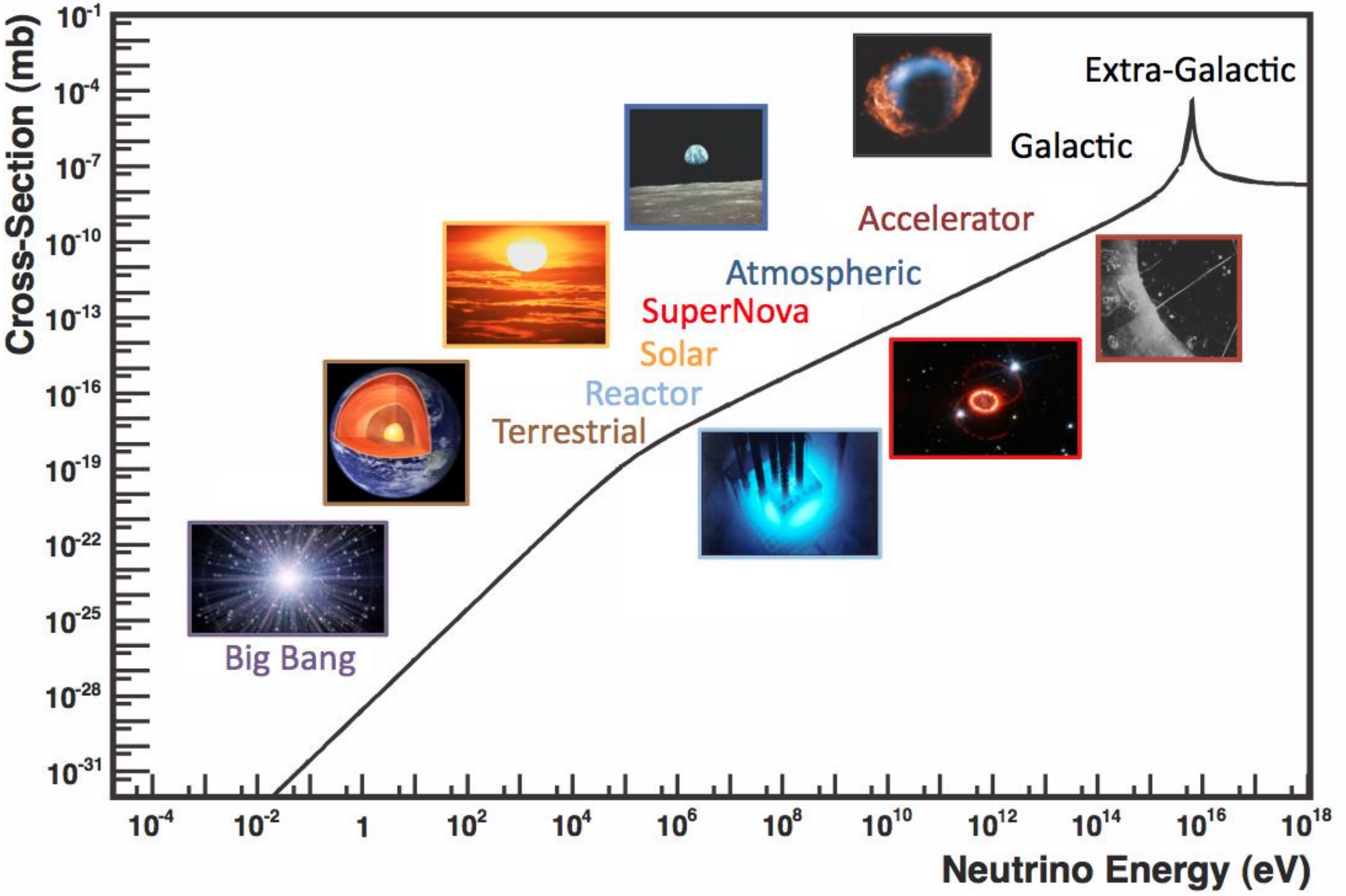
$$J_{CP}^{PMNS} = 0.035\; sin\delta_{CP}$$

$$J_{CP}^{CKM} \approx (3 \pm 1) \times 10^{-5}$$

Indications that the value of δ_{CP} is actually near maximal, and thus leptonic CP violation is large. This could explain the matter-antimatter asymmetry!

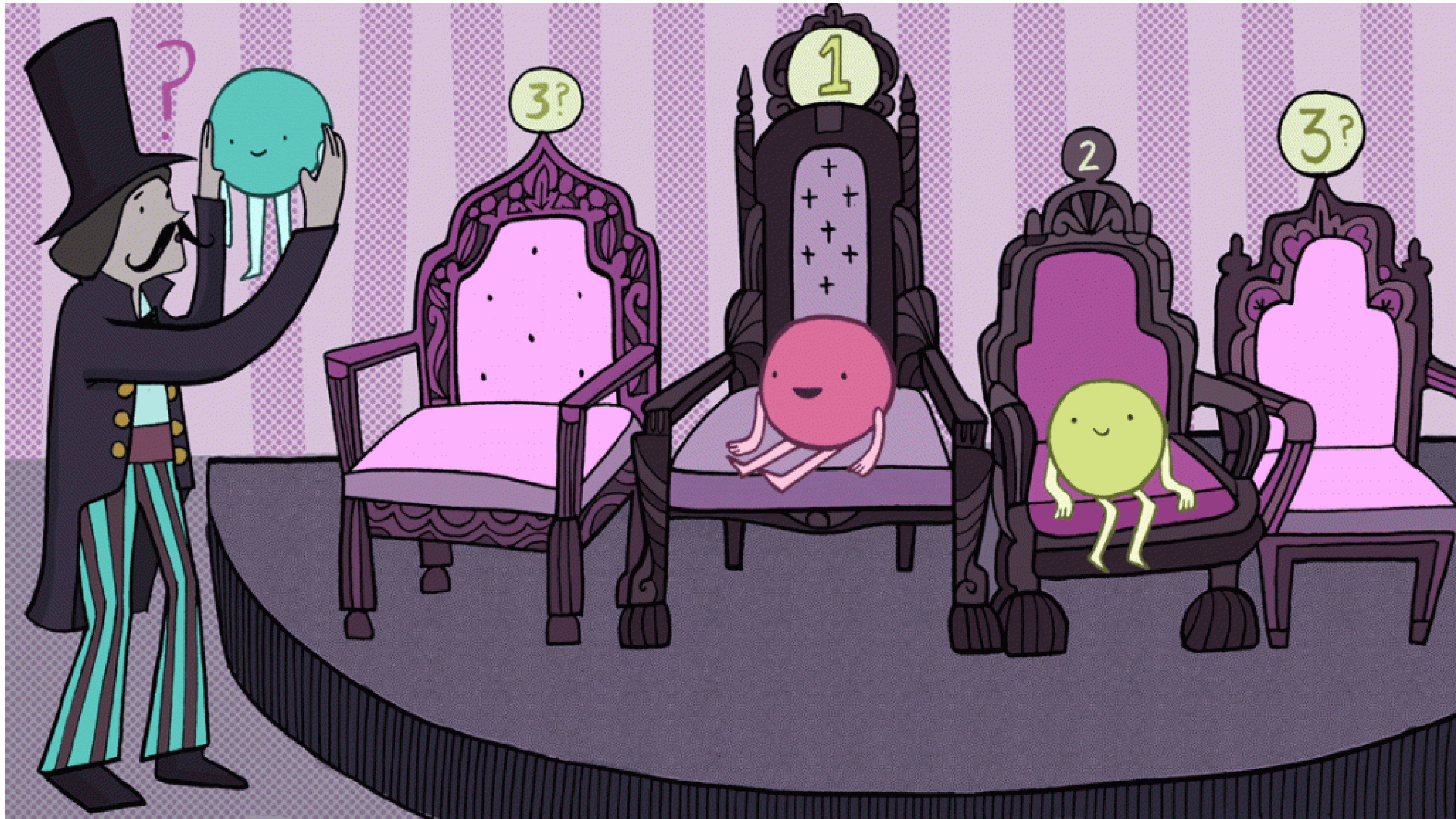
And this is perhaps only one of the most famous questions…

- Our neutrino detection capabilities span an incredible range of energies, with many neutrino experiments trying to achieve many different measurements.

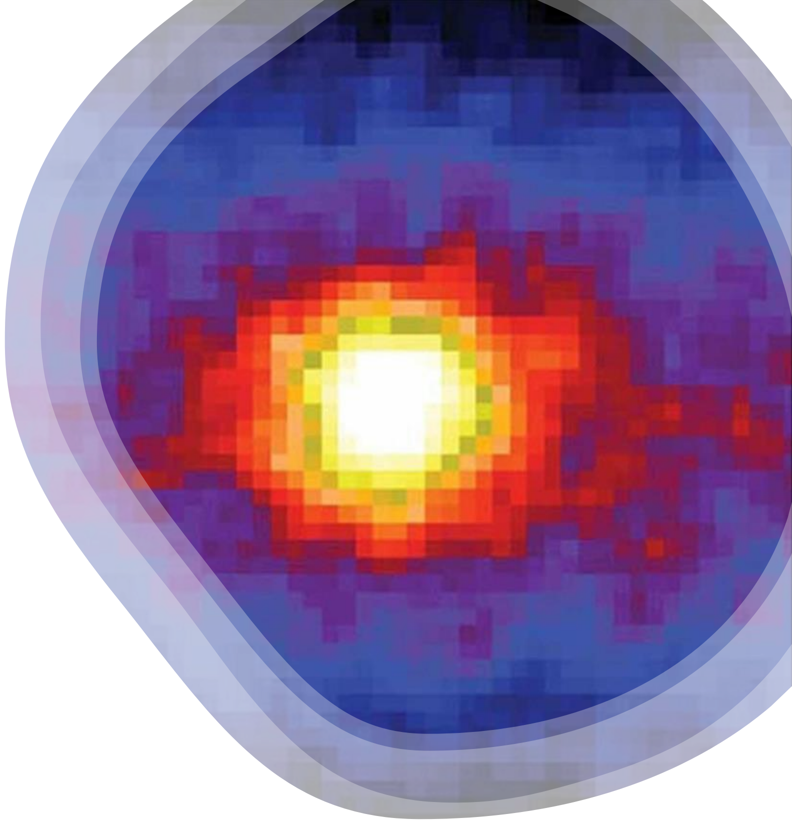


Other open questions

- Neutrinos have mass, where does this come from? (Neutrinoless double beta decay)
- Which order are these masses in? This is called neutrino mass ‘hierarchy’. (Matter effects)
- Is there another neutrino which does not interact like the rest? (Short baseline experiments)
- Is θ_{23} maximal or which side of $\frac{\pi}{4}$ does it lie.
- What is a neutrino’s speed? Looking at you OPERA.



SNEWS

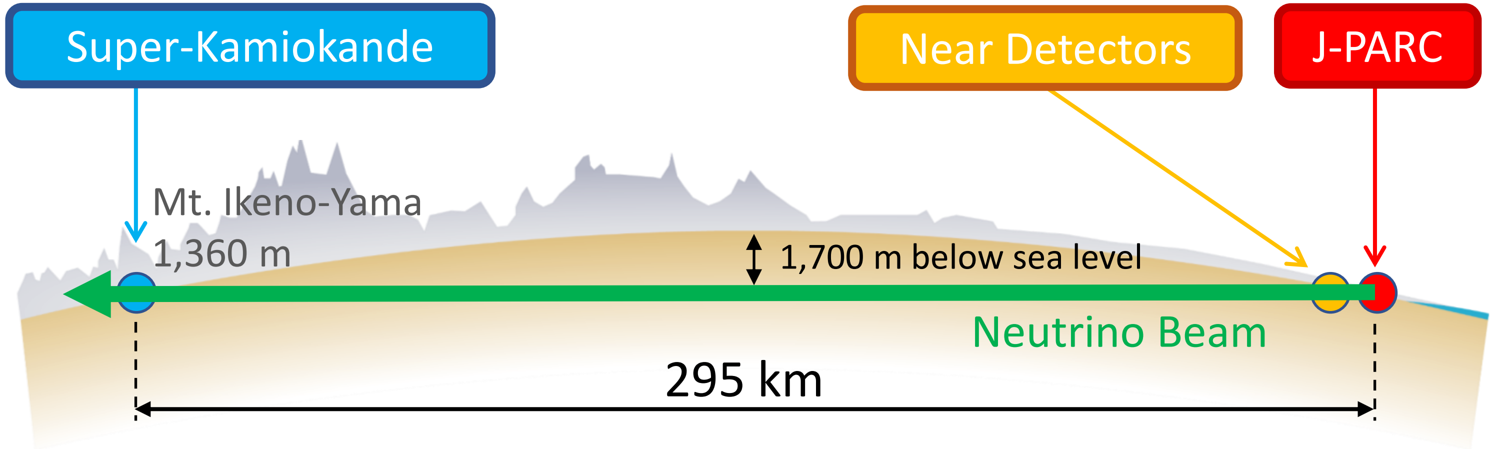


Neutrino Astronomy

Is it not a strange fate that we should suffer so much fear and doubt for so small a thing?- Boromir

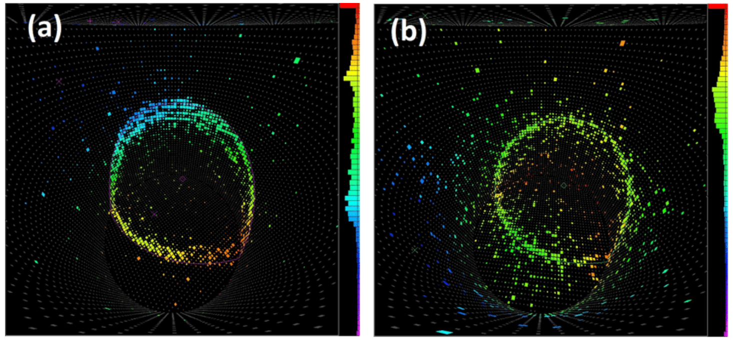
BACKUP

T2K (Tokai to Kamioka) is a long-baseline neutrino experiment that utilises multiple detectors in the goal of measuring the properties of neutrinos and their oscillations.



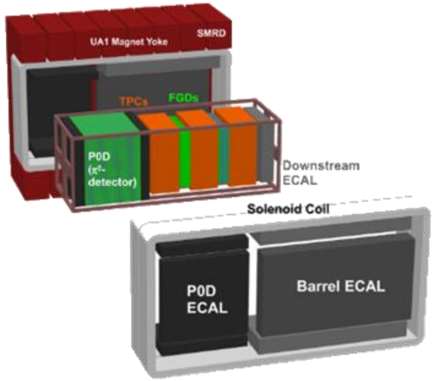
Super-K

- Off-axis water-based Cerenkov far detector.
- Topology based PID.
- CCQE dominant interactions.



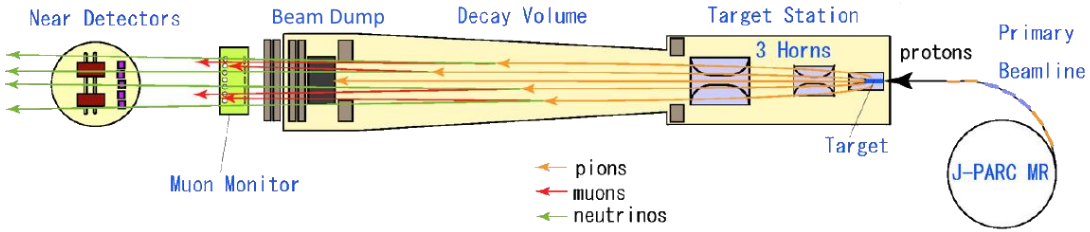
Cerenkov rings detected at SK. (a) is a muon event, (b) is an electron event.

ND280

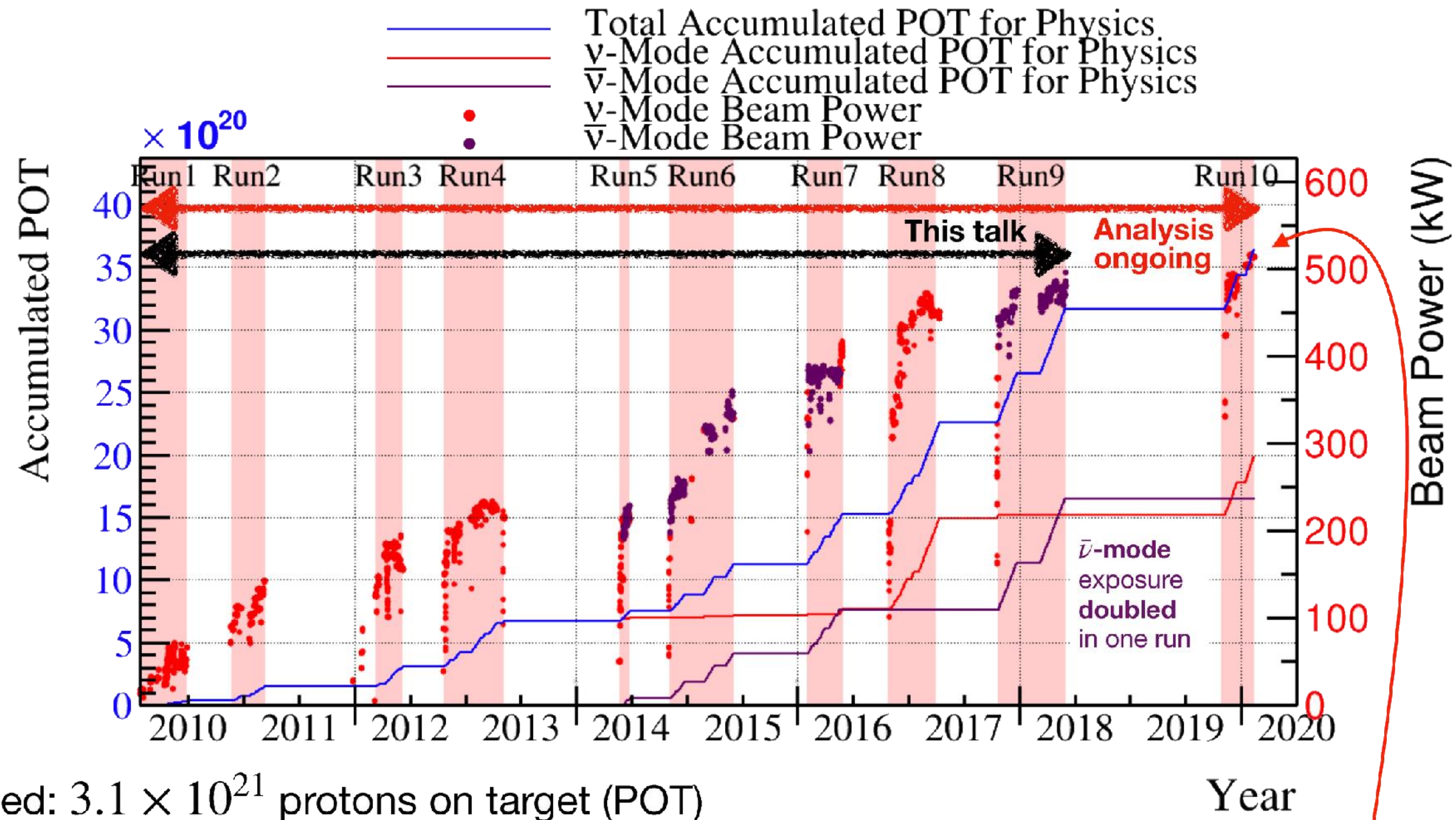


- Magnetised composite near detector.
- Off-axis (replicates SK energy spectra).
- Constrains flux and cross-section uncertainties.

J-PARC



- ‘Off axis’ beam tuned to 0.6 GeV for oscillation max at SK.
- Produces pure $\nu_\mu/\bar{\nu}_\mu$ flux.
- Able to be run in ν or $\bar{\nu}$ mode.



This talk

Analysis ongoing

$\bar{\nu}$-mode exposure **doubled** in one run

Analyzed: 3.1×10^{21} protons on target (POT)

ν-mode : $\bar{\nu}$-mode ~ 50 : 50

515 kW operation achieved recently!

33% increase of ν-mode data in upcoming analysis.

ND280 detector suite

- Same off-axis angle as Super-K (2.5 degrees).
- Measures ν_μ and ν_e spectrum before the oscillation → TPCs + FGDs
- Measure background processes to oscillation (NCπ0, NC1π, CC1π…)
- Compare Carbon and Oxygen interactions (FGD2 and P0D)

ND280 installed in ex-UA1 magnet (0.2 T) 3.5x3.6x7.3 m

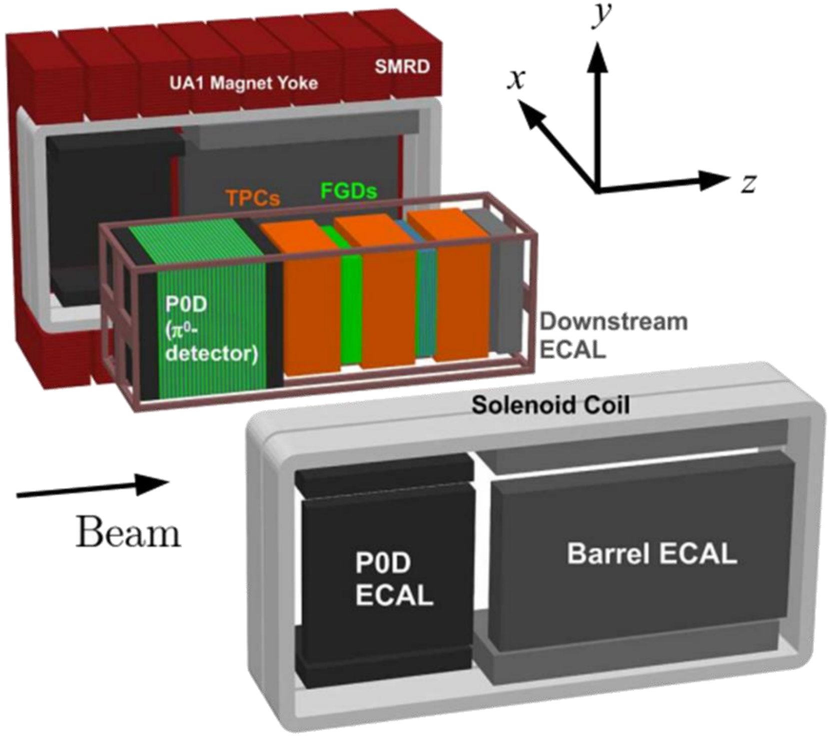
SMRD (Side Muon Range Detector): scintillator planes in magnet yokes. Measure high angle muons

2 FGDs (Fine Grained Detector): active target mass for the tracker, optimized for p/π separation Carbon+Water target in FGD2

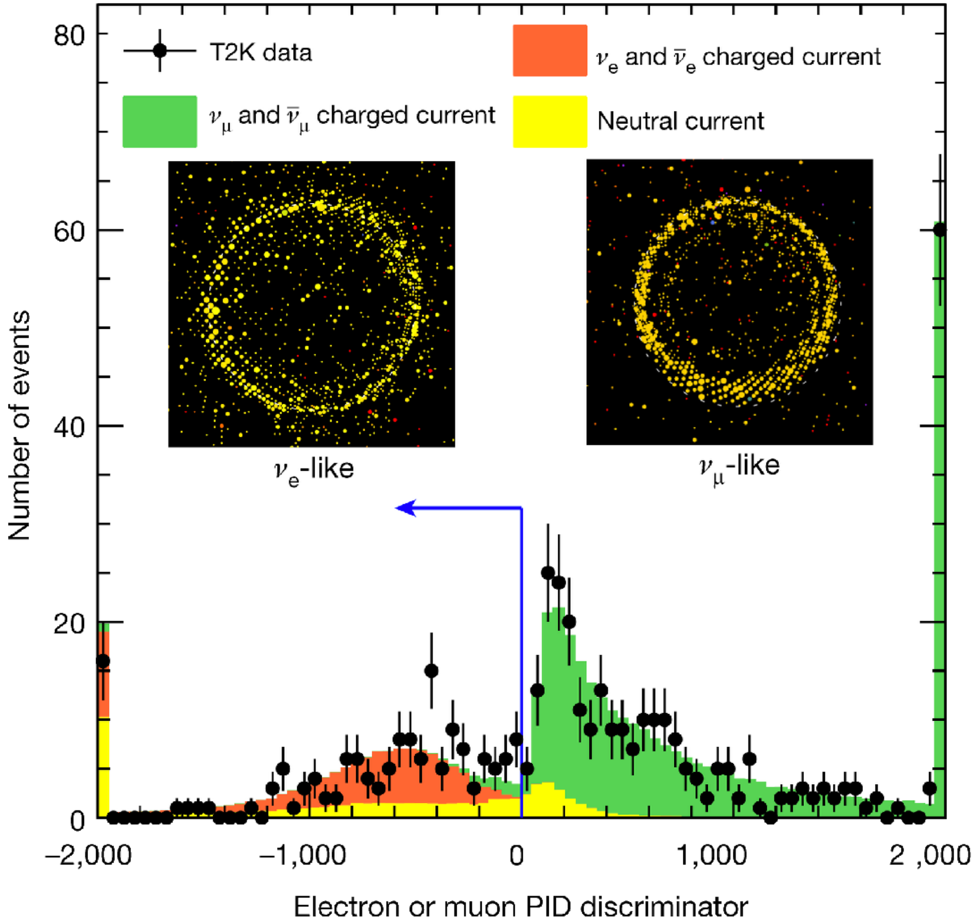
P0D (π0 detector): scintillator bars interleaved with fillable water target bags and lead and brass sheets. Optimised for γ detection

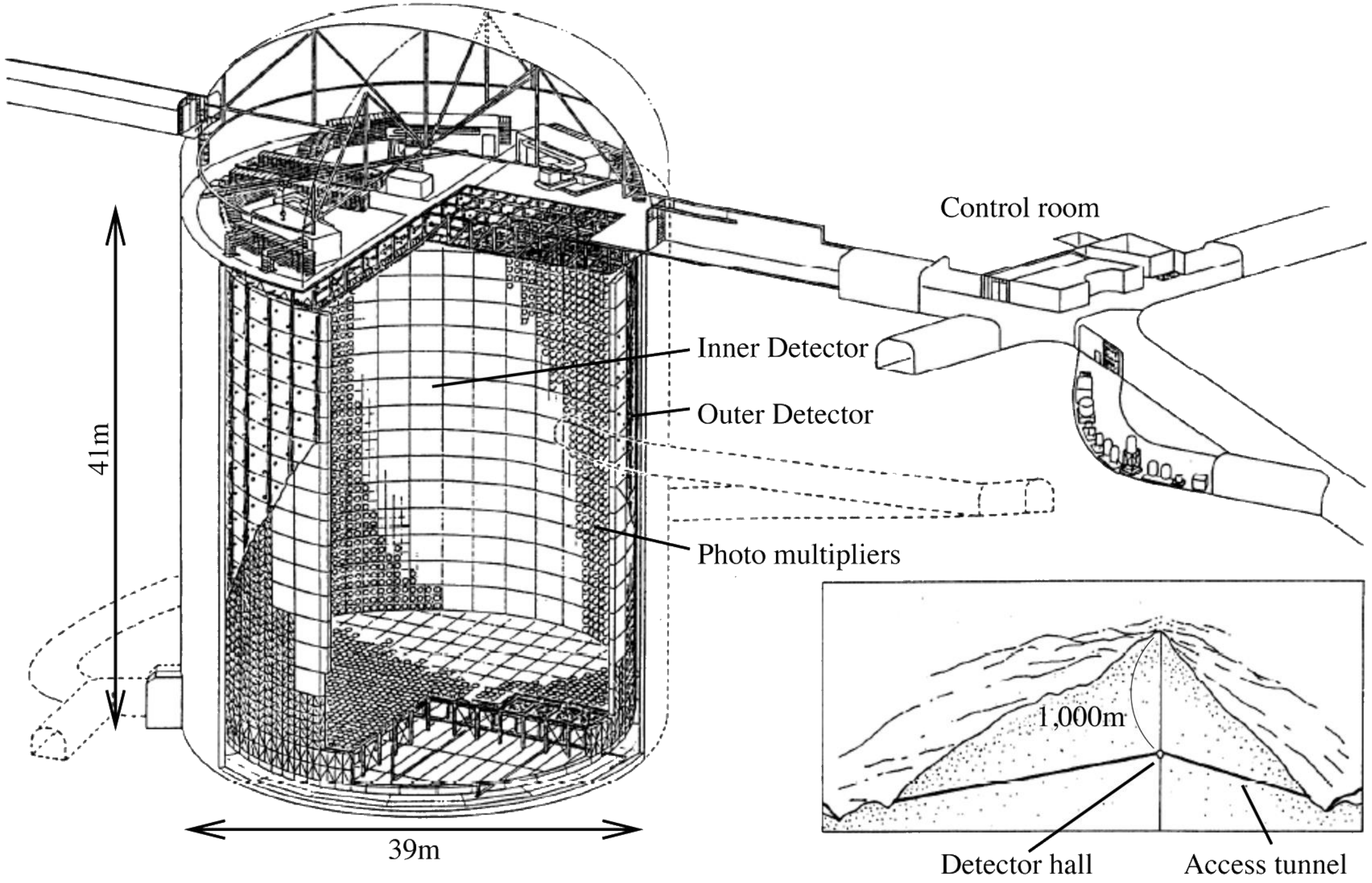
3 TPCs (Time Projection Chambers): measure momentum and charge of particles from FGD and P0D, PID capabilities through dE/dx

P0D, Barrel and Downstream ECAL: scintillator planes with radiator to measure EM showers



50kton water Cerenkov detector.
~11,000 20” PMTs
Vertex reconstruction
Mis-ID of less than 1%.
Super-K is located 1,000 m (3,300 ft) underground in the Mozumi Mine in Hida's Kamioka area.



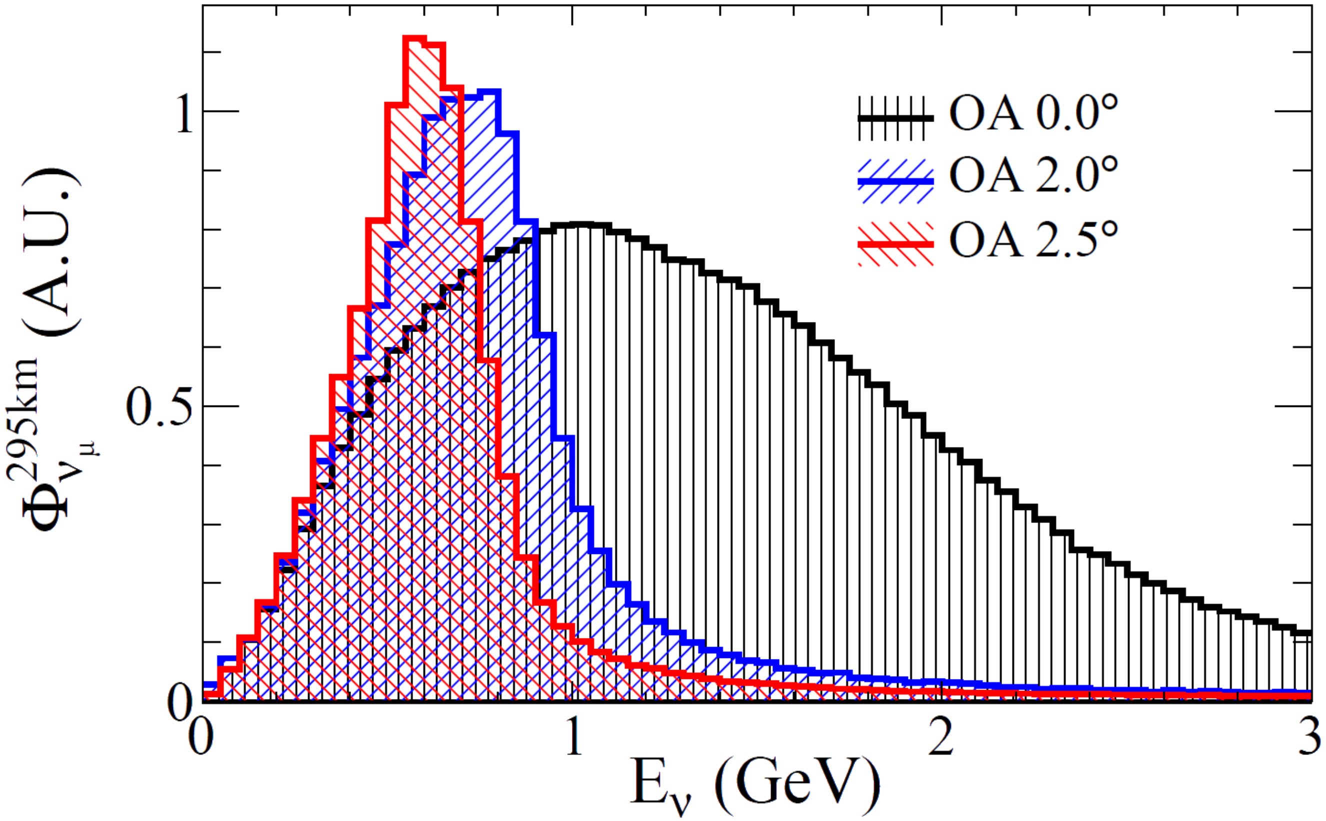


https://www-sk.icrr.u-tokyo.ac.jp/realtimemonitor/

Disclaimer: Almost all events are cosmic ray muon events.

$$E = \frac{E^*}{\gamma(1 - \beta \cos\theta)}$$

E* is neutrino energy in centre of mass frame of the decaying meson.
γ is the neutrino's Lorentz factor.
β is the neutrino's Lorentz velocity.
θ is the neutrino's angle in the lab frame.



T2K is the first experiment in which the off-axis concept was implemented. This decreases the amount of neutrinos at high energies (decreasing more complicated interaction types and also tightening flux around osc max).

Speaking of constraints…

These are official results that mirror those released at Neutrino 2020, and are from our T2K internal tech note.

- Binned log-likelihood method compares predicted and observed event spectra over parameter space.
- Systematics (and nuisance oscillation parameters) are marginalised over using their prior constraints.
- This leaves us with a likelihood dependant only on parameters of interest.
- Confidence intervals are constructed using const. ΔX^2 (left) or Feldman-Cousins (right).