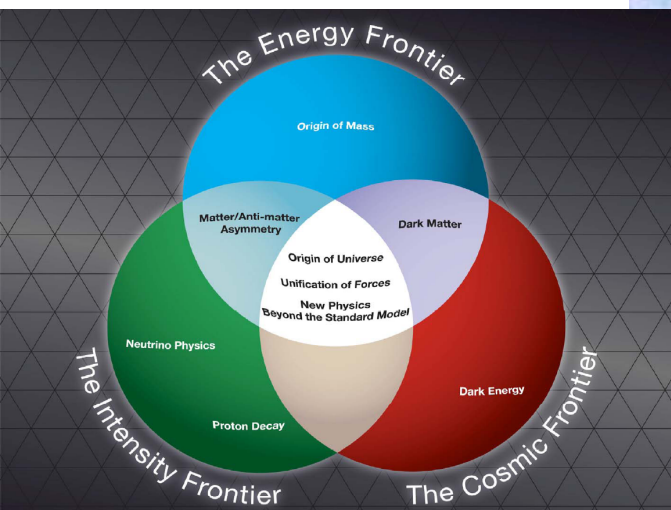


*II Workshop on MUON Precision Physics 2023 (MPP2023), Liverpool, 7-10 November 2023*

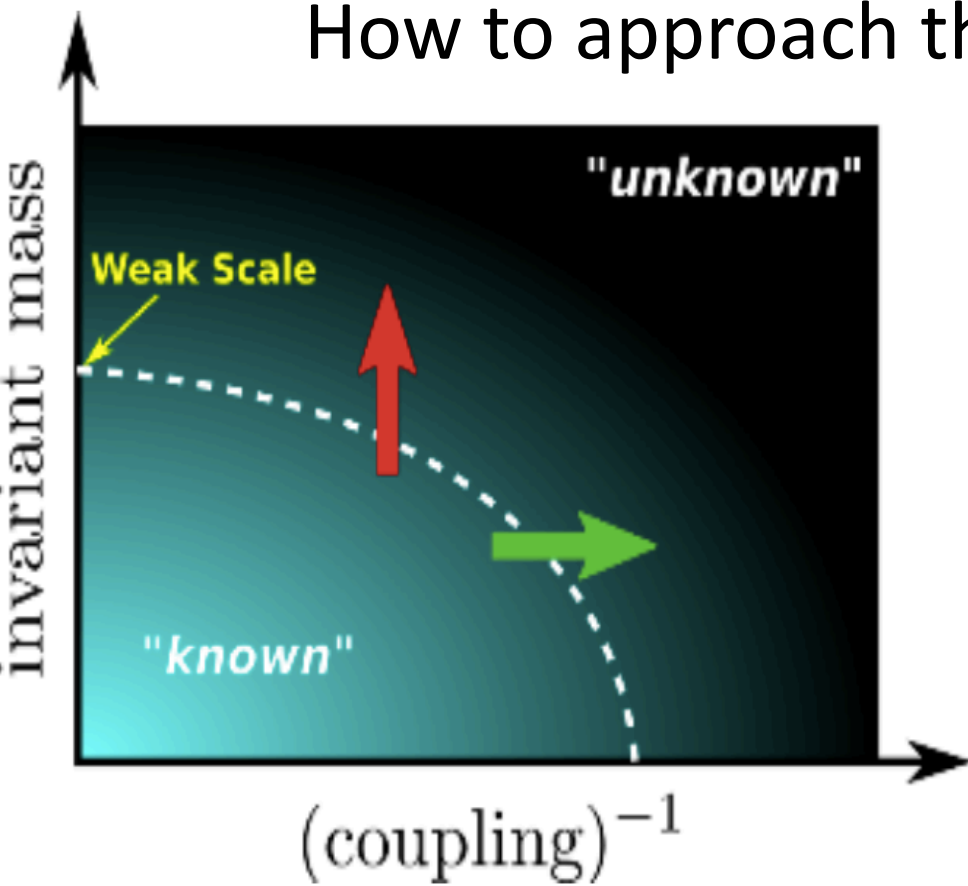


# The *Precision Frontier* of Particle Physics in 2023

Antonio Masiero

Univ. of Padova and INFN, Padova

# How to approach the “Unknown”

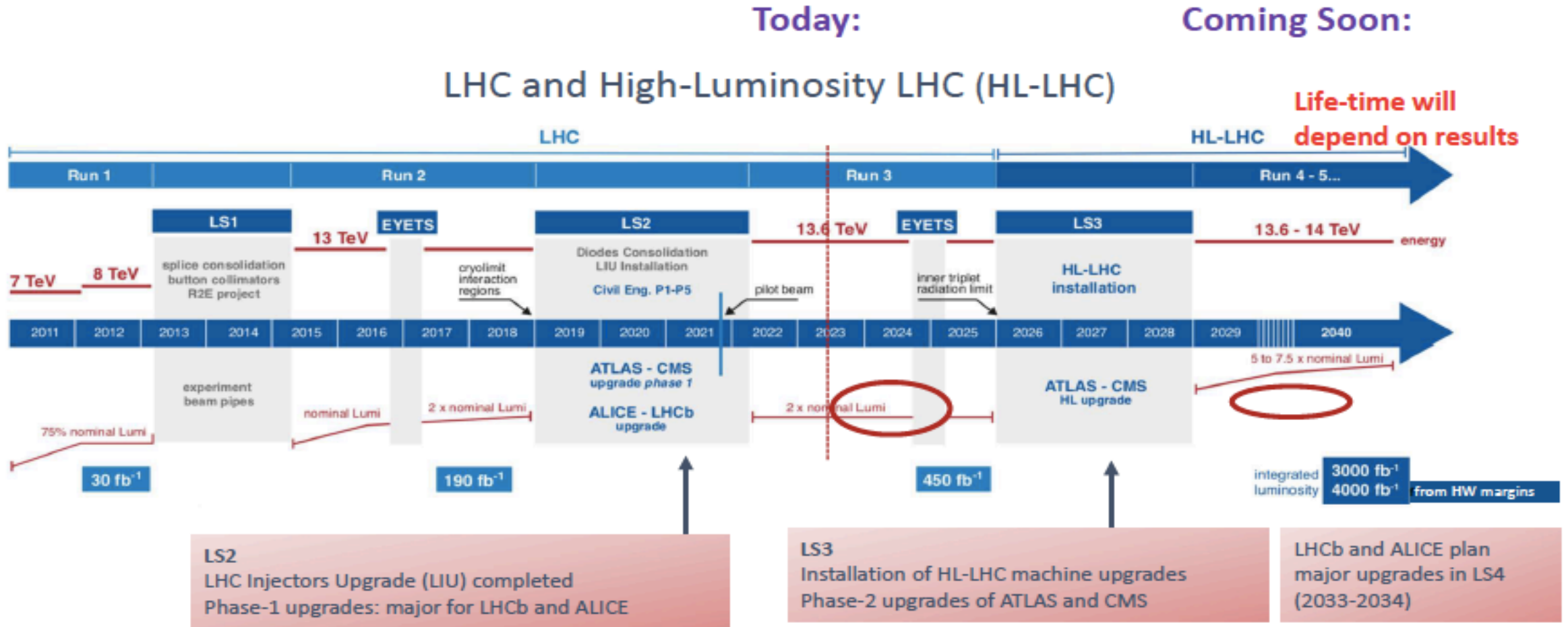


**High-Energy Frontier** →  
**New Physics (NP) at large energy scales**

**Next generation Astro-Particle experiments** (DM, DE, neutrinoless  $\beta\beta$ -decay, observational cosmology, multimessenger and multiwavelengths physics)  
→ **exploring large distances and times, feebly or very feebly coupled NP**

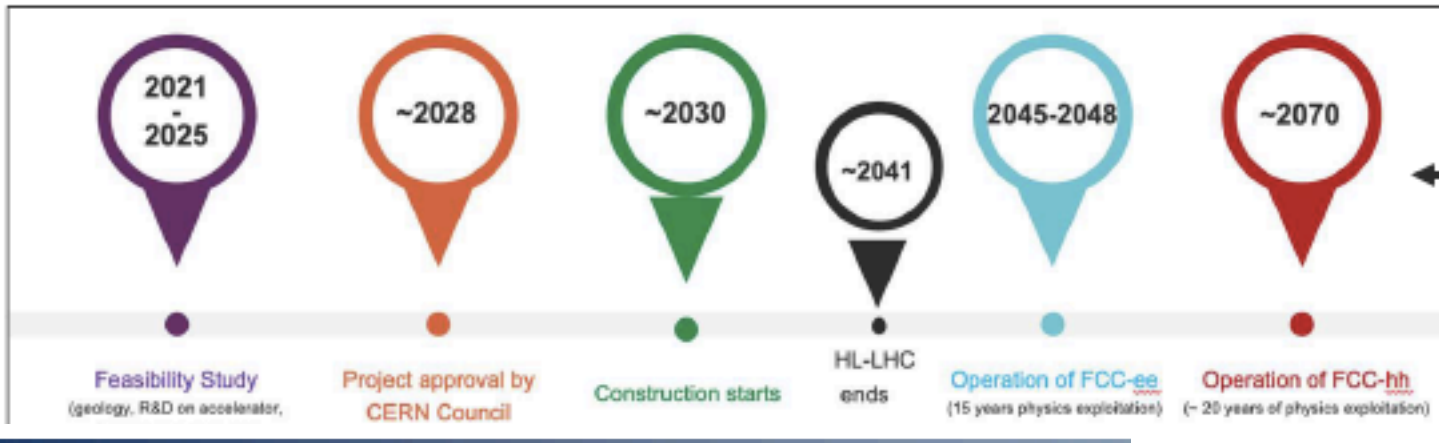
To proceed along the above two frontiers (in particular the HE one) → **LARGE-SCALE** experiments and Research Infrastructures on earth or in space (HE particle colliders, telescopes, large volume detectors) demanding increasingly **LARGER TIME-SCALES** (from the proposal to the operational phase) and **LARGER COSTS**

# The HE particle collider road: past, present and future



Expected integrated luminosity at the end of LHC (2025): > 450 fb<sup>-1</sup> (design target: 300 fb<sup>-1</sup>)

Luminosity target for HL-LHC: 3000 fb<sup>-1</sup> - "needed to observe HH production at ~ 5σ level in ATLAS and CMS"

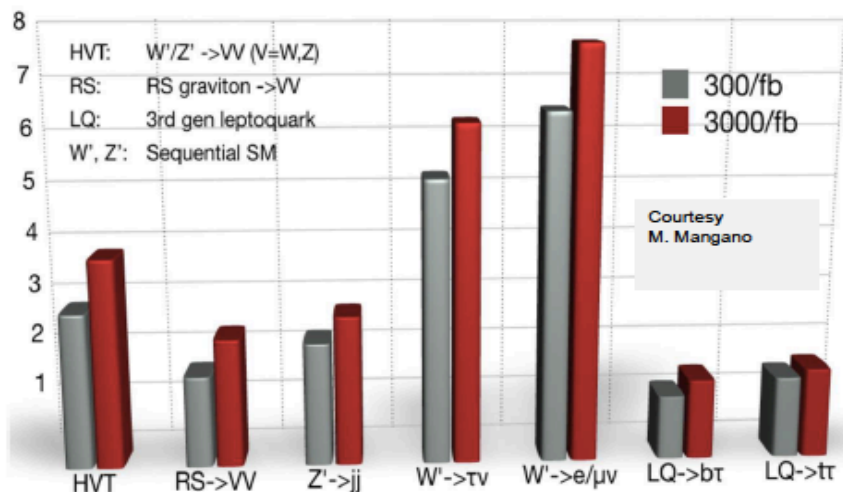


Realistic schedule takes into account:

- CERN Council approval timeline
  - past experience in building colliders at CERN
  - that HL-LHC will run until ~ 2041
- **ANY future collider at CERN cannot start physics operation before 2045-2048** (but construction will proceed in parallel to HL-LHC operation)

## Physics potential of HL-LHC 2

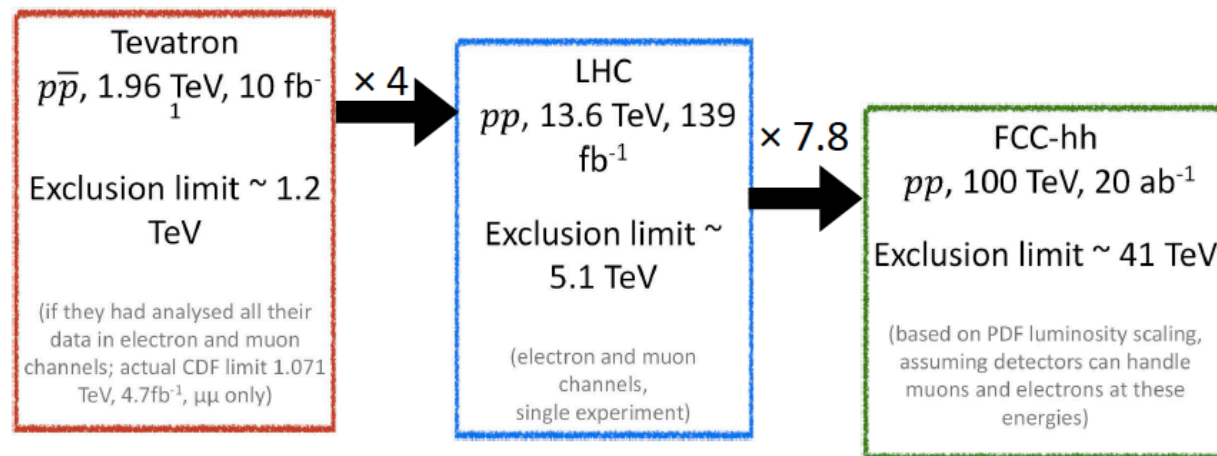
5 $\sigma$  discovery mass reach (TeV) for new particles:



## Beyond SM Searches

Following Gavin Salam, look at  $Z'_{SM}$  as a simple measure of progress

- perhaps not very "exciting", but simple and most experiments look for it



# Complementary (*not* ALTERNATIVE!) approach → **HIGH-PRECISION EXPS. in SMALL/MID-SCALE RIs**

Low-energy high-precision expts. can exploit :

- many recent *advances in experimental techniques and technologies* + (experimental as well as theoretical) *synergies* with adjacent areas of particle physics (atomic, molecular, optical, nuclear, particle physics)
- the relevant impact of *quantum mechanical virtual effects* on physical phenomena → access to the exploration of BSM new physics areas (large energy scales, very feebly coupled new particles, hidden sectors, etc.) difficult to be probed by traditional HE particle physics

**SYNERGY between small/mid-scale & large-scale experiments → casting a wider and tighter net for possible effects of BSM physics**

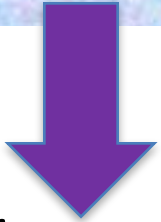
**Community Planning Exercise: Snowmass 2021 Blum, Winter et al.  
arXiv:2209.08041v2**



# The (intertwined) precision physics in this talk

- **Leptonic magnetic dipole moments**  $(g-2)_l$   $l = e, \mu, \tau$
- **Electric dipole moments (EDMs)**
- **Lepton Flavour Universality (LFU)**
- **Charged Lepton Flavour Violation (CLFV)**

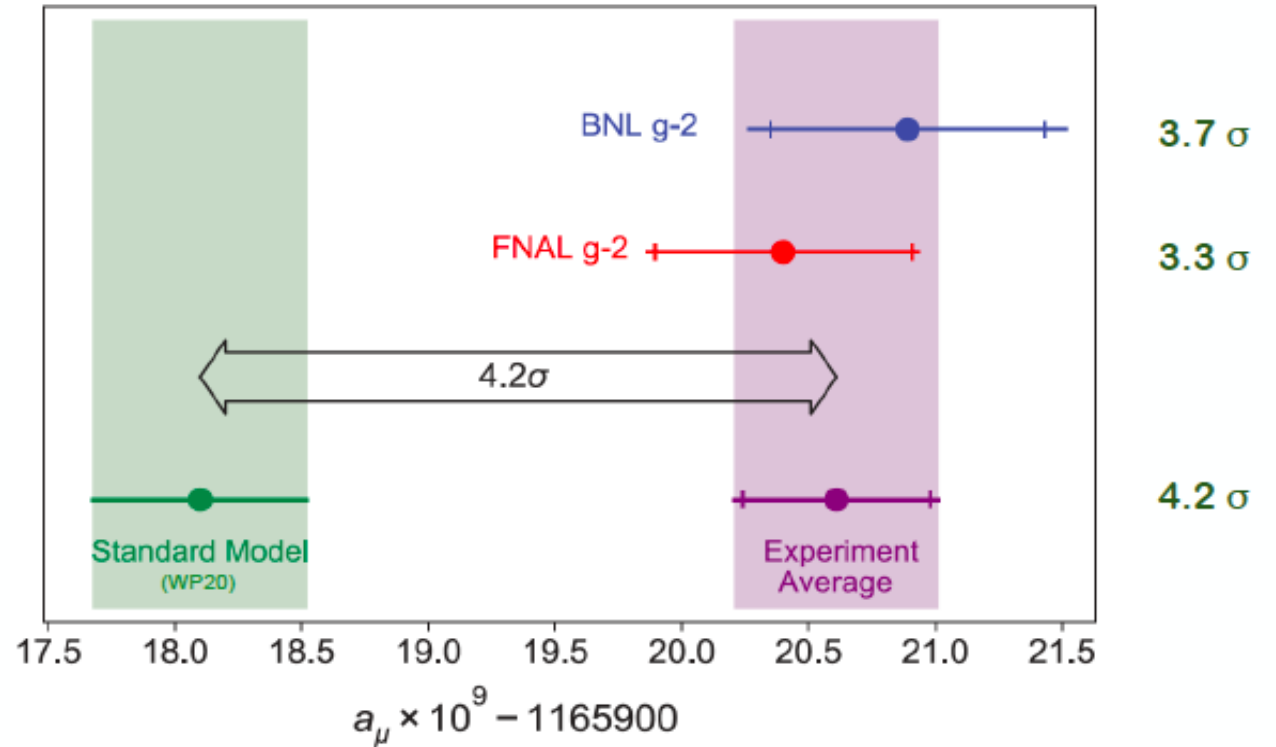
# The OLD $(g-2)_\mu$ puzzle



3-4 $\sigma$  discrepancy between exp. result and SM prediction

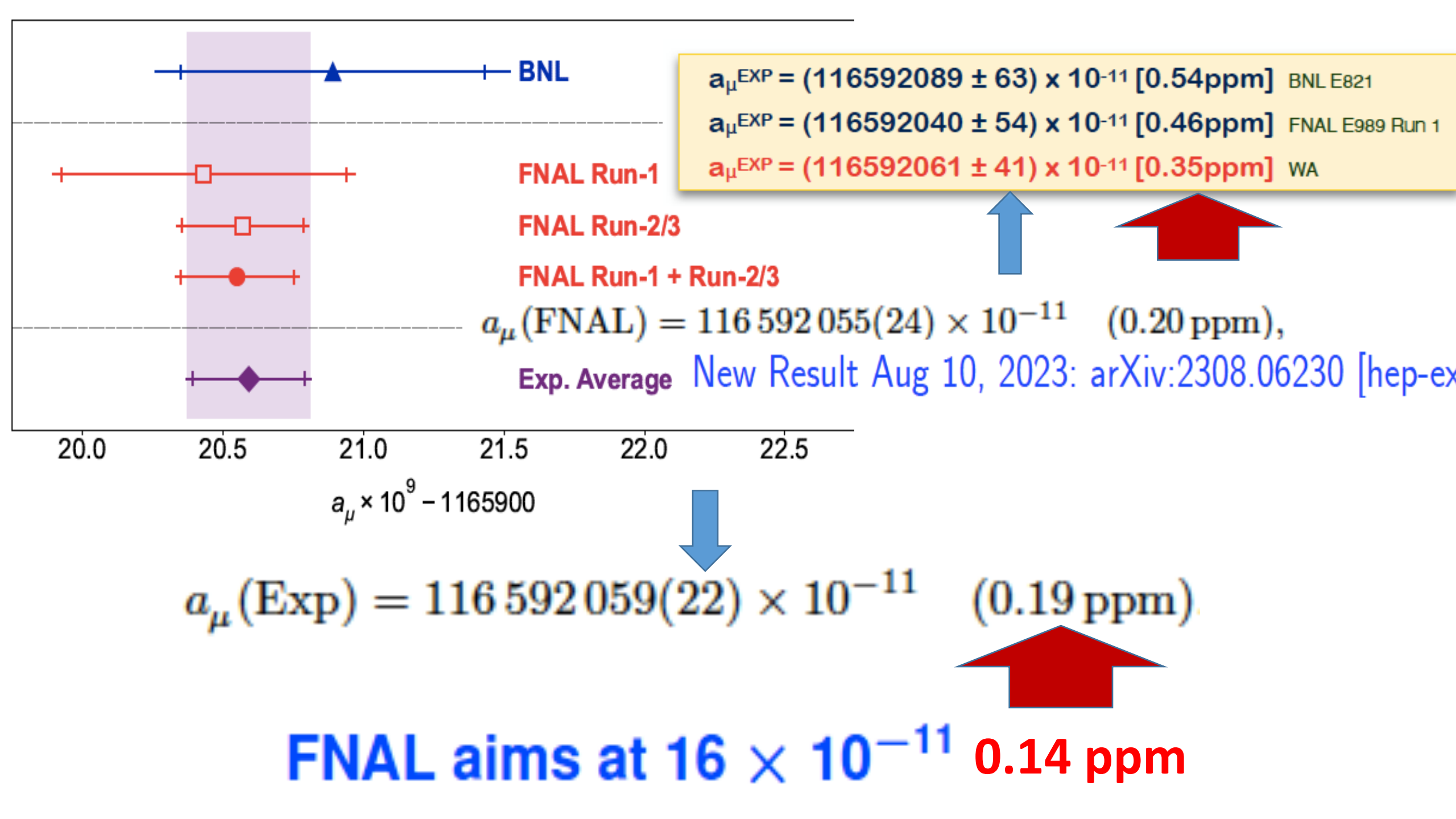
**CONSTANTLY** with us for a couple of decades!

(and last August it even became  $> 5\sigma$  comparing the exp. result with the SM expectation in the White Paper WP20 by the Muon g-2 Th. Initiative)



$a_\mu^{\text{EXP}} = (116592089 \pm 63) \times 10^{-11}$ [0.54ppm]	BNL E821
$a_\mu^{\text{EXP}} = (116592040 \pm 54) \times 10^{-11}$ [0.46ppm]	FNAL E989 Run 1
$a_\mu^{\text{EXP}} = (116592061 \pm 41) \times 10^{-11}$ [0.35ppm]	WA

- FNAL aims at  $16 \times 10^{-11}$ . First 4 runs completed, 5th in progress.
- Muon g-2 proposal at J-PARC: Phase-1 with  $\sim$  BNL precision.



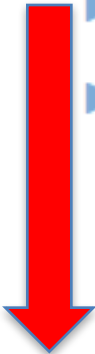


# NEW PHYSICS for the muon $g-2$ : at which scale?

$$\Delta a_\mu \equiv a_\mu^{\text{NP}} \approx (a_\mu^{\text{SM}})_{\text{weak}} \approx \frac{m_\mu^2}{16\pi^2 v^2} \approx 2 \times 10^{-9}$$


- ▶ A weakly interacting NP at  $\Lambda \approx v$  can naturally explain  $\Delta a_\mu \approx 2 \times 10^{-9}$
- ▶  $\Lambda \approx v$  favoured by the *hierarchy problem* and by a WIMP DM candidate.

On the other hand, HE experiments (LEP, Tevatron, LHC) have NOT provided any clue for the presence of new (charged) particles at the ELW. scale

- 
- ▶ NP is very light ( $\Lambda \lesssim 1$  GeV) and feebly coupled to SM particles.
  - ▶ NP is very heavy ( $\Lambda \gg v$ ) and strongly coupled to SM particles.

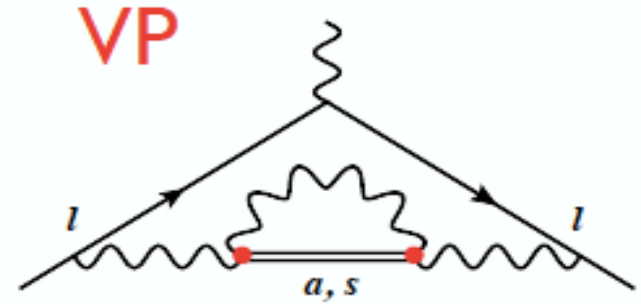
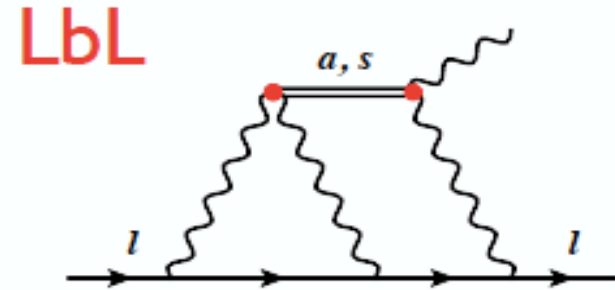
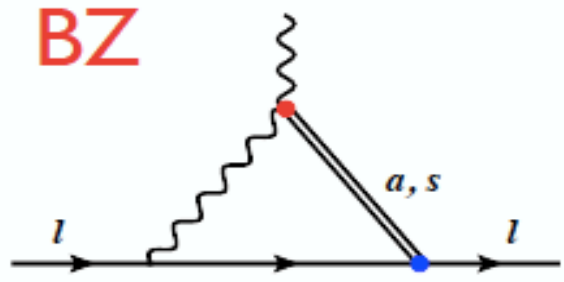
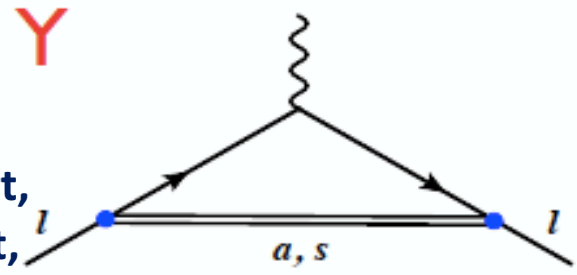
P. Paradisi, La Thuile 2021

**The case of AXION-LIKE PARTICLES (ALPs)**

# ALPs contributions to the muon g-2?

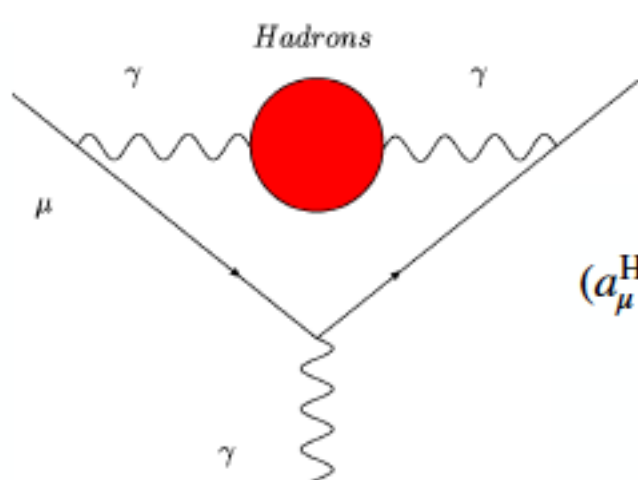


Marciano, AM, Paradisi,  
Passera '16; Bauer, Neubert,  
Thamm '17; Bauer, Neubert,  
Renner, Schnubel, Thamm '19;  
Cornella, Paradisi, Sumensari '19



- Both scalar and pseudoscalar ALPs can solve  $\Delta a_\mu$  for masses  $\sim [100\text{MeV}-1\text{GeV}]$  and couplings allowed by current experimental constraints.
- They can be tested at present low-energy  $e^+e^-$  experiments, via dedicated  $e^+e^- \rightarrow e^+e^- + \text{ALP}$  &  $e^+e^- \rightarrow \gamma + \text{ALP}$  searches.

# HVP: the major source of uncertainty in the muon g-2 SM computation



$$\text{Im} \left[ \text{wavy line} \cdot \text{red circle} \cdot \text{wavy line} \right] \sim \left| \text{wavy line} \rightarrow \text{red lines} \right|^2 \sim \sigma(e^+e^- \rightarrow \gamma^* \rightarrow \text{hadrons})$$

$$(a_\mu^{\text{HVP}})_{e^+e^-} = \frac{\alpha}{\pi^2} \int_{m_{\pi^0}^2}^{\infty} \frac{ds}{s} K(s) \text{Im} \Pi_{\text{had}}(s) = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \sigma_{\text{had}}(s)$$

dispersion relations

optical theorem

kernel function

$$K(s) \approx m_\mu^2/3s \quad \text{for} \quad \sqrt{s} \gg m_\mu$$

$$a_\mu^{\text{HLO}} = \frac{\alpha^2}{3\pi^2} \int_{4m_\mu^2}^{\infty} \frac{ds}{s} K(s) R(s)$$

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{had})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

$$a_\mu^{\text{HLO}} = 6895 (33) \times 10^{-11}$$

F. Jegerlehner, arXiv:1711.06089

$$= 6939 (40) \times 10^{-11}$$

Davier, Hoecker, Malaescu, Zhang, arXiv:1908.00921

$$= 6928 (24) \times 10^{-11}$$

Keshavarzi, Nomura, Teubner, arXiv:1911.00367





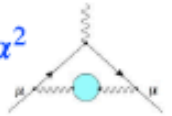
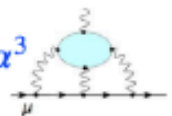
$$= 6931 (40) \times 10^{-11} (0.6\%)$$

WP20 value

WP20 = White Paper of the Muon g-2 Theory Initiative: arXiv:2006.04822

# Standard Model Contribution: Calculating the Anomaly

$$a_\mu = a_\mu(QED) + a_\mu(EW) + a_\mu(hadronic)$$

	contribution	error <sup>2</sup>
QED		
EW		
HVP		
HLbL		

Contribution	Order	Value	Uncertainty
QED	+... (5 loops)	$116\,584\,718.9(1) \times 10^{-11}$	0.001 ppm
EW	+... x	$153.6(1.0) \times 10^{-11}$	0.01 ppm
HVP	+... (NNLO)	$6845(40) \times 10^{-11}$	0.34 ppm [0.6%]
HLbL	+... (NLO)	$92(18) \times 10^{-11}$	0.15 ppm [20%]

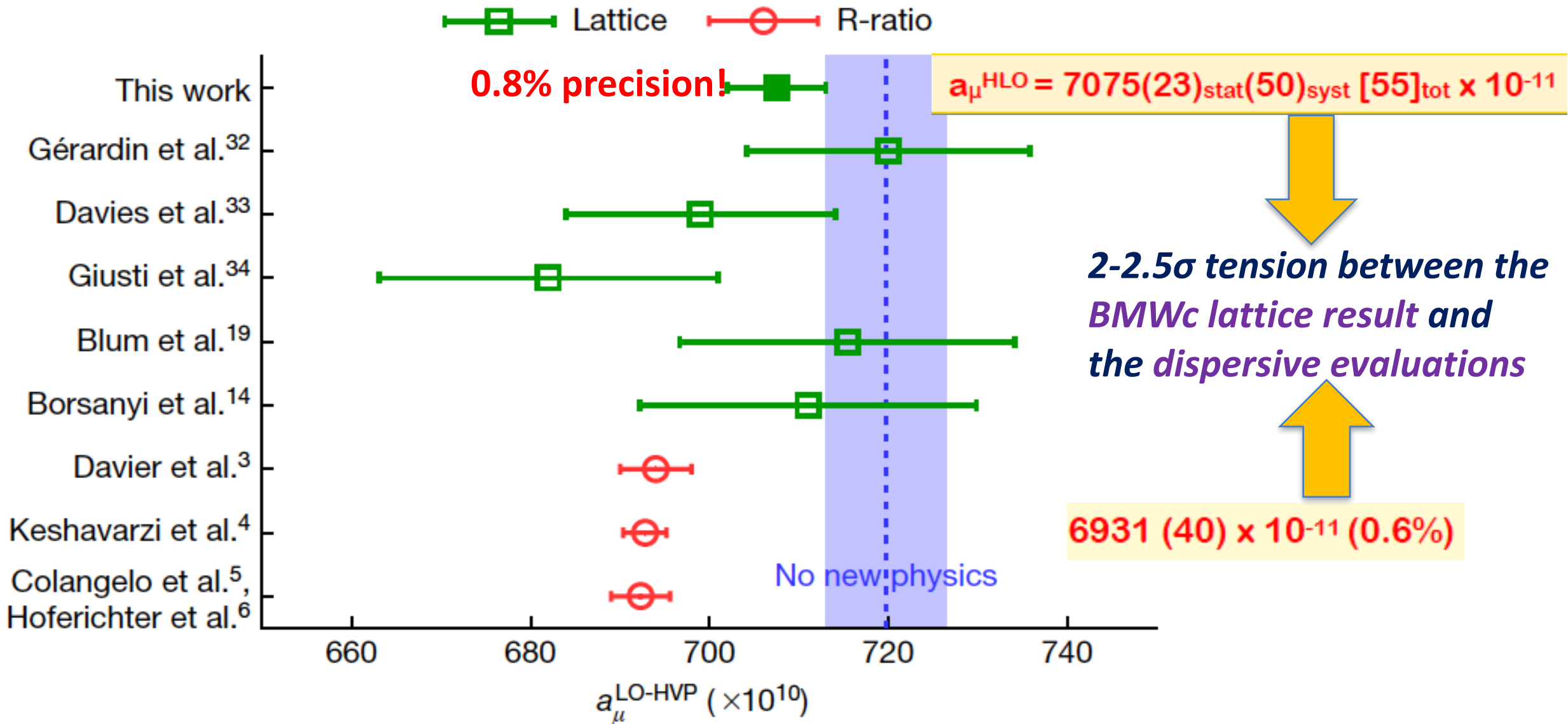
Well-known (QED, EW)

Non-perturbative (Data-driven & lattice QCD) (HVP, HLbL)

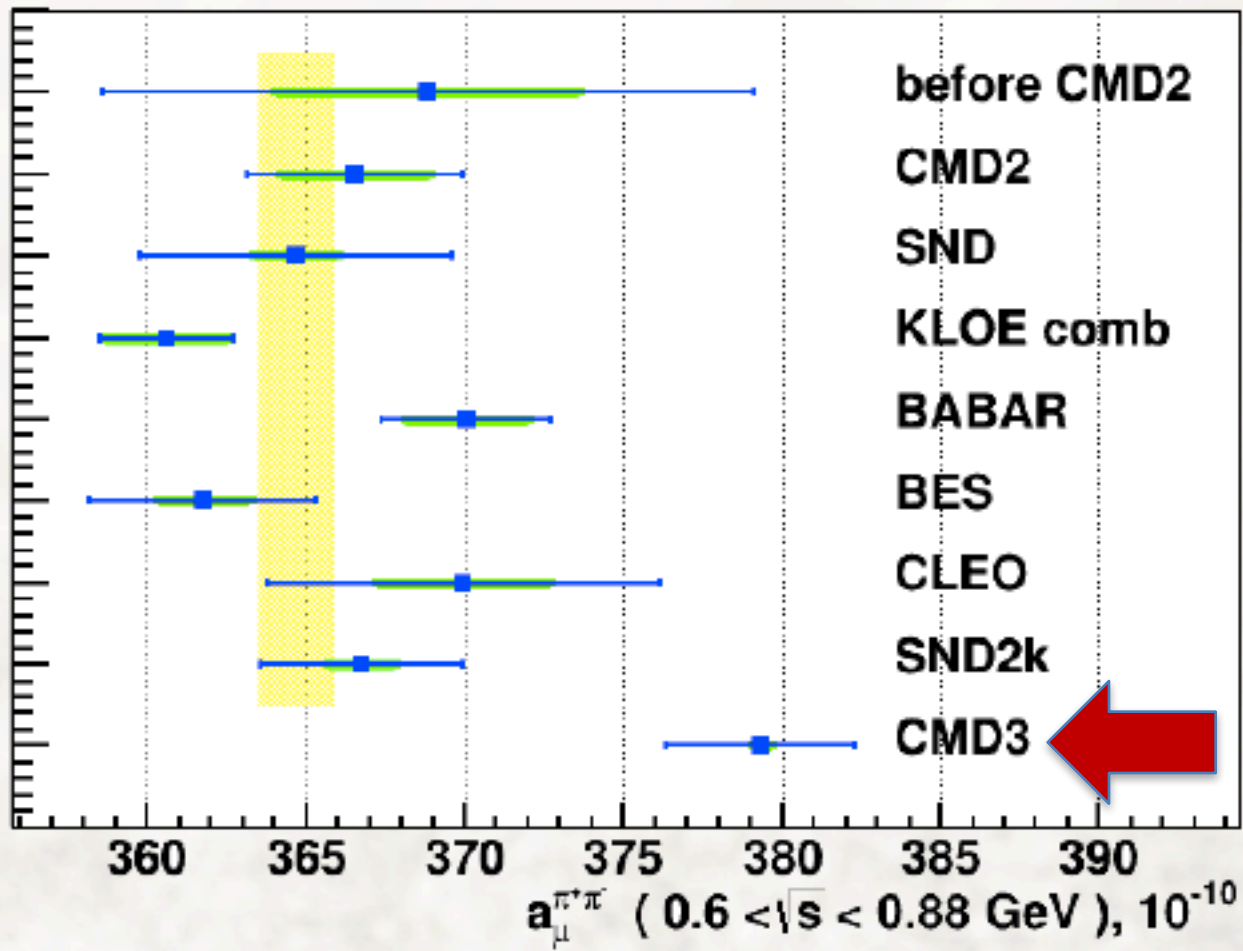
- QED and EW contributions are very well-known with small uncertainties
- Hadronic contribution error dominates the uncertainty budget
- HVP needs to be on the 0.5% precision to keep up with the experiment uncertainties
- HLbL precision demand is less than HVP, only 10% would be good enough
- Refining the SM calculations means refining the HVP calculation
- **Muon g-2 Theory Initiative** was formed to determine SM value of  $a_\mu$ . Produce a single consensus theoretical value which is comparable to the experimental value.

E. Barsal-Yucel, Lepton Photon 2023

BMWc20: S. Borsanyi et al. 2002.12347, published on Nature, April 7, 2021  
 first published lattice result with **sub-percent precision!**



$$a_{\mu}^{had,LO} = \frac{m_{\mu}^2}{12\pi^3} \int_{4m_{\pi}^2}^{\infty} \frac{\sigma_{e^+e^- \rightarrow \gamma^* \rightarrow hadrons}(s) K(s)}{s} ds$$



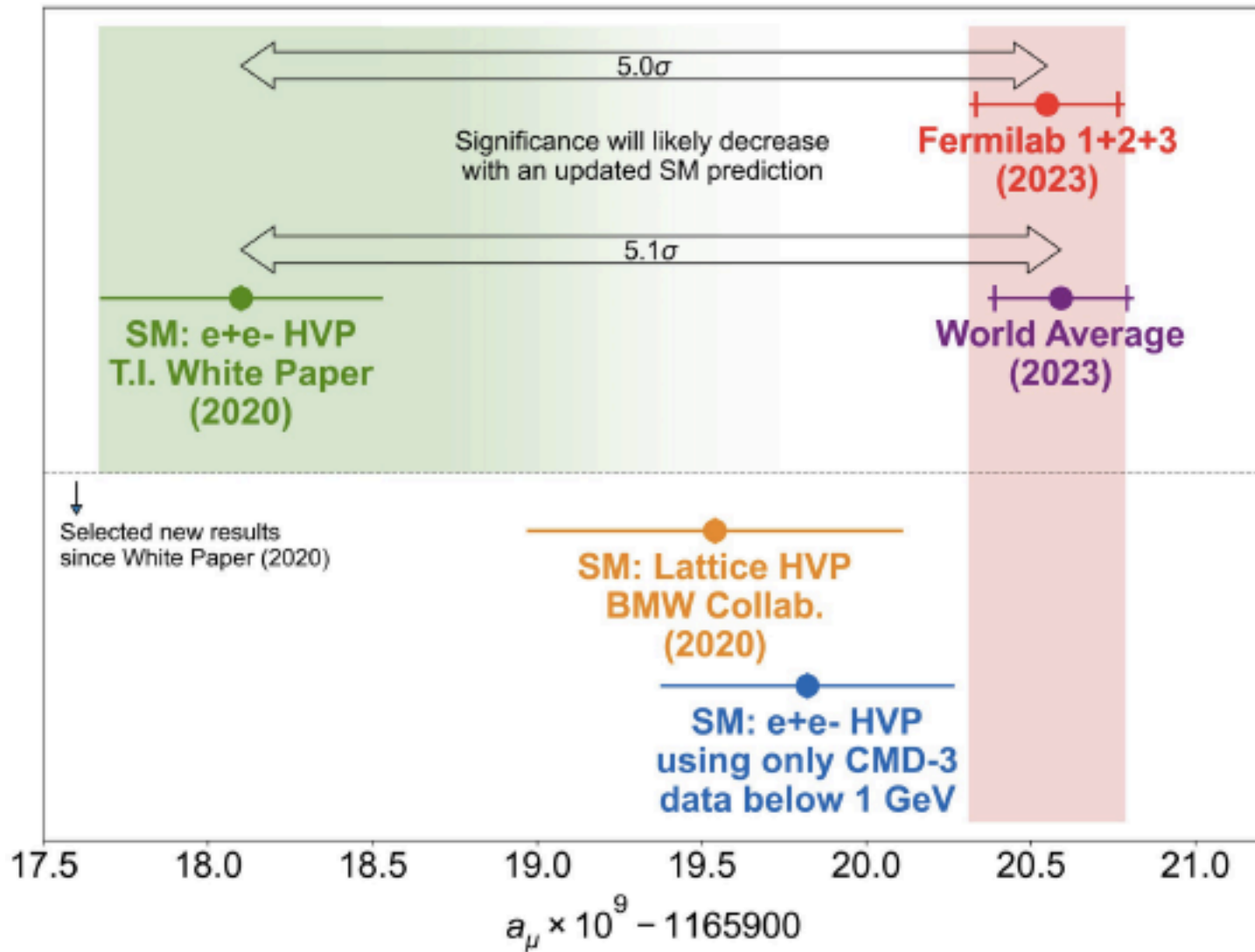
## New result on R(s) from **CMD3** (VEPP – 2000 Novosibirsk)

$0.6 < \sqrt{s} < 0.88 \text{ GeV}$

	$a_{\mu}^{\pi\pi, LO}, 10^{-10}$
before CMD2	$368.8 \pm 10.3$
CMD2	$366.5 \pm 3.4$
SND	$364.7 \pm 4.9$
KLOE	$360.6 \pm 2.1$
BABAR	$370.1 \pm 2.7$
BES	$361.8 \pm 3.6$
CLEO	$370.0 \pm 6.2$
SND2k	$366.7 \pm 3.2$
CMD3	$379.3 \pm 3.0$



F. Ignatov (CMD-3 Coll.), 6<sup>th</sup> Plenary Workshop TI, Bern, Sept. 4 2023

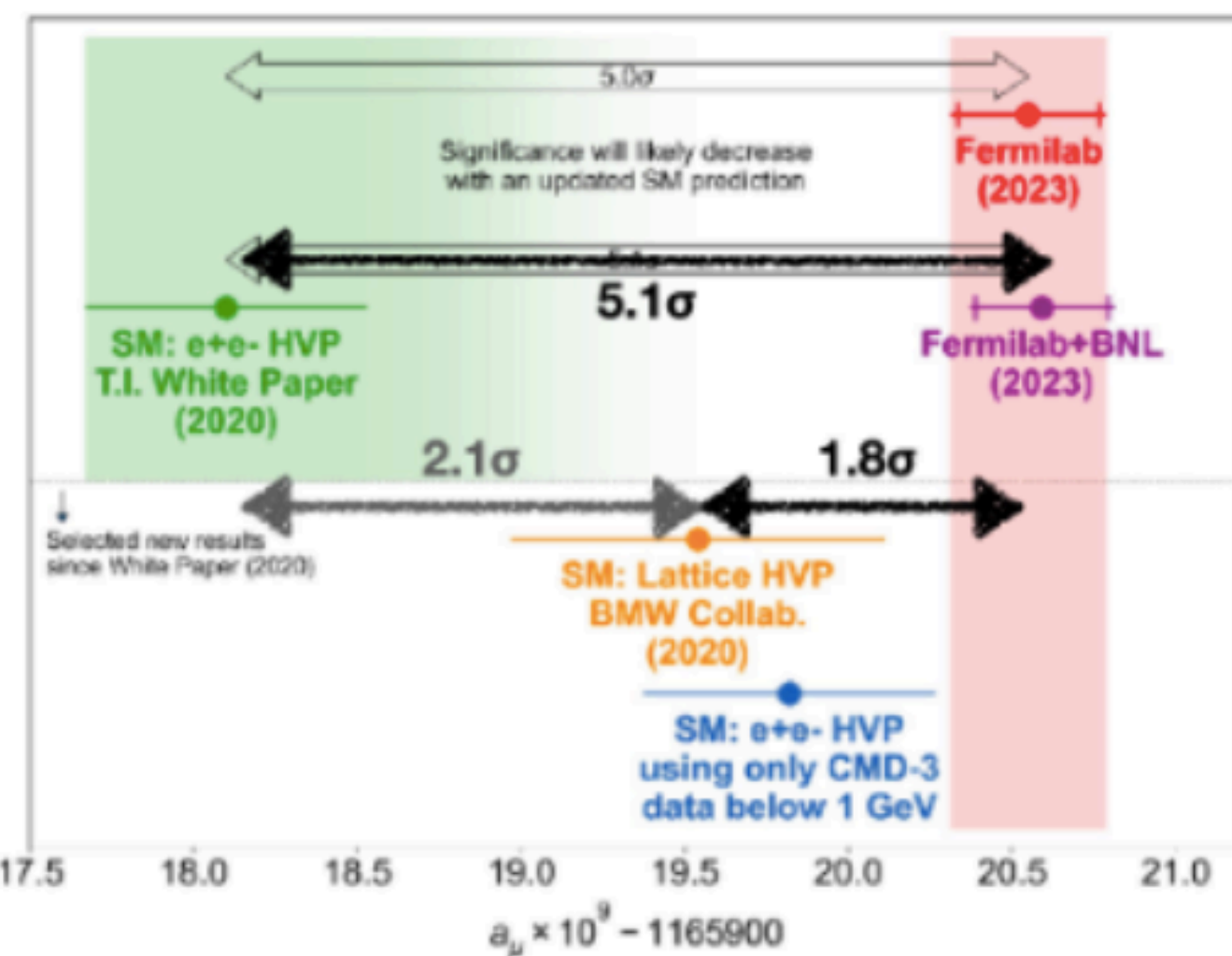


The CMD-3 is only one now over many other e+e- experiments (BaBar, KLOE, BES, CMD-2, SND, ...)

Unfortunately at the moment, we don't know the reasons of the disagreement between different experiments.

James Mott: <https://indico.fnal.gov/event/60738/>

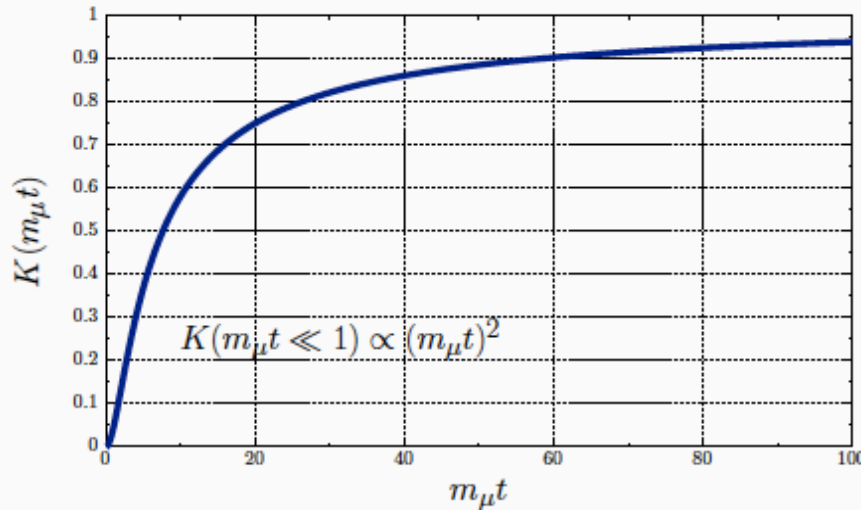
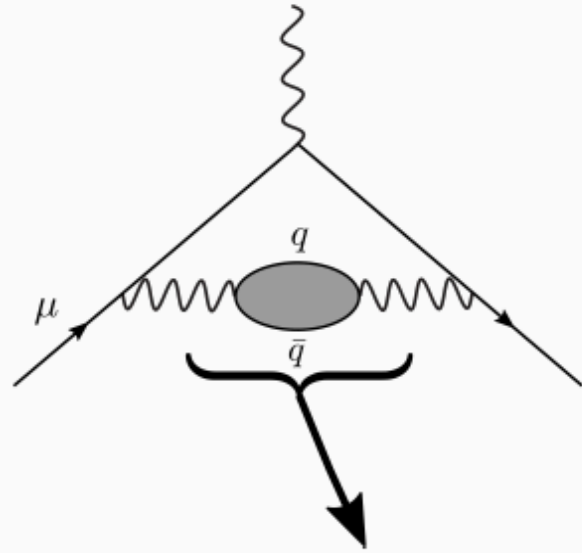
Alex Keshavarzi: <https://indico.fnal.gov/event/57249/contributions/271581/>



From  
 G. Venanzoni, EPS-HEP2023, Hamburg,

The **CMD-3** data in  $e^+e^- \rightarrow \pi\pi$  provides an R-ratio result compatible with the lattice one

# LO-HVP from Lattice QCD



$$\Pi_{\mu\nu}(Q) = \int d^4x e^{iQ \cdot x} \langle J_\mu(x) J_\nu(0) \rangle = (\delta_{\mu\nu} Q^2 - Q_\mu Q_\nu) \Pi(Q^2)$$

$$a_\mu^{\text{LO-HVP}} = 4\alpha_{em}^2 \int_0^\infty dQ^2 \frac{1}{m_\mu^2} f\left(\frac{Q^2}{m_\mu^2}\right) \cdot (\Pi(Q^2) - \Pi(0)).$$

**Time-Momentum representation (Bernecker & Meyer, 2011)**

$$a_\mu^{\text{LO-HVP}} = 2\alpha_{em}^2 \int_0^\infty dt t^2 K(m_\mu t) V(t), \quad V(t) \equiv \frac{1}{3} \sum_{i=1,2,3} \int d\vec{x} \langle J_i(\vec{x}, t) J_i(0) \rangle_i$$

Colangelo, El-Khadra, Hoferichter, Keshavarzi, Lehner, Stoffer, Teubner, arXiv:2205.12963v2 (2022)

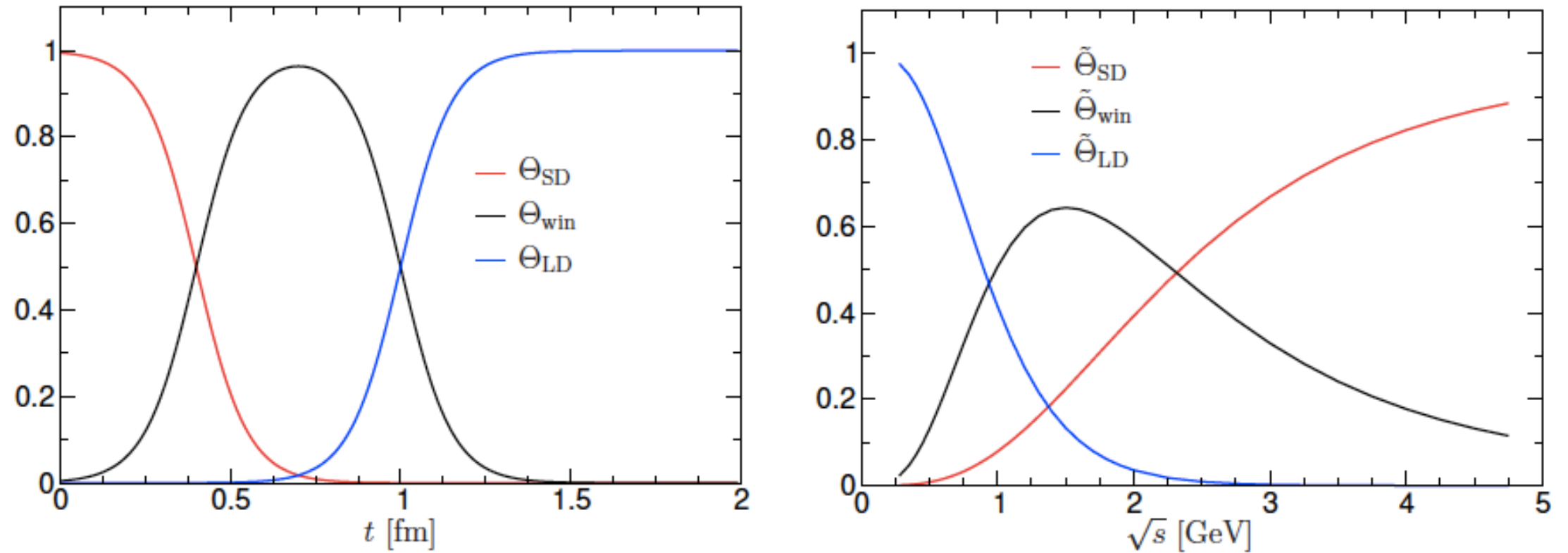
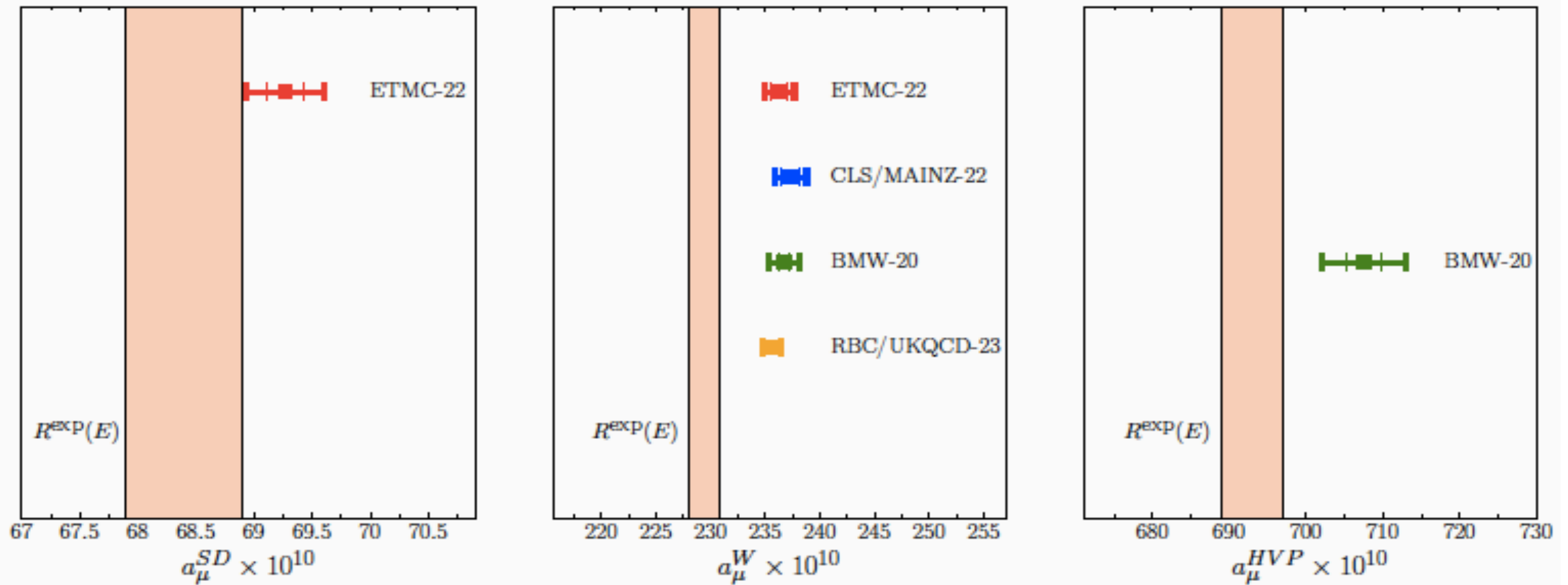


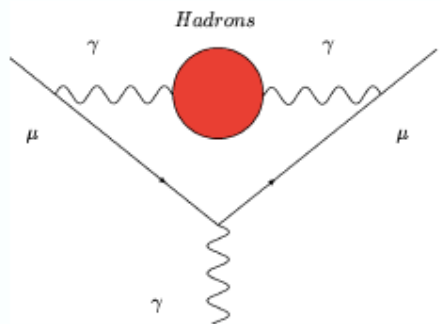
Figure 1: Short-distance, intermediate, and long-distance weight functions in Euclidean time (left), and their correspondence in center-of-mass energy (right).

The experimental ( $R^{\text{exp}}(E)$ -based) and the SM lattice QCD determination of the intermediate window are in significant tension [without CMD-3].



- Tension in  $a_\mu^W$  is larger than  $4\sigma$  (depending on how lattice results are combined).
- Substantial agreement for  $a_\mu^{SD}$  suggests that the tension is localized at intermediate/low energies  $E$ .

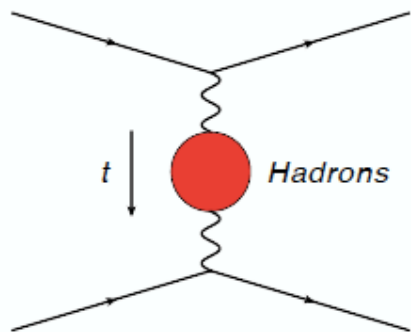
- At present, the leading hadronic contribution  $a_\mu^{\text{HLO}}$  is computed via the **timelike** formula:



$$a_\mu^{\text{HLO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{\infty} ds K(s) \sigma_{\text{had}}^0(s)$$

$$K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m_\mu^2)}$$

- Alternatively, exchanging the  $x$  and  $s$  integrations in  $a_\mu^{\text{HLO}}$



$$a_\mu^{\text{HLO}} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{\text{had}}[t(x)]$$

$$t(x) = \frac{x^2 m_\mu^2}{x-1} < 0$$

Lautrup, Peterman, de Rafael, 1972

$\Delta\alpha_{\text{had}}(t)$  is the hadronic contribution to the running of  $\alpha$  in the **spacelike region:  $a_\mu^{\text{HLO}}$  can be extracted from scattering data!**

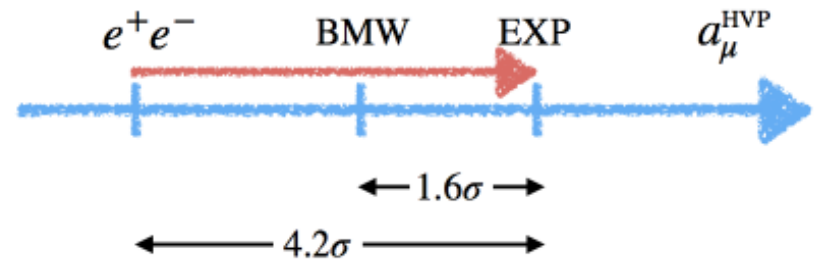


# New Physics to solve the **new** muon $g-2$ puzzle ?

Not including CMD3  $\rightarrow (a_\mu^{\text{HVP}})_{\text{EXP}} = a_\mu^{\text{EXP}} - a_\mu^{\text{SM, rest}}$

$$(a_\mu^{\text{HVP}})_{e^+e^-}^{\text{WP20}} = 6931(40) \times 10^{-11}$$

$$(a_\mu^{\text{HVP}})_{\text{BMW}} = 7075(55) \times 10^{-11}$$



NP in  $\sigma_{\text{had}}(e^+e^- \rightarrow \text{hadrons})$  such that

1.  $(a_\mu^{\text{HVP}})_{e^+e^-}^{\text{WP20}} \approx (a_\mu^{\text{HVP}})_{\text{EXP}}$
2. the approximate agreement between BMW and EXP is not spoiled
3. w/o a direct contribution  $a_\mu^{\text{NP}}$  (i.e. NP not in muons)

# Can $\Delta a_\mu$ be due to a missing contribution in $\sigma_{\text{had}}$ ?

[Marciano, Passera, Sirlin 2008 & 2010;  
 Keshavarzi, Marciano, Passera, Sirlin 2020.  
 See also Crivellin, Hoferichter, Manzari, Montull 2020;  
 Malaescu, Schott 2020;  
 Colangelo, Hoferichter, Stoffer 2020]

→ a upward shift of  $\sigma_{\text{had}}$  induces an increase of  $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$

$$\alpha(M_Z) = \frac{\alpha}{1 - \Delta\alpha(M_Z) - \Delta\alpha_{\text{had}}^{(5)}(M_Z) - \Delta\alpha_{\text{top}}(M_Z)}$$

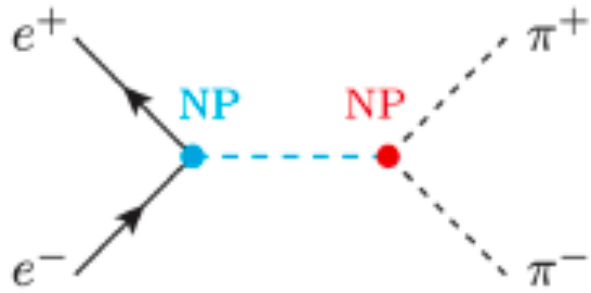
$$a_\mu^{\text{HLO}} \simeq \frac{m_\mu^2}{12\pi^3} \int_{4m_\pi^2}^{\infty} ds \frac{\sigma(s)}{s}, \quad \Delta\alpha_{\text{had}}^{(5)} = \frac{M_Z^2}{4\pi\alpha^2} \int_{4m_\pi^2}^{\infty} ds \frac{\sigma(s)}{M_Z^2 - s}$$

$$\text{Im} \left[ \text{wavy line} \bullet \text{wavy line} \right] \sim \left| \text{wavy line} \rightarrow \text{hadrons} \right|^2 \sim \sigma(e^+e^- \rightarrow \gamma^* \rightarrow \text{hadrons})$$

Shifts  $\Delta\sigma(s)$  to fix  $\Delta a_\mu$  are possible,  
 but conflict with the EW fit if they occur above  $\sim 1$  GeV

**Keshavarzi, Marciano, Passera,  
 Sirlin, PRD 2020 (updated 2021)**

Alternatively, one could invoke **NP intervening in Bhabha scattering**, see Darmé, Grilli di Cortona and Nardi, arXiv 2112.09139



**NP coupled both to hadrons and electrons**

but **not** directly to the **muons**

$$\text{Im} \left[ \text{wavy line} \cdot \text{red circle} \cdot \text{wavy line} \right] \sim \left| \text{wavy line} \rightarrow \text{hadrons} \right|^2 \sim \sigma(e^+e^- \rightarrow \gamma^* \rightarrow \text{hadrons})$$

$$(a_\mu^{\text{HVP}})_{e^+e^-} = \frac{\alpha}{\pi^2} \int_{m_{\pi^0}^2}^{\infty} \frac{ds}{s} K(s) \text{Im} \Pi_{\text{had}}(s) = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \sigma_{\text{had}}(s) \quad \sigma_{\text{had}} = \sigma_{\text{had}}^{\text{SM}} + \Delta\sigma_{\text{had}}^{\text{NP}}$$

**SUBTRACTION** since NP does **NOT** contribute to the HVP at the LO, but it **DOES** contribute to the cross-section at the LO

$$\sigma_{\text{had}} - \Delta\sigma_{\text{had}}^{\text{NP}}$$

a **POSITIVE** SHIFT on

$(a_\mu^{\text{HVP}})_{e^+e^-}$  requires  $\Delta\sigma_{\text{had}}^{\text{NP}} < 0$  (negative interference)

The unique scenario to obtain such a **SIZEABLE NEGATIVE interference**

- **SIZEABLE** → **TREE-LEVEL** contribution to modify  $\sigma_{\text{had}}$  at  $\sqrt{s} < 1 \text{ GeV}$  (hence, **sub-GeV mediator** coupling to the hadronic and electron currents at tree-level)
- **NEGATIVE INTERF.** → NP particle couples via a **VECTOR** current to the u, d quarks (given the dominance of the  $\pi^+\pi^-$  channel)

$$\mathcal{L}_{Z'} \supset (g_V^e \bar{e}\gamma^\mu e + g_V^q \bar{q}\gamma^\mu q) Z'_\mu \quad q = u, d \quad m_{Z'} \lesssim 1 \text{ GeV}$$

→ a light spin-1 mediator with vector couplings to first generation SM fermions

$$\frac{\sigma_{\pi\pi}^{\text{SM+NP}}}{\sigma_{\pi\pi}^{\text{SM}}} = \left| 1 + \frac{g_V^e (g_V^u - g_V^d)}{e^2} \frac{s}{s - m_{Z'}^2 + im_{Z'}\Gamma_{Z'}} \right|^2$$

However, **severe constraints on the  $Z'$  couplings** to electrons and to hadrons

- for  $m_{Z'} \lesssim 0.3 \text{ GeV}$  ( $Z' \rightarrow e^+e^-$  is the main decay mode)

$$e^+e^- \rightarrow \gamma Z' \text{ @ BaBar} \quad \longrightarrow \quad g_V^e \lesssim 2 \cdot 10^{-4}$$

- for  $m_{Z'} \gtrsim \text{MeV}$

$$\text{electron } g-2 \quad \longrightarrow \quad |g_V^e| \lesssim 10^{-2} (m_{Z'}/0.5 \text{ GeV})$$

$e^+e^- \rightarrow q\bar{q}$  has been measured with per-cent accuracy at LEP-II

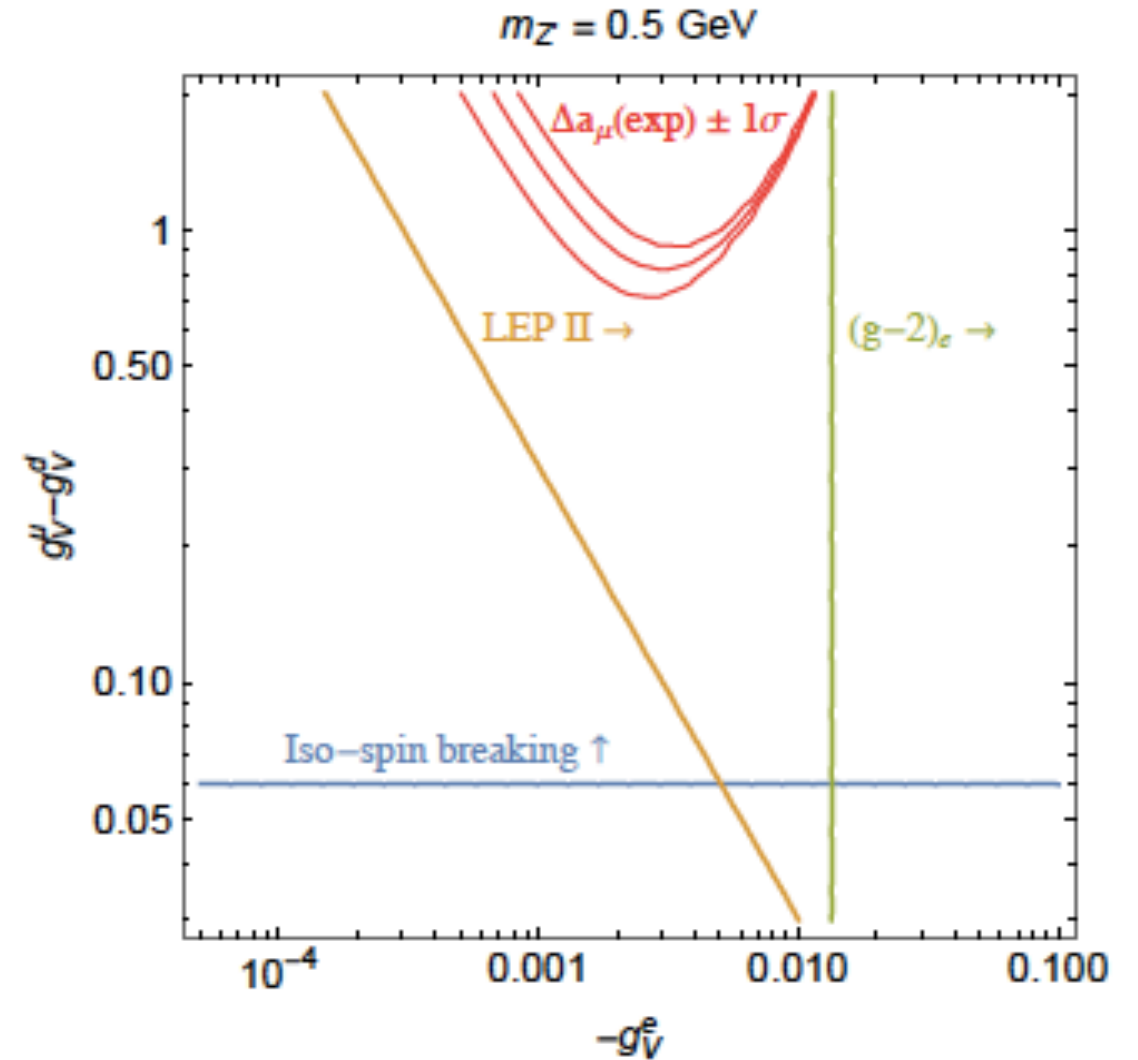
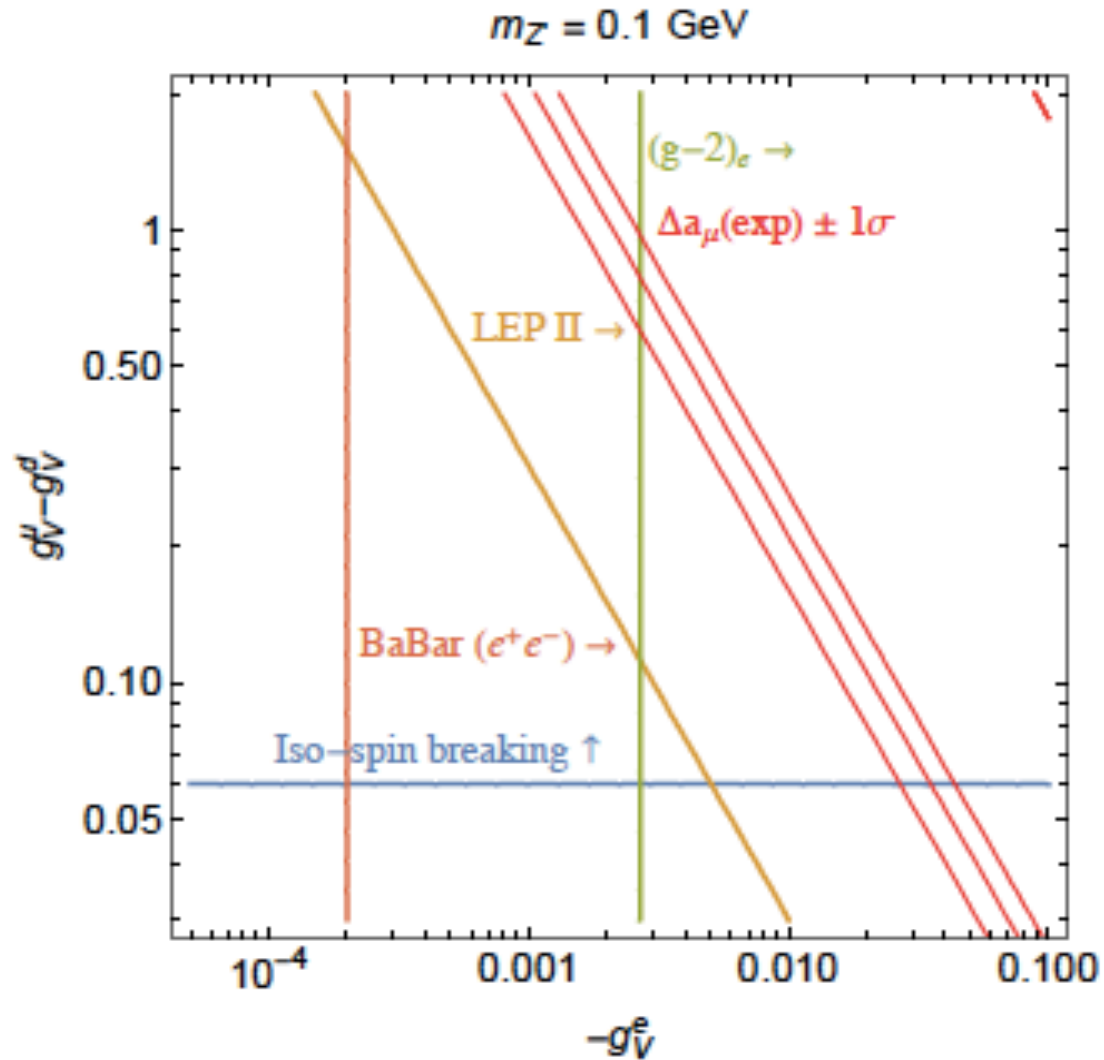
$$\frac{\sigma_{qq}^{\text{SM+NP}}}{\sigma_{qq}^{\text{SM}}} \approx 1 + 2 \frac{g_V^e g_V^q}{e^2 Q_q} \quad \longrightarrow \quad |g_V^e g_V^q| \lesssim 4.6 \cdot 10^{-4} |Q_q| \quad (\epsilon \lesssim 3.3 \cdot 10^{-3})$$

Iso-spin breaking observables

$$\longrightarrow \quad |g_V^u - g_V^d| \lesssim 0.06$$

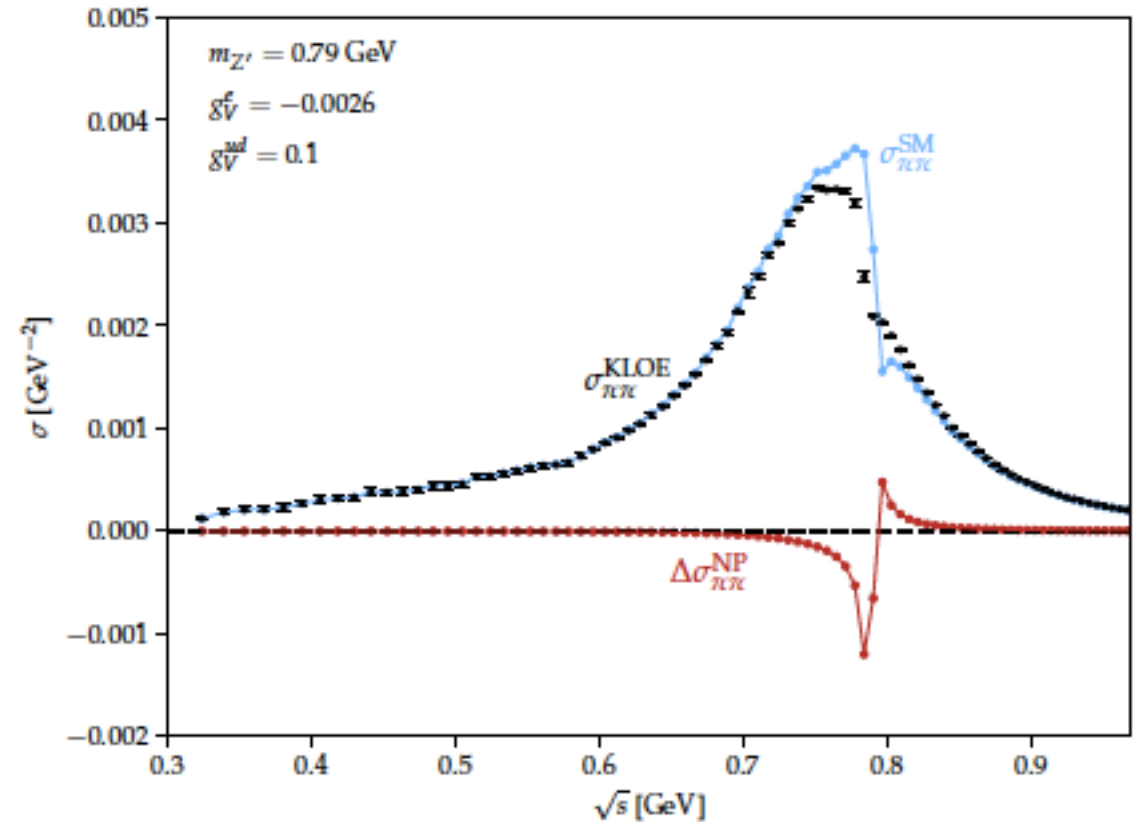
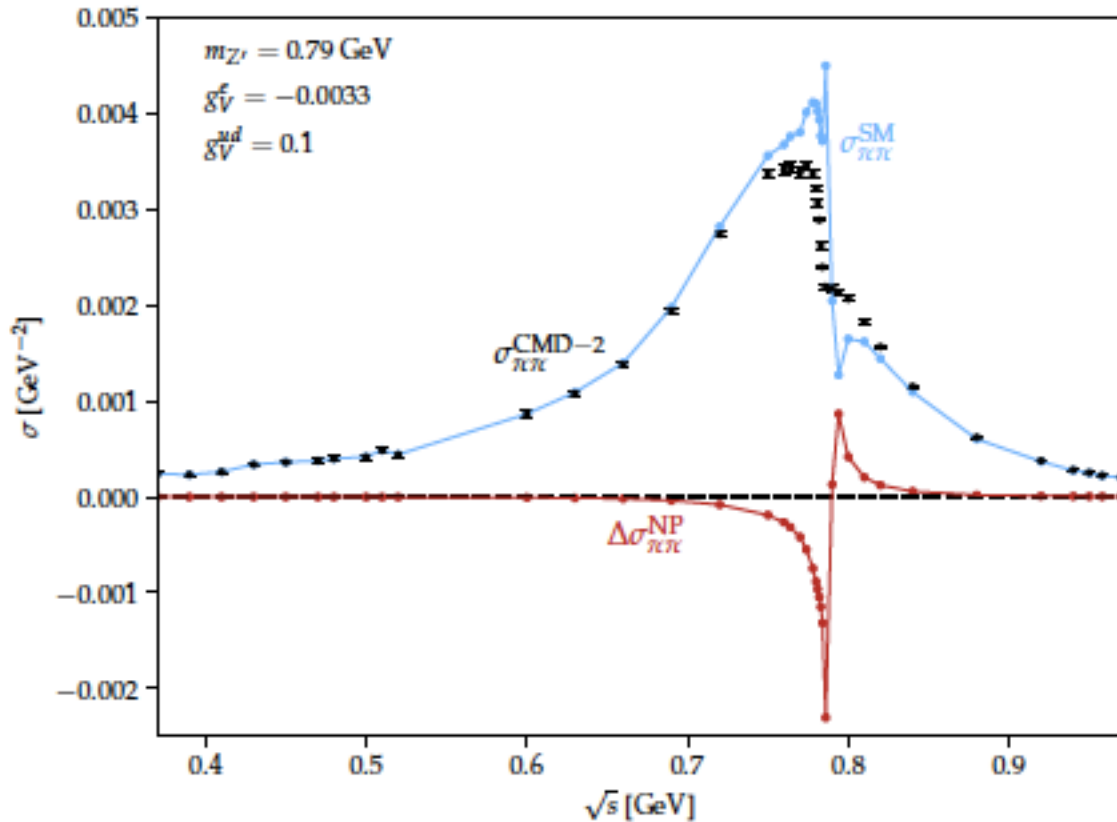
charged vs. neutral pion mass<sup>2</sup> difference  $\Delta m^2 = m_{\pi^+}^2 - m_{\pi^0}^2$  (rescaling the lattice QCD calculation of Frezzotti, Gagliardi, Lubicz, Martinelli, Sanfilippo and Simula 2112.01066)

At least **TWO independent bounds prevent** to get a sizeable contribution to  $\Delta a_\mu$  modifying  $\sigma_{\text{had}}$  via  $Z'$  exchange to **solve** the “**new**”  $\mu$   $g-2$  puzzle





However, Coyle and Wagner have recently claimed that **it is possible** to overcome the mentioned obstruction (in particular the isospin breaking constraint) by taking **a large  $g_V^u - g_V^d$  with a  $Z'$  mass near the  $\rho$  resonance mass of 770 MeV** – a lattice-QCD calculation is needed to provide a more precise evaluation of the isospin breaking replacing the massless photon with a massive  $Z'$  boson



# Sensitivity of other physical observables to

$$[\delta a_{\mu}^{\text{HVP}}]_{\text{NP}} = [a_{\mu}^{\text{HVP}}]_{\text{LQCD\&DR,CDM3}} - [a_{\mu}^{\text{HVP}}]_{\text{DR,WP202}}$$



If and to which extent the discrepancy between the leading HVP to the muon  $g-2$  computed, on one side, making use of the lattice QCD result by the BMW collaboration as well as the recent exp. results by the CMD-3 collaboration and, on the other side, using the low-energy  $e^+e^- \rightarrow$  hadrons data used by the Muon  $g-2$  Theory Initiative **can be tested via:**

- the Electron  $g-2$  ( $a_e$ )
- the HyperFine Splitting (**HFS**) in the muonium system
- the Tau  $g-2$  ( $a_{\tau}$ )

# Measurement of the Electron Magnetic Moment

X. Fan,<sup>1,2,\*</sup> T. G. Myers,<sup>2</sup> B. A. D. Sukra,<sup>2</sup> and G. Gabrielse<sup>2,†</sup>

<sup>1</sup>*Department of Physics, Harvard University, Cambridge, Massachusetts 02138, USA*

<sup>2</sup>*Center for Fundamental Physics, Department of Physics and Astronomy, Northwestern University, Evanston, Illinois 60208, USA*

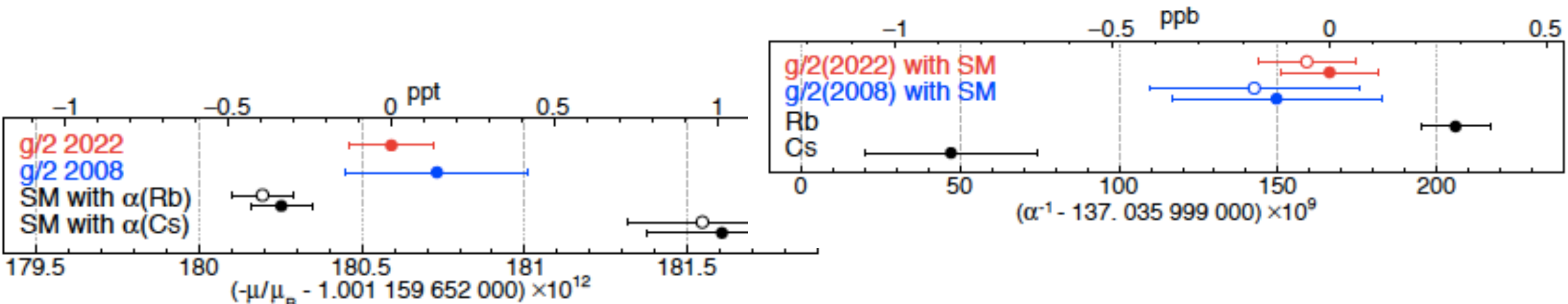
(Dated: December 8, 2022)

The electron magnetic moment,  $-\mu/\mu_B = g/2 = 1.001\,159\,652\,180\,59(13)$  [0.13 ppt], is determined 2.2 times more accurately than the value that stood for 14 years. The most precisely determined property of an elementary particle tests the most precise prediction of the Standard Model (SM) to 1 part in  $10^{12}$ . The test would improve an order of magnitude if the uncertainty from discrepant measurements of the fine structure constant  $\alpha$  is eliminated since the SM prediction is a function of  $\alpha$ . The new measurement and SM theory together predict  $\alpha^{-1} = 137.035\,999\,166(15)$  [0.11 ppb] with an uncertainty ten times smaller than the current disagreement between measured  $\alpha$  values.

$$a_e^{\text{EXP}} = 0.00115965218059(13)$$

$$\delta a_e^{\text{EXP}} = 1.3 \times 10^{-13}$$

In **2008** Gabrielse et al. had obtained  $\delta a_e^{\text{EXP}} = 2.8 \times 10^{-13}$



$$a_e^{\text{SM}} = a_e^{\text{QED}} + a_e^{\text{had+weak}} \quad \text{Jegerlehner EPJ Web Conf. 218 (2019)}$$

$$a_e^{\text{QED}} = \frac{\alpha}{2\pi} - 0.32847844400254(33) \left(\frac{\alpha}{\pi}\right)^2 + 1.181234016816(11) \left(\frac{\alpha}{\pi}\right)^3 - 1.91135(182) \left(\frac{\alpha}{\pi}\right)^4 + 7.791(580) \left(\frac{\alpha}{\pi}\right)^5$$

Improved, complete QED 5-loop contribution, expected soon, Laporta work in progress

Kinoshita et al. PRD 2017; Laporta PLB 2017

$$a_e^{\text{had+weak}} = 172.3(12) \times 10^{-14}$$

$\alpha = 1/137.035999046(27)$  (from Cs) Parker et al. Science 2018  
 $\alpha = 1/137.035999206(11)$  (from Rb) Morel et al. Nature 2020

Kinoshita et al. PRD 2017 – however S. Volkov PRD 2019 disagrees by an amount of  $\sim 7 \times 10^{-14}$

$$\delta a_e^{\text{QED5}} = 0.1 \times 10^{-13}$$

$$\delta a_e^{\text{HAD}} = 0.1 \times 10^{-13}$$

$$\delta a_e^{\delta\alpha} = 0.9 \times 10^{-13}$$

Using the most precise determination of  $\alpha$  from Rb Morel et al. Nature 2020

$$a_e^{\text{SM}} = 11596521816.1(0.9) \times 10^{-13}$$

$$a_e^{\text{HVP, LO}} = 186.08 \pm 0.66 \times 10^{-14}$$

Keshavarzi, Nomura and Teubner PRD 2020

$$a_e^{\text{EXP}} = 0.00115965218059(13)$$

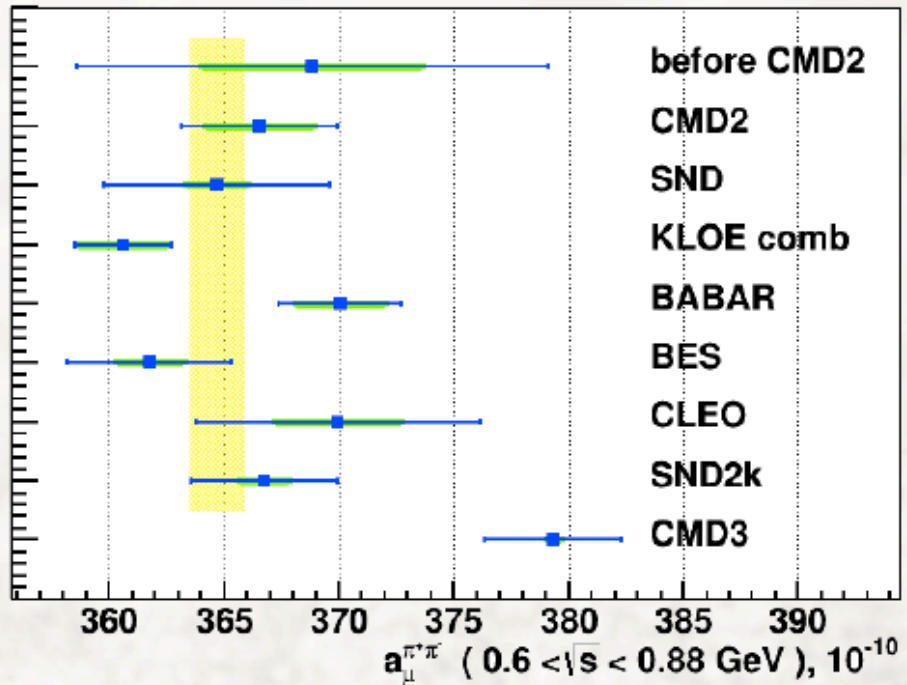


$$\delta a_e^{\text{EXP}} = 1.3 \times 10^{-13}$$

$a_e^{\text{HVP, LO}} = 186.08 \pm 0.66 \times 10^{-14}$  based on low-energy  $e^+e^- \rightarrow \text{hadrons}$  data **WITHOUT** the CMD-3 result

Impact of taking the CMD-3 result for the low-energy  $e^+e^- \rightarrow \text{hadrons}$  on  $a_e^{\text{HVP, LO}}$

$$a_\mu^{\text{had, LO}} = \frac{m_\mu^2}{12\pi^3} \int_{4m_\pi^2}^{\infty} \frac{\sigma_{e^+e^- \rightarrow \gamma^* \rightarrow \text{hadrons}}(s) K(s)}{s} ds$$



$$\delta a_\mu^{\text{CMD-3}} \simeq 150 \times 10^{-11}$$

$$\delta a_e^{\text{CMD-3}} \quad ?$$

Naïve scaling valid, for instance, in MFV models

$$\delta a_e^{\text{CMD-3}} = \delta a_\mu^{\text{CMD-3}} \left( \frac{m_e}{m_\mu} \right)^2 \simeq 3.5 \times 10^{-14}$$

$$\delta a_e^{\text{EXP}} = 1.3 \times 10^{-13}$$

$$\delta a_e^{\delta\alpha} = 0.9 \times 10^{-13}$$

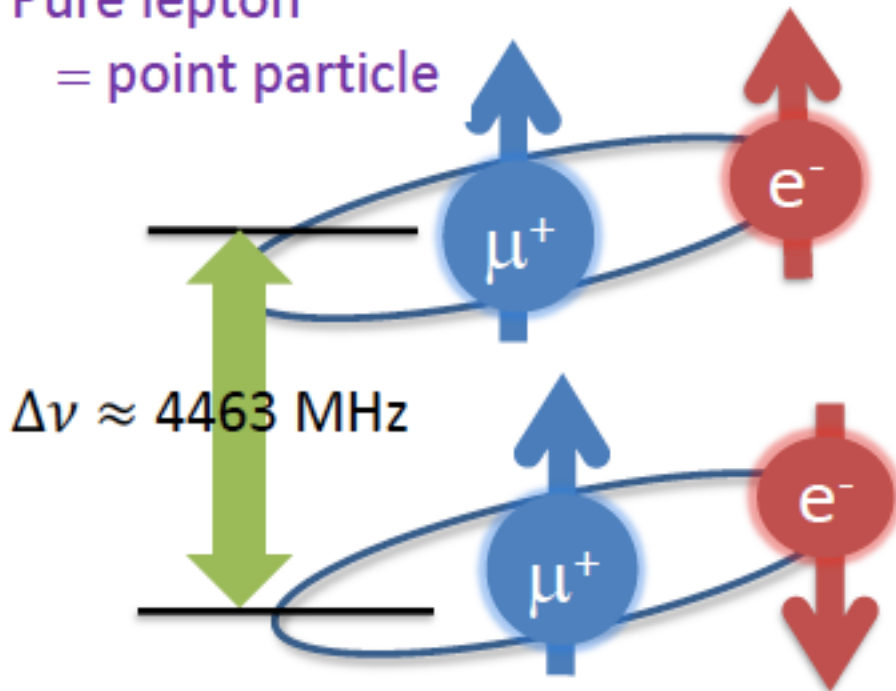
To check the  $a_\mu^{\text{HVP}}$  **BMW+CMD3**  $\longleftrightarrow$  **Muon g-2 TI** tension through the electron g-2 we need: a **THEORETICAL PREDICTION** & an **EXPERIMENTAL MEASUREMENT** of  $a_e$  at the level of  **$O(10^{-14})$**

preliminary Di Luzio, Keshavarzi, A.M., Paradisi and Passera, work in progress

# Independent muon g-2 determination from MUONIUM SPECTROSCOPY

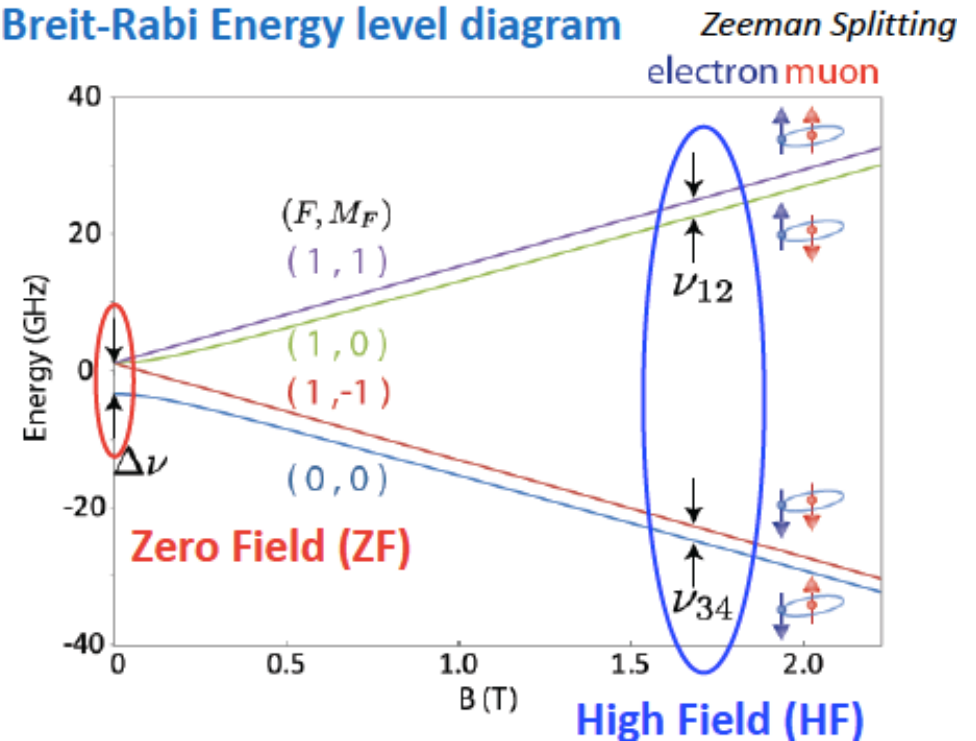
**Muonium: bound state of  $\mu^+$  and  $e^-$**

Pure lepton  
= point particle



$\Delta\nu$  : Muonium Hyperfine Structure

**Breit-Rabi Energy level diagram**



$$\nu_{12} + \nu_{34} = \Delta\nu$$

$$\nu_{12} - \nu_{34} \propto \mu_\mu / \mu_p$$



# Most Precise Test of Bound-State QED

## Experiment:

LAMPF Experiment (1999)

W. Liu *et al.*, Phys. Rev. Lett. **82** (1999) 711

$\nu_{\text{HFS}}(\text{exp})$	4463.302 765 (53) MHz	[12 ppb]
	$\mu_{\mu}/\mu_p = 3.18334524(37)$	[120ppb]
	$m_{\mu}/m_e = 206.768277(24)$	[120ppb]

## Theory:

M. I. Eides Phys. Lett. B **795** (2019) 113

$\nu_{\text{HFS}}(\text{theory})$	4463.302 868 (515) MHz	[120 ppb]
$\nu_{\text{HFS}}(\text{QED})$	4463.302 720 (511) (70) (2) MHz	
	$(m_{\mu}/m_e)(\text{QED})(\alpha)$	
$\nu_{\text{HFS}}(\text{weak})$	-65 Hz	
$\nu_{\text{HFS}}(\text{had. v.p.})$	232 (1) Hz	
$\nu_{\text{HFS}}(\text{had. h.o.})$	5 (2) Hz	

$$\nu(\text{Fermi}) = \frac{16}{3} \alpha^2 c R_{\infty} \frac{m_e}{m_{\mu}} \left( 1 + \frac{m_e}{m_{\mu}} \right)^{-3}$$

QED calculation: Effort for 10 Hz accuracy in progress (by Eides et al.) Also Laporta is computing the full 3-loop QED contribution

$$\Delta\nu_{\text{Mu}}^{\text{had, VP}} = (232.04 \pm 0.82) \text{ Hz}$$

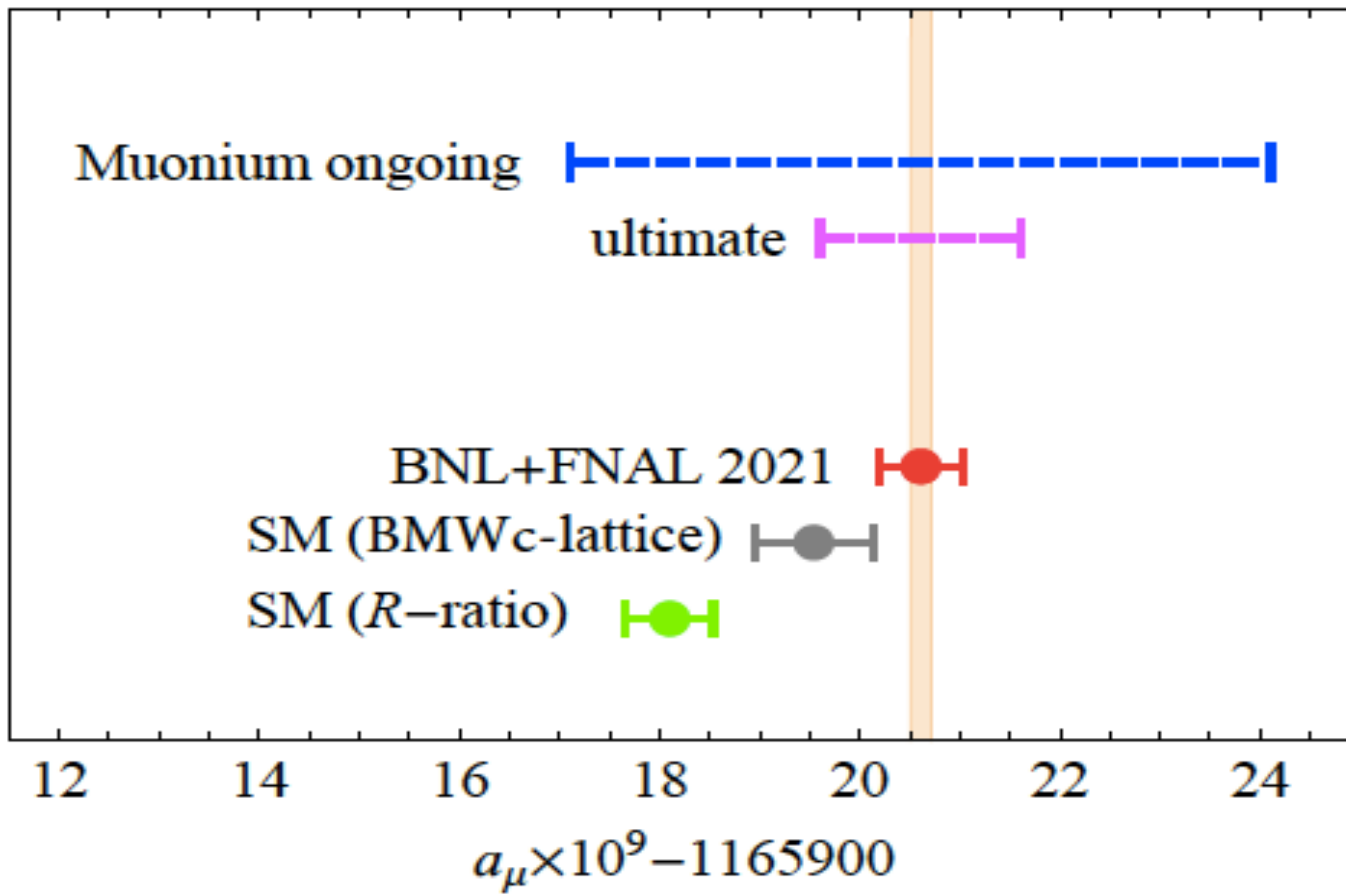
Dominated by the  $\pi^+\pi^-$  channel  $\Delta\nu_{\text{Mu}}^{\text{had, VP}} = (159.64 \pm 0.60) \text{ Hz}$

Ongoing indicates the milestones set by the **Mu-MASS (Muonium LAsEr SpectroScopy) exp. at PSI** and **MuSEUM (Muonium Spectroscopy Experiment Using Microwave) at J-**

P

$\nu_i$ (unit)	quantity	$u_r$			parameter (unit)	quantity	$u_r$		
		current	ongoing	ultimate			current	ongoing	ultimate
1S-2S (ppt)	QED	8.1	5.7	0.7	$m_e/m_\mu$ (ppb)	$\nu_{1S-2S}(\text{exp})$	825	0.84	0.34
	HVP	$\mathcal{O}(10^{-2})$				QED(1S-2S)	1.7	1.2	0.1
	$R_\infty$	1.9	0.65			$R_\infty$	0.40	0.13	
	$\alpha$	$\mathcal{O}(10^{-3})$				total	825	1.5	0.37
	exp	$3.99 \times 10^3$	4.1	1.6					
HFS (ppb)	QED	16	2.2	0.2	$a_\mu$ (ppm)	$\nu_{1S-2S}(\text{exp})$	708	0.73	0.29
	HVP	0.33	0.18			$\nu_{\text{HFS}}(\text{exp})$	10	1.9	0.77
	$\alpha$	0.30	0.16			QED(1S-2S)	1.4	1.0	0.07
	$R_\infty$	$\mathcal{O}(10^{-3})$				QED(HFS)	14	1.9	0.2
	exp	12	2.2	0.90		HVP(HFS)	0.29	0.16	
				$R_\infty$	0.35	0.13			
				$\alpha$	0.26	0.14			
				total	708	3.0	0.88		

Delaunay, Ohayon, Soreq PRL 2021



Delaunay, Ohayon, Soreq PRL 2021

To start testing the NP in the mentioned  $a_\mu^{\text{HVP}}$  discrepancy: [

$$\delta a_\mu^{\text{HVP}}]_{\text{NP}} = [a_\mu^{\text{HVP}}]_{\text{LQCD\&DR,CDM3}} - [a_\mu^{\text{HVP}}]_{\text{DR,WP202}} \quad (\sim 3\% \text{ shift})$$

→ probably **we need to reach for**  $\Delta v_\mu^{\text{exp}}$  and  $\Delta v_\mu^{\text{SM}}$  to precisions of O(few – 10) Hz

Di Luzio, Keshavarzi, A.M., Paradisi, Passera work in progress

# An extreme challenge: testing the NP of the **muon g-2 puzzle** through an accurate TH. and EXP. determination of the **TAU MAGNETIC MOMENT**

$$-0.007 < a_{\tau}^{\text{BSM}} < 0.005 \quad [2\sigma] \quad \text{from global analysis of LEP and SLD data in EFT}$$

Gonzalez-Sprinberg, Santamaria, Vidal NPB 2000

well above the Schwinger's 1-loop QED contribution!

What would be needed to be sensitive to the NP accounting for the muon g-2 tension: **rescaling with  $(m_{\tau}/m_{\mu})^2 \rightarrow a_{\tau}^{\text{NP}} \simeq 10^{-6}$**

ArXiv:2111.10378v2

PSI-PR-21-27, ZU-TH 56/21

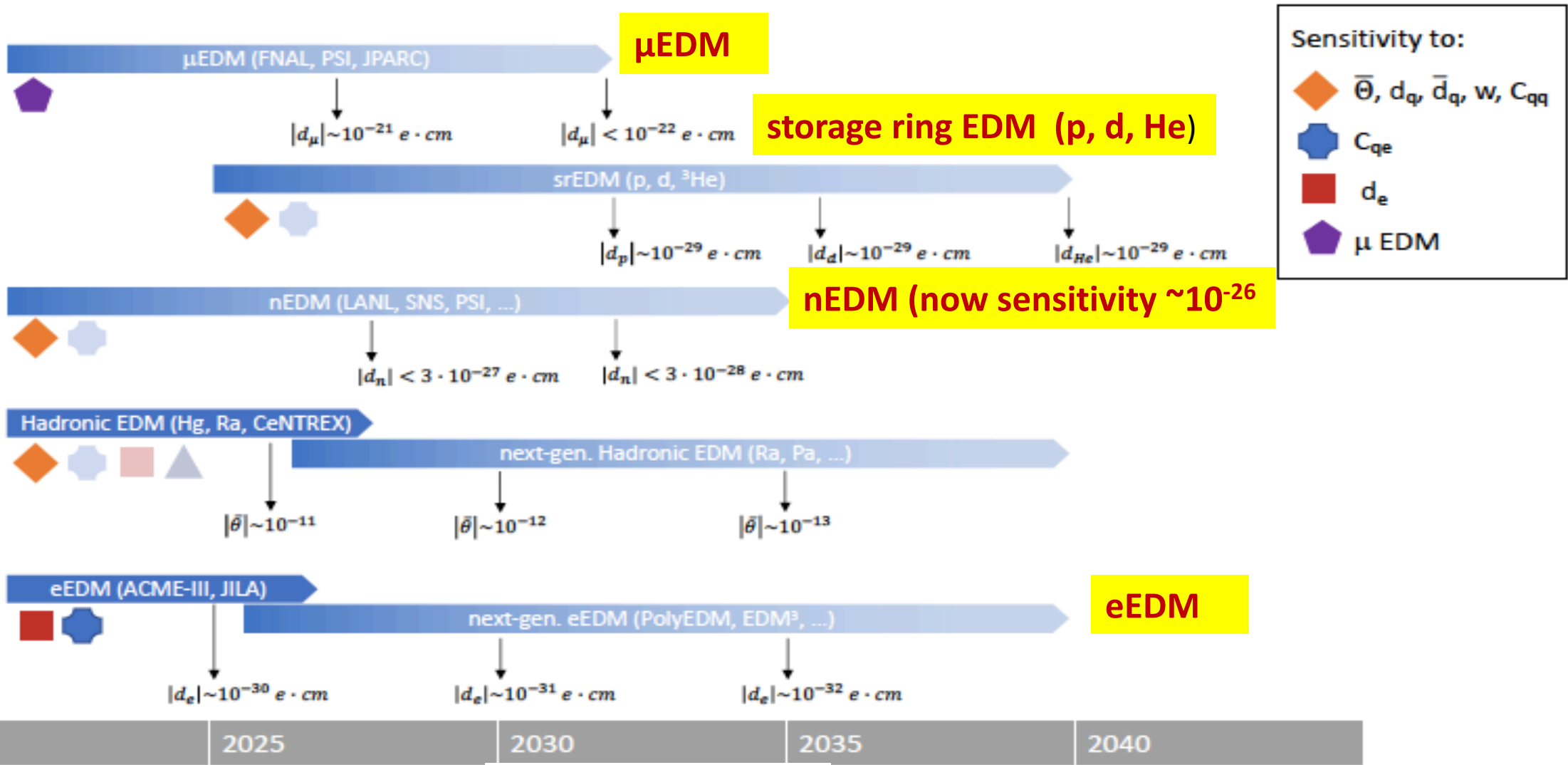
Towards testing the magnetic moment of the tau at one part per million

Andreas Crivellin,<sup>1,2</sup> Martin Hoferichter,<sup>3</sup> and J. Michael Roney<sup>4,5</sup>

**Belle II symmetry measurements with an important polarization upgrade of the SuperKEKB**

The impressive potentialities to explore the  
“**UNKNOWN**” **BSM physics**  
through the study of the **EDMs**

- **New science opportunities** in the (experimental and theoretical) current and near-future exploration of EDMs for various physical systems : **electron, muon, neutron, proton, atom, molecule**
- Coordinated program (with different scientific communities) of complementary EDM searches in **AMO** (Atomic Molecular Optical), **NUCLEAR** and **PARTICLE** physics
- An exceptionally sensitive way to explore the **NEW source(s) of CP VIOLATION** necessary to develop a cosmic asymmetry between matter and anti-matter starting with a symmetric early universe
- Feasible to achieve in a few years **relevant improvements** (from **one to even 3-4 orders of magnitude**) **on EDM sensitivities** – in particular AMO physics considers it realistic to achieve 1, 2-3, 4-6 orders of magnitude improvements in the few, 5-10 and 15-20 year time-scales, respectively



Fundamental Physics in Small Experiment T. Blum, P. Winter

**τEDM** BelleII: now  $\sim 10^{-18} e \cdot cm$  with beam polarization upgrade at SuperKEKB → **reach  $\sim 10^{-20} e \cdot cm$**

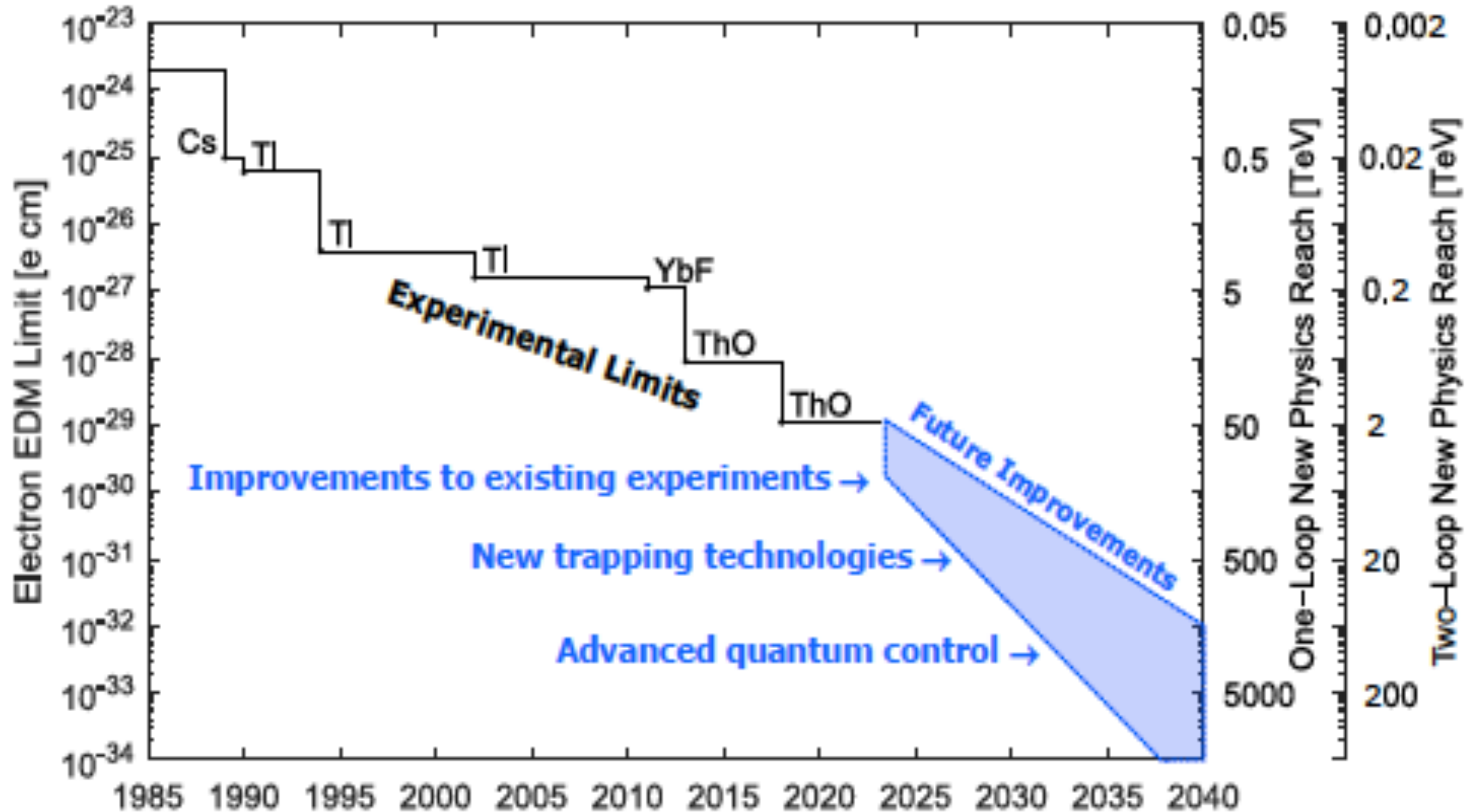
**COMMUNITY PLANNING EXERCISE: SNOWMASS 2021**

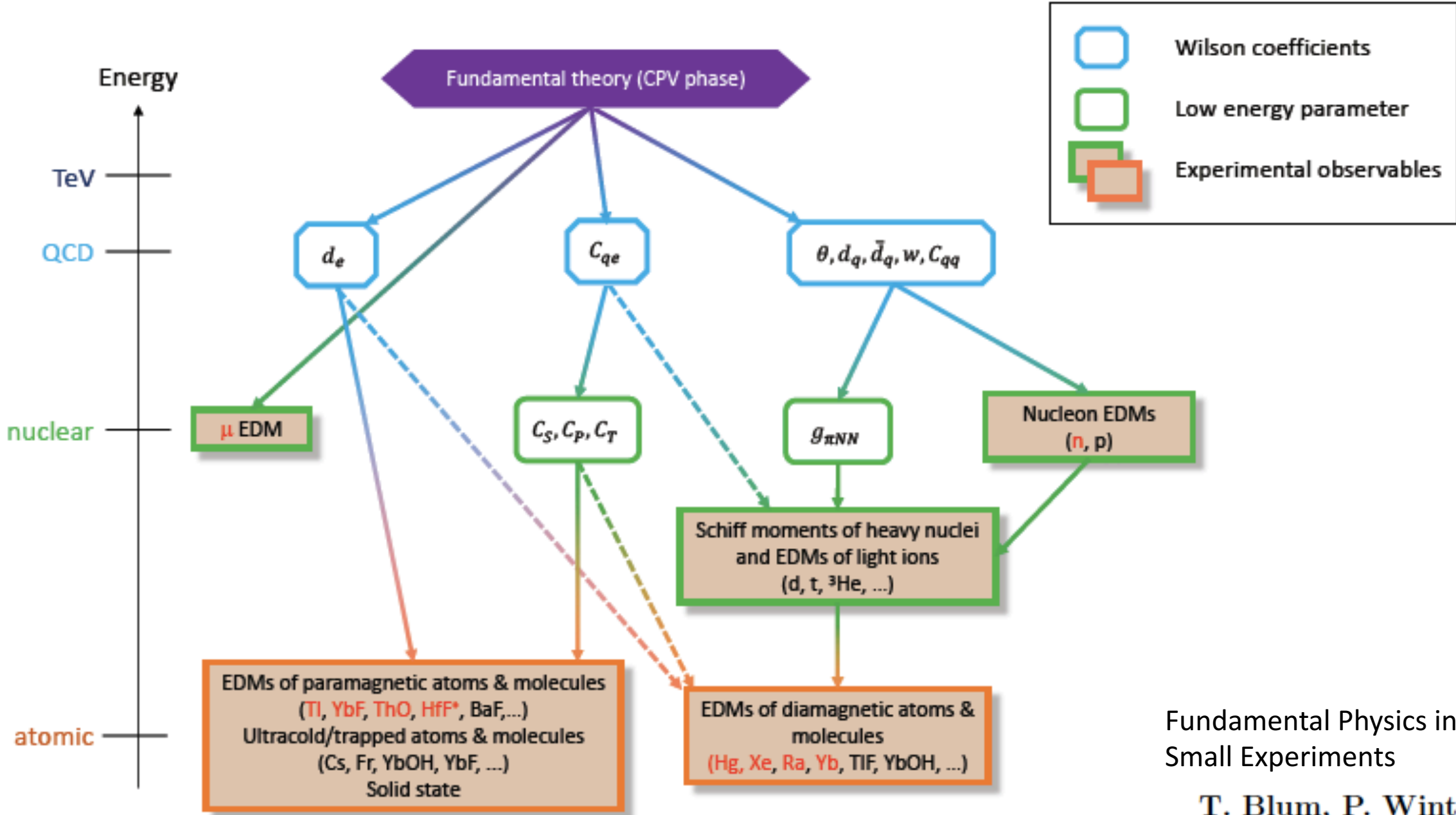
T. Bhattacharya, T.Y. Chen, V. Cirigliano, D. DeMille, A. Geraci, N.R. Hutzler, T.M. Ito, D. Kaplan, O. Kim, R. Lehnert, W.M. Morse, Y.K. Semertzidis



# great prospects for the (exp. & th.) progress in the **electron EDM physics**

COMMUNITY PLANNING EXERCISE: SNOWMASS 2021





Fundamental Physics in Small Experiments

T. Blum, P. Winter

# LFV, $(g - 2)_{\text{lept}}$ and $(\text{EDM})_{\text{lept}}$ correlations in Effective Theories

- $\text{BR}(\ell_i \rightarrow \ell_j \gamma)$  vs.  $(g - 2)_\mu$

Giudice, Paradisi and Passera JHEP 2012

$$\text{BR}(\mu \rightarrow e \gamma) \approx 3 \times 10^{-13} \left( \frac{\Delta a_\mu}{3 \times 10^{-9}} \right)^2 \left( \frac{\theta_{e\mu}}{10^{-5}} \right)^2$$

$$\text{BR}(\tau \rightarrow \mu \gamma) \approx 4 \times 10^{-8} \left( \frac{\Delta a_\mu}{3 \times 10^{-9}} \right)^2 \left( \frac{\theta_{\mu\tau}}{10^{-2}} \right)^2$$

- EDMs vs.  $(g - 2)_\mu$

$$d_e \approx \left( \frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 10^{-29} \left( \frac{\phi_e^{\text{CPV}}}{10^{-5}} \right) e \text{ cm},$$

$$d_\mu \approx \left( \frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 2 \times 10^{-22} \phi_\mu^{\text{CPV}} e \text{ cm},$$

- Main messages:

- ▶  $\Delta a_\mu \approx (3 \pm 1) \times 10^{-9}$  requires a nearly flavor and CP conserving NP
- ▶ Large effects in the muon EDM  $d_\mu \sim 10^{-22} e \text{ cm}$  are still allowed!

$$\frac{\Delta a_e}{\Delta a_\mu} = \frac{m_e^2}{m_\mu^2} \iff \Delta a_e = \left( \frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 0.7 \times 10^{-13}$$

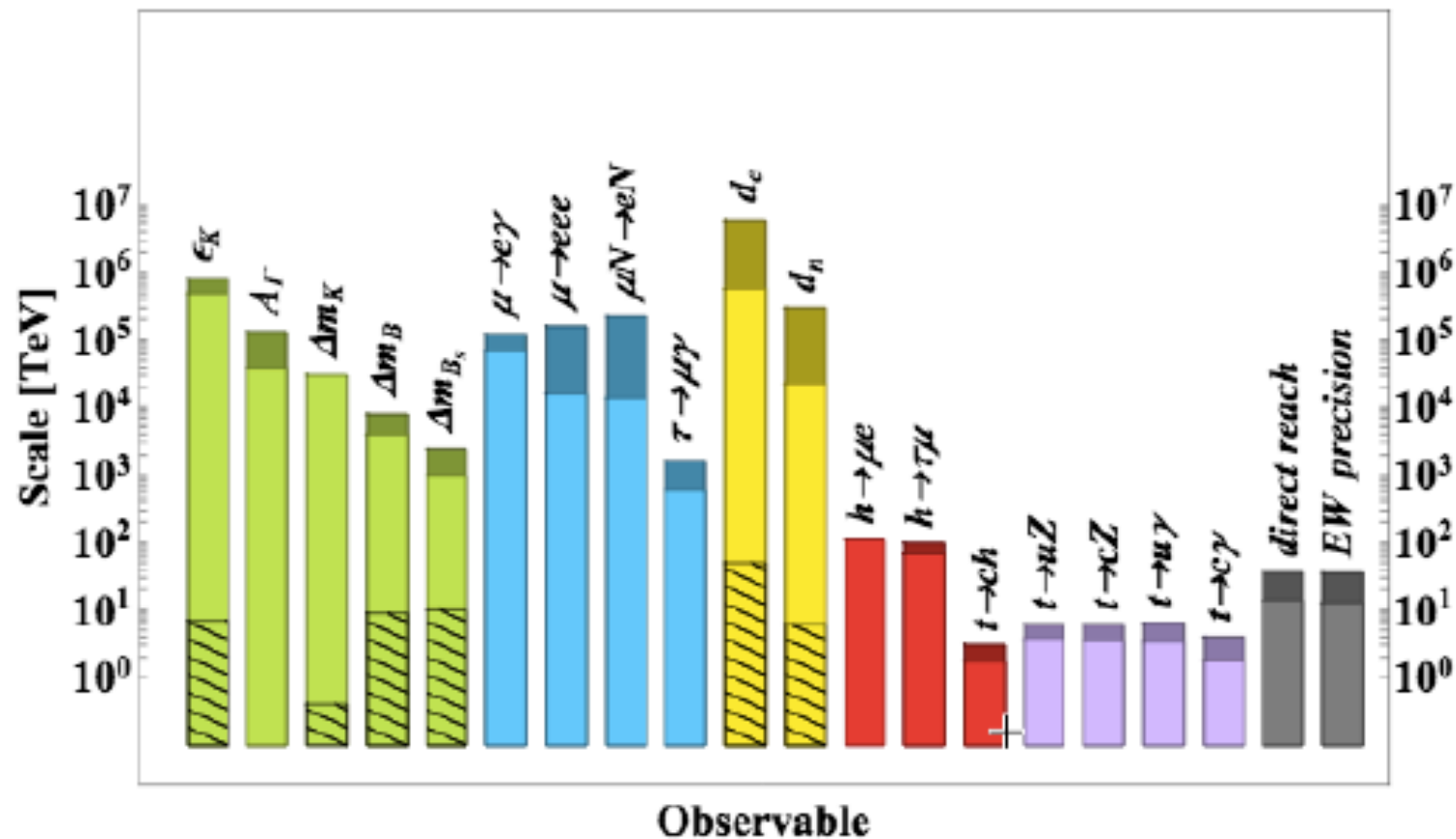
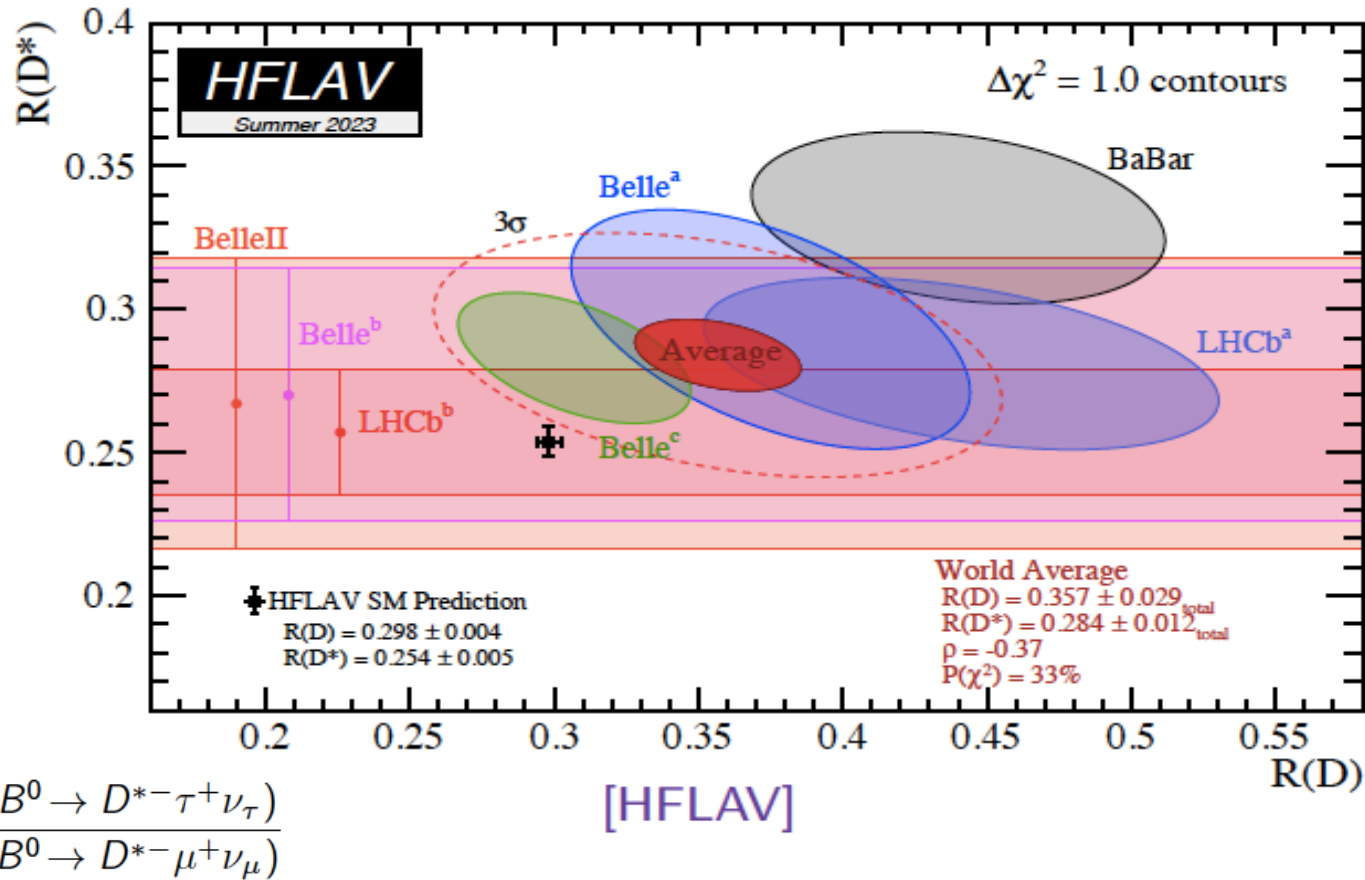


Fig. 5.1: Reach in new physics scale of present and future facilities, from generic dimension six operators. Colour coding of observables is: green for mesons, blue for leptons, yellow for EDMs, red for Higgs flavoured couplings and purple for the top quark. The grey columns illustrate the reach of direct flavour-blind searches and EW precision measurements. The operator coefficients are taken to be either  $\sim 1$  (plain coloured columns) or suppressed by MFV factors (hatch filled surfaces). Light (dark) colours correspond to present data (mid-term prospects, including HL-LHC, Belle II, MEG II, Mu3e, Mu2e, COMET, ACME, PIK and SNS).



$$\mathcal{R}(D^{*-}) = \frac{\mathcal{B}(B^0 \rightarrow D^{*-} \tau^+ \nu_\tau)}{\mathcal{B}(B^0 \rightarrow D^{*-} \mu^+ \nu_\mu)}$$

[HFLAV]

Tension of about  $3.3\sigma$  between average of measurements and SM predictions

A new anomaly?

A new player in the room!

$$R_{\nu\nu}^{K^{(*)}} = \mathcal{B}(B \rightarrow K^{(*)} \nu \bar{\nu}) / \mathcal{B}(B \rightarrow K^{(*)} \nu \bar{\nu})^{\text{SM}}$$

Belle II at EPS conference, 2023

$$R_{\nu\nu}^K < 3.6 \quad (90\% \text{ C.L.}),$$

$$R_{\nu\nu}^{K^*} < 2.7 \quad (90\% \text{ C.L.})$$

Glazov at EPS 2023

Privately produced comparison

Belle, 1702.03224

Belle II 2023

$$\mathcal{B}(B^\pm \rightarrow K^\pm \nu \bar{\nu}) = 2.40(67) \times 10^{-5}$$

2.9  $\sigma$  larger than SM prediction

Searching for explanation

Bause et al., 2309.00075

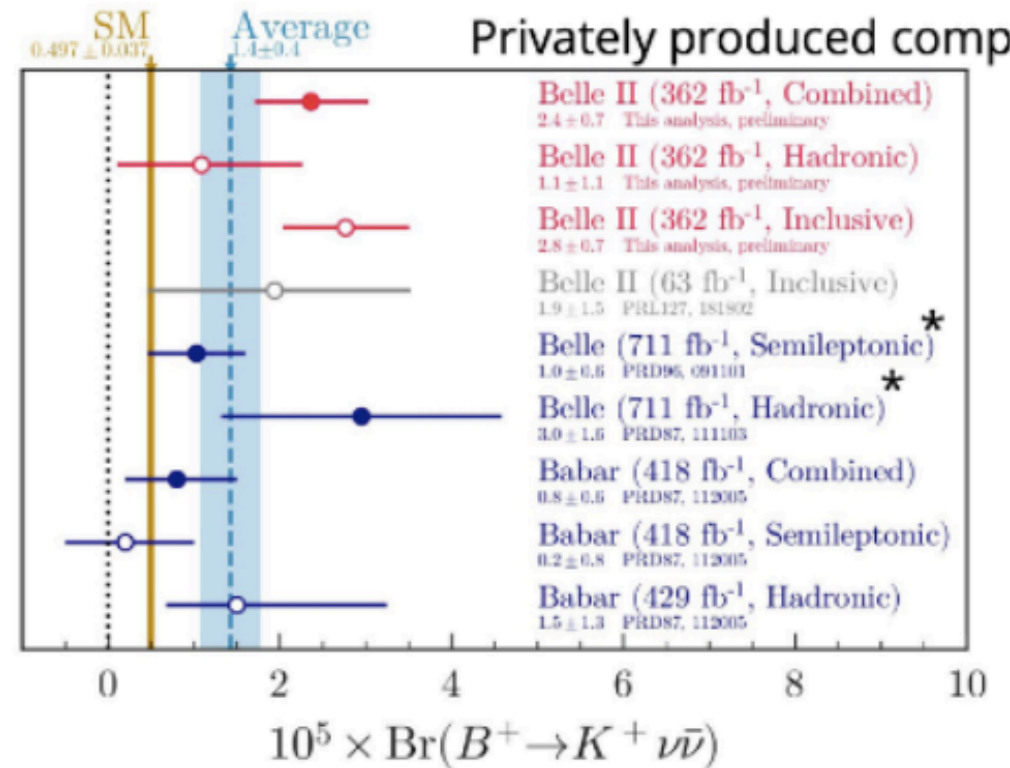
Allwicher et al., 2309.02246

Felkl et al., 2309.02940,

He et al., 2309.12741

...

$$R_{\nu\nu}^K = 5.4 \pm 1.5$$



S. Fajfer, Workshop on Implications of LHCb measurements and future prospects, CERN, 25-27 Oct.2023



# An **exciting** (**challenging** and **promising**) **era** for the **precision frontier physics**

- The experimental and theoretical precision physics community has entered an era of **unprecedented precision experiments**
- From Snowmass 2021: “**While relatively small in size and cost compared to their energy frontiers cousins, they are large in reach and discovery potential**”
- Very relevant (I’d say, necessary) to efficiently coordinate the **many experimental and theoretical efforts** in the area through a convinced **synergy** among the **various communities** operating in precision physics in (very) **different experimental, technological and theoretical environments**