Proton final-states and kinematic imbalance in neutrino cross sections with the MicroBooNE detector

Phys. Rev. Lett. 131, 101802 (2023), Phys. Rev. D 108, 053002 (2023), arXiv:2310.06082

Afroditi Papadopoulou <u>apapadopoulou@anl.gov</u> Liverpool, 9/1/2024



High-Precision Neutrino Measurements

Question: why matter dominated universe ?

(Potential) answer: Neutrino parameters lead to preferred matter production

Experimental test: Forthcoming DUNE experiment will measure neutrino (matter) and anti-neutrino (anti-matter) interaction rates

Challenge: unprecedented understanding of neutrino-nucleus interactions





Neutrinos 101

Flavor eigenstates

Mass eigenstates



Experimental Overview

Two-neutrino approximation

$$P(\nu_{\alpha} \rightarrow \nu_{\beta}) \sim \sin^2(2\theta) \sin^2\left(\frac{\Delta m_{ij}^2 L}{4E}\right)$$

L: distance (baseline) E: neutrino energy Δm²: neutrino mass splitting

$$U_{\rm PMNS} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{\rm CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{\rm CP}} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad \begin{array}{c} c_{\rm ij} = \cos\theta_{\rm ij} \\ s_{\rm ij} = \sin\theta_{\rm ij} \\ \theta = {\rm mixing \ angle} \end{array}$$

Experimental Overview

Optimizing baseline L for a given neutrino source of energy E Mass eigenstate splitting Δm^2 and mixing angles θ known to few %-level







Neutrino Oscillation Experiments



 $N_{ND}^{\alpha}(E_{rec})$: Number of events of flavor α

 E_{rec} = Reconstructed v energy

Near Detector



$$\begin{split} \underset{\text{Notice of the selection}}{\text{Neutrino flux}} & \underset{\text{selection efficiency}}{\text{Neutrino efficiency}} \\ N_{\text{ND}}^{\alpha}(E_{\text{rec}}) \sim \Phi_{\text{ND}}(E_{\nu}) \ \sigma(E_{\nu}) \ \epsilon_{\text{ND}}(E_{\nu}) \\ & \underset{\text{Neutrino cross}}{\text{Neutrino cross}} \end{split}$$

 $E_v = True v energy$ Also dependence on detector observables

Neutrino Oscillation Experiments



 $N_{FD}^{\ \beta}(E_{rec})$: Number of events of flavor β

$$E_{rec}$$
 = Reconstructed v energy

Far Detector





 $E_v =$ True v energy Also dependence on detector observables

Far Detector



- Smearing relating E_v to E_{rec}
- Neutrino signal topologies
- Neutrino backgrounds for BSM
- ...

Far Detector



$$N_{FD}^{\ \beta}(E_{rec}) \sim \Phi_{FD}(L, E) \sigma(E_{\nu}) \epsilon_{FD}(E_{\nu}) \mathbf{P}(\mathbf{v}_{\alpha} \rightarrow \mathbf{v}_{\beta})$$
 Oscillation
Need for high precision

v cross section measurements

Challenging ...

- Broad neutrino spectra
- Various complex interaction mechanisms

Any mismodeling in neutrino event generator simulation predictions can limit experimental sensitivity

Х

Quasi-elastic (QE)

Х

Meson Exchange

Current (MEC)



Resonance (RES)

Deep Inelastic

Scattering (DIS)

Future Experiments

50% CP Violation Sensitivity



- Mismodeling can impact required run time of forthcoming flagship experiments
- But ... head start with Short-Baseline Neutrino (SBN) Program (<u>MicroBooNE</u>, SBND, ICARUS)

DUNE CDR, <u>arXiv:1512.06148</u>



85 tonne Liquid Argon Time Projection Chamber (LArTPC) JINST 12, P02017 (2017)

LArTPC Operation Principle





MicroBooNE

- 3 wire planes
- 8192 gold coated wires
- 3 mm wire spacing
- 32 PMTs

MicroBooNE Data Events



- Excellent spatial resolution
- Low detection thresholds
- Precise calorimetric information
- Powerful particle identification

MicroBooNE Data Events



• Largest available neutrino-argon data set with ~500k recorded neutrino interactions

• 15 released and more than 30 active MicroBooNE cross section analyses

• Multiple topologies investigated

Already Public Results

CC inclusive

- 1D ν_µ CC inclusive @ BNB
 <u>Phys. Rev. Lett. 123, 131801 (2019)</u>
- 1D ν_µ CC E_ν @ BNB <u>Phys. Rev. Lett. 128, 151801 (2022)</u>
- 3D CC E_v @ BNB arXiv:2307.06413, submitted to PRL
- 1D v_e CC inclusive @ NuMI <u>Phys. Rev. D105, L051102 (2022)</u> <u>Phys. Rev. D104, 052002 (2021)</u>

Pion production

• ν_µ NCπ⁰ @ BNB <u>Phys. Rev. D 107, 012004 (2023)</u>

Rare channels

- η production @ BNB, submitted to PRL <u>arXiv:2305.16249</u>
- Λ production @ NuMI <u>Phys. Rev. Lett. 130, 231802 (2023)</u>

CC0π

- 1D ν_e CCNp0π @ BNB Phys. Rev. D 106, L051102 (2022)
- 1D & 2D ν_{μ} CC1p0 π Transverse Imbalance @ BNB <u>Phys. Rev. Lett. 131, 101802 (2023)</u>
 - Phys. Rev. D 108, 053002 (2023)
- 1D & 2D ν_{μ} CC1p0 π Generalized Imbalance @ BNB arXiv:2310.06082, submitted to PRD
- 1D ν_µ CC1p0π @ BNB
 <u>Phys. Rev. Lett. 125, 201803 (2020)</u>
- 1D ν_μ CC2p @ BNB <u>arXiv:2211.03734</u>
- 1D ν_µ CCNp0π @ BNB
 Phys. Rev. D102, 112013 (2020)

15 cross section publications and way more to come!





Already Public Results

CC inclusive

- 1D ν_µ CC inclusive @ BNB
 Phys. Rev. Lett. 123, 131801 (2019)
- 1D ν_µ CC E_ν @ BNB
 <u>Phys. Rev. Lett. 128, 151801 (2022)</u>
- 3D CC E_v @ BNB <u>arXiv:2307.06413</u>, submitted to PRL
- 1D v_e CC inclusive @ NuMI <u>Phys. Rev. D105, L05110</u> <u>Phys. Rev. D104, 052002</u>

$CC0\pi$

- 1D ν_e CCNp0π @ BNB Phys. Rev. D 106, L051102 (2022)
- 1D & 2D v_{μ} CC1p0 π Transverse Imbalance @ BNB <u>Phys. Rev. Lett. 131, 101802 (2023)</u>
 - Phys. Rev. D 108, 053002 (2023)
- 1D & 2D v_{μ} CC1p0 π Generalized Imbalance @ BNB

Opportunity to extensively benchmark neutrino event generator predictions

Pion production

• ν_µ NCπ⁰ @ BNB <u>Phys. Rev. D 107, 012004 (2023)</u> arX1v:2211.03/3

1D ν_µ CCNp0π @ BNB
 Phys. Rev. D102, 112013 (2020)

Rare channels

- η production @ BNB, submitted to PRL <u>arXiv:2305.16249</u>
- Λ production @ NuMI <u>Phys. Rev. Lett. 130, 231802 (2023</u>)

15 cross section publications and way more to come!

µBooNE

Nuclear Effects in Event Generators

Rev. Mod. Phys. 89, 045002 (2017)

Struck nucleon motion in argon



Double-Differential Single-Proton Knockout



- First double-differential single-proton cross section measurement on argon
- Identified kinematic variables and phase-space regions with sensitivity to Fermi motion & final state interactions
- Uses ~50% of available MicroBooNE data sets & Booster Neutrino Beam (BNB)

Phys. Rev. Lett. 131, 101802 (2023) Phys. Rev. D 108, 053002 (2023) arXiv:2310.06082

Single-Proton Knockout



- Dominated by Charged Current Quasi-elastic (CCQE) interactions
- Simple single muon-proton events
- Dominant at MicroBooNE energies

CC1p0π Quasielastic-like Signal Definition

Ranges driven by minimum track length, track containment, hadronic reinteractions, and systematics

• 1 muon

 $100 < P_{u} < 1200 \text{ MeV/c}$

• 1 proton

 $300 < P_p < 1000 \text{ MeV/c}$

- No π^{\pm} with $P_{\pi} > 70 \text{ MeV/c}$
- No π^0 or heavier mesons
- Any number of neutrons

9051 CC1p0π candidate data events CC1p0π ~10% efficiency ~70% purity



MC: GENIE v3.0.6 G18_10a_02_11b + tune* Nieves QE & MEC, Berger Sehgal RES

<u>Phys. Rev. Lett. 131, 101802 (2023)</u> <u>Phys. Rev. D 108, 053002 (2023)</u> * <u>Phys. Rev. D 105, 072001 (2022)</u>







Transverse missing momentum $\delta \mathbf{p}_{\mathrm{T}} = | \mathbf{p}_{\mathrm{T}}^{\mu} + \mathbf{p}_{\mathrm{T}}^{p} | = 0$

Transverse projections equal and opposite due to momentum conservation





Transverse missing momentum $\delta p_T = | \mathbf{p}_T^{\mu} + \mathbf{p}_T^{p} | > 0$

Broad distribution due to initial nucleon motion and other nuclear effects







Transverse Missing Momentum δp_{T}





- S = Signal, B = Background
- **QE** dominance in peak below Fermi momentum (~250 MeV/c)
- MEC/RES mainly in high momentum tail

GENIE v3.0.6 G18_10a_02_11b + tune* Nieves QE & MEC, Berger Sehgal RES

32

Transverse Orientation $\delta \alpha_{_{\rm T}}$

* <u>Phys. Rev. D 105, 072001 (2022)</u>





- + $\delta \alpha_{_{\rm T}}$ asymmetry due to proton FSI
- MEC/RES fractional contribution enhanced in ~180° region

GENIE v3.0.6 G18_10a_02_11b + tune* Nieves QE & MEC, Berger Sehgal RES ³³

Transverse Orientation $\delta \alpha_{_{\rm T}}$

Phys. Rev. D 108, 053002 (2023)

* Phys. Rev. D 105, 072001 (2022)





Need to move from event distributions to cross sections→ unfolding

Cross Section Extraction with Wiener SVD Unfolding

JINST 12 P10002 (2017)

Input Quantities

- Measurement (Data)
- Background (Cosmics + MC)
- Response Matrix (MC)
- Total Covariance Matrix (MC)



Cross Section Extraction with Wiener SVD Unfolding JINST 12 P10002 (2017)

Input Quantities

- Measurement (Data)
- Background (MC)
- Response Matrix (MC)
- Total Covariance Matrix (MC)

Probability that a generated event is reconstructed and selected

Diagonal matrix with flat ~6% efficiency


Cross Section Extraction with Wiener SVD Unfolding

Input Quantities

- Measurement (Data)
- Background (MC)
- Response Matrix (MC)
- Total Covariance Matrix (MC)

Includes information on statistical and systematic uncertainties



Uncertainties



- + Statistical (1.5%)
- + Number of argon targets (1%)

Total (11%)

Systematics-dominated analysis

Cross Section Extraction with Wiener SVD Unfolding

- Output quantities in regularized space
- Unfolded data spectrum
- Smearing Matrix A_C
 *Applied on theory predictions and included in data release



Cross Section Extraction with Wiener SVD Unfolding

- Output quantities in regularized space
- Unfolded data spectrum
- Smearing Matrix A_c

*Applied on theory predictions and included in data release

0.0	-0.03	-0.06.0.05.0.09.0	0.08 0.03	0.00	0.10	0.22	0.27	1	L
0.8 ت	0.15 0.14 0.07 0.08	3 0.14 0.19 0.20 0	0.16 0.10	0.12	0.27	0.38	0.42	-().8
0.6	- 0. 00 0.06 0.16 0.19	0.13 0.04 0.00 (0.02 0.15	0.33	0.40	0.24	0.16	-() 6
	-0.05-0.02 0.05 0.09	0.05 -0.02 -0.01	0.12 0.33	0.39	0.25	0.05	-0.04		
$d_{0} = 0.4$	0.13 0.12 -0.03-0.05	5 0.05 0.15 0.26 (0.00 0.13 0.19 (0.25 0.09 0.10 -0.03	-0.03 -0.10	-0.04 -0.07	0.01	0.04	-().4
	0.03 -0.03 -0.08 -0.00 0.02 0.04 0.07 0.20	0 0.19 0.27 0.21 (0 0.30 0.19 0.11 (0.10 0.04 0.05 0.02	0.02 0.03	0.04 0.04	0.06 0.03	0.06 0.03	-0).2
<u>ц</u> 0.2	0.01 0.13 0.34 0.41 -0.01 0.15 0.34 0.25 0.22 0.32 0.20 0.07	0.21 0.03 -0.01 (0.05 -0.04-0.05- 0.04 0.01 0.01 (0.00 0.04 0.02 0.04 0.01 -0.01	0.04 0.04 -0.01	0.01 0.01 -0.00	-0.03 -0.03 0.01	-0.04 -0.04 0.01	_()
С	0.61 0.52 0.20 0.08	2.0.12 0.15 0.14	0.07 -0.01	-0.04	-0.02	0.01	0.03		
\tilde{T} rue δp_{T} [GeV/c]									

Transverse Missing Momentum δp_{T} Cross Section



High Statistics→Into the Multiverse!

- Extension to 2D for the first time on argon
- Probe regions with greater model discrimination power



High Statistics→Into the Multiverse!

- Extension to 2D for the first time on argon



High Statistics→Into the Multiverse!

QE-dominated region

Phys. Rev. Lett. 131, 101802 (2023)

* Phys. Rev. D 105, 072001 (2022)





- No high transverse missing momentum tail
- Ideal part of phase-space to study Fermi motion
- Results consistent with local Fermi gas distribution

G18 = GENIE v3.0.6 G18_10a_02_11b + tune* GiBUU = GiBUU 2021

High Statistics→Into the Multiverse! MEC/RES/FSI-dominated





- FSI predictions in good agreement with data
- Minimal no-FSI contributions at high δp_T
- High $\delta \alpha_T \&$ high δp_T part of phase-space ideal to test FSI / multinucleon effects

G18 = GENIE v3.0.6 G18_10a_02_11b + tune* GiBUU = GiBUU 2021

CC1p0π TKI Summary

- First single- and double- differential neutrino-argon cross section measurements in TKI
- Fermi motion studied with 2D measurement in $\delta p_{\rm T}$ with $\delta \alpha_{\rm T} < 45^o$
- FSI & multinucleon effects studied with 2D measurement in δp_T with $135^\circ < \delta \alpha_T < 180^\circ$
- Way more single- and double-differential results in <u>Phys. Rev. Lett. 131, 101802 (2023)</u> and <u>Phys. Rev. D 108, 053002 (2023)</u>!



CC1p0π TKI Summary



Phys. Rev. C 95, 065501 (2017)

arXiv:2310.06082

• Extension to 3D by considering longitudinal component of missing momentum and calorimetric assumption on the incoming energy

48

• Extension to 3D by considering longitudinal component of missing momentum and calorimetric assumption on the incoming energy

<u>Phys. Rev. C 95, 065501 (2017)</u>

arXiv:2310.06082

BE = 30.9 MeV

$$E_{\text{cal}} = E_{\mu} + K_{p} + B$$
$$\vec{q} = E_{\text{cal}}\hat{z} - \vec{p}_{\mu}$$
$$p_{L} = p_{L}^{\mu} + p_{L}^{p} - E_{\text{cal}}$$

• Extension to 3D by considering longitudinal component of missing momentum and calorimetric assumption on the incoming energy



$$p_n = |\vec{p}_n| = \sqrt{p_L^2 + \delta p_T^2}$$
$$\alpha_{3D} = \cos^{-1} \left(\frac{\vec{q} \cdot \vec{p}_n}{|\vec{q}| |\vec{p}_n|} \right)$$

arXiv:2310.06082

Phys. Rev. C 95, 065501 (2017)

- Extension to 3D by considering longitudinal component of missing momentum and calorimetric assumption on the incoming energy
- Extensively tested again several event generators and model configurations





Name	Generator / Configuration
Gv2	GENIE v2.12.10
G18	GENIE v3.0.6 G18_10a_02_11a
G18T	G18 with tune
G21	GENIE v3.2.0 G21_11b_00_000
GiBUU	GiBUU 2021
NuWro	NuWro v19.02.1
NEUT	NEUT v5.4.0

- Extension to 3D by considering longitudinal component of missing moment and calorimetric assumption on the incoming energy
- Extensively tested again several event generators and model configurations





Name	Generator / Configuration
Gv2	GENIE v2.12.10
G18	GENIE v3.0.6 G18_10a_02_11a
G18T	G18 with tune
G21	GENIE v3.2.0 G21_11b_00_000
GiBUU	GiBUU 2021
NuWro	NuWro v19.02.1
NEUT	NEUT v5.4.0

Selected comparisons shown next using same CC1p0π selection

Missing momentum GENIE v3.0.6 G18_10a_02_11a (no tune)



- QE dominance due to $CC1p0\pi$ signal definition
- p_n pushes non-QE component to higher values

arXiv:2310.06082

Into the GKI multiverse!

QE-dominated region



- Tail significantly suppressed
- Consistent with local Fermi gas
- G18T results in lowest χ^2

Into the GKI multiverse!

MEC/RES/FSI-dominated



• Sharply peaked distribution to the right

- Driven by FSI
- GiBUU yields best result

CC1p0π GKI Summary

- Introduction of generalized kinematic imbalance (GKI) variables in 3D space
- Enhanced sensitivity to nuclear effects
- First single- and double-differential cross section GKI measurement ever with MicroBooNE
- G18T results in good description in QE-dominated regions
- GiBUU yields best performance in FSI-dominated regions
- Way more results in <u>arXiv:2310.06082</u>!



CC1p0π GKI Summary



57

SBN Experiments (SBND & ICARUS)

Common model configuration used by both experiments, systematics under development High statistics cross section measurements planned using both BNB & NuMI beamlines



DUNE Near Detector

2x2 Near Detector prototype Simulation studies on track multiplicity and inclusive selection Cross section measurement targeting specific interactions



Electron Complementarity





Backup Slides

Backup Slides: Table of Contents

MicroBooNE / Unfolding: 66-77 TKI PRD: 78-127 TKI PRL: 128-159 GKI arxiv: 160-182 e4v: 183-280



Figure 3.23: Expected sensitivity of DUNE to determination of the neutrino mass hierarchy (top) and discovery of CP violation, i.e., $\delta_{CP} \neq 0$ or π , (bottom) as a function of exposure in kt · MW · year, assuming equal running in neutrino and antineutrino mode, for a range of values for the ν_e and $\bar{\nu}_e$ signal normalization uncertainties from 5% \oplus 3% to 5% \oplus 1%. The sensitivities quoted are the minimum sensitivity for 100% of δ_{CP} values in the case of mass hierarchy and 50% (bottom left) or 75% (bottom right) of δ_{CP} values in the case of CP violation. The two bands on each plot represent a range of potential beam designs: the blue hashed band is for the CDR Reference Design and the solid green band is for the Optimized Design. Sensitivities are for true normal hierarchy; neutrino mass hierarchy and θ_{PP} octant are assumed to be unknown

3.6.3 Effect of Variation in Uncertainty

Figure 3.23 shows DUNE sensitivity to determination of neutrino mass hierarchy and discovery of CP violation as a function of exposure for several levels of signal normalization uncertainty. As seen in Figure 3.23, for early phases of DUNE with exposures less than 100 kt · MW · year, the experiment will be statistically limited. The impact of systematic uncertainty on the CP-violation sensitivity for large exposure is obvious in Figure 3.23; the ν_e signal normalization uncertainty must be understood at the level of $5\% \oplus 2\%$ in order to reach 5σ sensitivity for 75% of $\delta_{\rm CP}$ values with exposures less than ~900 kt · MW · year in the case of the Optimized Design. Specifically, the absolute normalization of the ν_{μ} sample must be known to ~5% and the normalization of the ν_e sample, relative to the $\bar{\nu}_e$, ν_{μ} , and $\bar{\nu}_{\mu}$ samples after all constraints from external, near detector, and far detector data have been applied, must be determined at the few-percent level. This level of systematic uncertainty sets the capability and design requirements for all components of the experiment, including the beam design and the near and far detectors.

Volume 2: The Physics Program for DUNE at LBNF

LBNF/DUNE Conceptual Design Report



MicroBooNE's physics program

Non-standard Cross-section Detector neutrino oscillations measurements physics, R&D

Beyond Standard Model physics!





120 GeV protons

- ~680 m from MicroBooNE
- ~8° off-axis from MicroBooNE

- 8 GeV protons
- ~470 m from MicroBooNE

• On-axis



 MicroBooNE collected BNB and NuMI data between 2015 and 2020

 ~50% of the dataset (Runs 1-3) used in first wave of results



Readout electronics and field response removal

<u>JINST 13, P07006 (2018)</u> <u>JINST 13, P07007 (2018)</u> From raw hits to particle reconstruction

Pandora Pattern Recognition

Eur. Phys. J. C78, 1, 82 (2018)





Eur. Phys. J. C 79 673 (2019)



TABLE IV. Tuned parameter values and uncertainties after fitting to T2K CC0 π data for the nominal simulation and three tunes that build to the final four parameter tune. Note that postfit χ^2 values are quoted here only for the 58 bins included in the fit (excluding the highest muon momentum bin in each cos θ bin), and using diagonal elements of the covariance matrix only. In the text and figures, pre- and postfit χ^2 comparisons are also quoted for the full T2K dataset of 67 bins. "Norm." is an abbreviation for normalization.

	MaCCQE fitted value	CC2p2h Norm. fitted value	CCQE RPA Strength fitted value	CC2p2h Shape fitted value	$\frac{\text{T2K}}{\chi^2_{\text{diag}}/\text{N}_{\text{bins}}}$
Nominal (untuned)	0.961242 GeV	1	100%	0	106.7/58
Fit MaCCQE + CC2p2h Norm.	$1.14\pm0.07~{ m GeV}$	1.61 ± 0.19	100% (fixed)	0 (fixed)	71.8/58
Fit MaCCQE + CC2p2h Norm + CCQE RPA Strength	$1.18\pm0.08~\text{GeV}$	1.12 ± 0.38	$(64 \pm 23)\%$	0 (fixed)	69.7/58
Fit MaCCQE + CC2p2h Norm + CCQE RPA Strength + CC2p2h Shape	$1.10\pm0.07~\text{GeV}$	1.66 ± 0.19	$(85\pm20)\%$	$1^{+0}_{-0.74}$	52.5/58



FIG. 7. Correlations between parameters after fitting to T2K $CC0\pi$ data.

Phys. Rev. D 105, 072001 (2022)


Unfolding problem

• In practice, the data unfolding problem starts with

$$\chi^2(s) = (\boldsymbol{m} - \boldsymbol{r} \cdot s)^T Cov^{-1}(\boldsymbol{m} - \boldsymbol{r} \cdot s)$$

- m : measured spectrum, m-dimensional vector
- *s* : unknown spectrum, to be unfolded, *n*-dimensional vector
- r: smearing (response) matrix, m X n and $m \ge n$
- Cov : covariance matrix containing all statistical and systematic uncertainties associated with m and r.
- Cholesky decomposition: $Cov^{-1} = Q^T Q$, Q is a lower triangular matrix

$$\chi^2(s) = (M - R \cdot s)^T \cdot (M - R \cdot s)$$

Pre-scaling

•
$$M \coloneqq Q \cdot m$$

• $R \coloneqq Q \cdot r$

Solution (direct inversion) $\hat{s} = (R^T R)^{-1} R^T M$ $\hat{s} = (R^T R)^{-1} R^T (R \cdot s_{true} + N)$

The response matrix R is unnecessary to be a square matrix

Unfolding problem



This is one unbiased solution (direct inversion) to an unfolding problem. However, it has catastrophic oscillations, i.e. huge variance, in the unfolded spectrum.

- Decrease the number of bins to suppress the "oscillation" --> Nyquist theorem
- Trade-off bias and variance to suppress the "oscillation" --> e.g. regularization [unfolding method]

Wiener-SVD unfolding

 10^{3}

variance

 To automatically minimize the Mean Square Error (MSE) given a model S

$$MSE = E\left[\left(\hat{S} - S\right)^2\right] = E\left[\left(F \cdot \frac{M}{R} - S\right)^2\right] = E\left[\left(F \cdot S + F \cdot \frac{N}{R} - S\right)^2\right]$$

$$= E\left[\left((F - I) \cdot S\right)^2 + \left(F \cdot \frac{N}{R}\right)^2\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

$$F = "filter" = additional smearing matrix = regularization w C_0^{+} + F \cdot \frac{N}{R}\right]$$

0.6 0.8 1

1 √τ

75

bias versus variance with minimum MSE



FIG. 1. The log-likelihood ratio (LLR) particle identification (PID) score distribution used to tag the muon and proton candidates.



FIG. 2. (Top) the proton candidate LLR PID score distribution, illustrating the fitness of a cut at LLR PID < 0.05 to reject cosmic and non-CC1p0 π background events. (Bottom) the muon candidate LLR PID score distribution, illustrating a peak close to one. Only statistical uncertainties are shown on the data. The bottom panel shows the ratio of data to prediction.



FIG. 1. (Left) the proton candidate LLR PID score distribution, illustrating the particle composition of the variable. (Right) the muon candidate LLR PID score distribution, illustrating a peak close to one. Only statistical uncertainties are shown on the data. The bottom panel shows the ratio of data to prediction.



FIG. 3. Muon momentum reconstruction (top) before and (bottom) after the application of the muon momentum quality cut using contained muon tracks.






































































Calorimetric Energy



$$E^{Cal} = E_{\mu} + T_p + BE$$

•
$$E_{\mu}$$
 = muon energy

- T_p = proton kinetic energy
- BE = 40 MeV binding energy
- Peak at ~0.7 GeV

Calorimetric Energy Cross Section



$$E^{Cal} = E_{\mu} + T_p + BE$$

•
$$E_{\mu} = muon energy$$

- T_p = proton kinetic energy
- BE = 40 MeV binding energy
- All generators yield good agreement

116



- QE dominated region
- All generators yield good agreement



arXiv:2301.03700

- MEC/RES dominated region
- Similar shapes
- Normalization differences
- Still reasonable χ^2 !





-Total-NuWro-Stat-LY -TPC

-G4 -Flux -Dirt-POT-NTarget-MCStat

-Total-NuWro-Stat-LY -TPC -SCERecomb2-XSec

120

-SCERecomb2-XSec



Parameter	Description	CV	1σ Uncertainty	Contributing Uncertainty (%)
Quasi-Elastic Parameters	3			
MaCCQE	CCQE axial mass	$1.10~{\rm GeV}$	$\pm 0.1 \ {\rm GeV}$	0.038
RPA CCQE	Strength of the RPA correction	0.151	± 0.4	2.094
MaNCEL	Axial mass for NCEL	$0.961242~{\rm GeV}$	$\pm 25\%$	0.348
EtaNCEL	Empirical parameter used to account for sea quark contribution to NCEL form factor	0.12	$\pm 30\%$	0.010
AxFFCCQEshape	Parametrisation of the nucleon axial form factor	Dipole	z-expansion	0.022
VecFFCCQEshape	Parametrisation of the nucleon vector form factors	BBA07	Dipole	0.051
MEC Parameters				
NormCCMEC	Energy-independent normalization for CCMEC	1.66	± 0.5	1.832
NormNCMEC	Energy-independent normalization for NCMEC	1	$\pm 20\%$	0.129
FracPNCCMEC	Fraction of initial nucleon pairs that are pn $(0 = \text{Valencia})$	0	$\pm 20\%$	0.041
FracDeltaCCMEC	Relative contribution of Δ diagrams to total MEC cross section (0 = Valencia)	0	$\pm 30\%$	0.124
XSecShape CCMEC	Changes shape of differential cross section	1.0	0.0	2.273
DecayAngMEC	Changes angular distribution of nucleon cluster	Isotropic	$\cos^2\vartheta$ in rest frame	0.693
Resonant Parameters				
MaCCRES	CCRES axial mass	$1.120~{\rm GeV}$	± 0.2	0.986
MvCCRES	Shape-only CCRES axial mass	$0.840~{\rm GeV}$	± 0.1	0.775
MaNCRES	NCRES axial mass	$1.120~{\rm GeV}$	± 0.2	0.969
MvNCRES	NCRES vector mass.	$0.840~{\rm GeV}$	± 0.1	0.395
ThetaDelta2Npi	Interpolates angular distribution for $\Delta \rightarrow N + \pi$	Rein-Sehgal	Isotropic	1.533
ThetaDelta2NRad	Interpolates angular distribution for $\Delta \rightarrow N + \gamma$	Rein-Sehgal	$\cos^2 \vartheta$	0.016

Parameter	Description	CV	1σ Uncertainty	Contributing Uncertainty (%)
Non-Resonant Parameters				
NonRESBGvpNC1pi	Non-resonant background normalization for $\nu p~{\rm NC}1\pi$	0.1	± 0.5	0.041
NonRESBGvpNC2pi	Non-resonant background normalization for $vp~{\rm NC}2\pi$	1	± 0.5	0.096
NonRESBGvnNC1pi	Non-resonant background normalization for $\nu n~{\rm NC}1\pi$	0.3	± 0.5	0.390
NonRESBGvnNC2pi	Non-resonant background normalization for $\nu n~{\rm NC}2\pi$	1	± 0.5	0.022
NonRESBGvbarpNC1pi	Non-resonant background normalization for $\bar\nu p~{\rm NC}1\pi$	0.3	± 0.5	0.010
NonRESBGvbarpNC2pi	Non-resonant background normalization for $\bar\nu p~{\rm NC}2\pi$	1	± 0.5	0.010
NonRESBGvbarnNC1pi	Non-resonant background normalization for $\bar\nu n~{\rm NC}1\pi$	0.1	± 0.5	0.010
NonRESBGvbarnNC2pi	Non-resonant background normalization for $\bar\nu n~{\rm NC}2\pi$	1	± 0.5	0.010
NonRESBGvpCC1pi	Non-resonant background normalization for vp CC1 $\!\pi$	0.007713	± 0.5	0.014
NonRESBGvpCC2pi	Non-resonant background normalization for $\nu p~{\rm CC}2\pi$	0.787999	± 0.5	0.059
NonRESBGvnCC1pi	Non-resonant background normalization for $\nu n~{\rm CC1}\pi$	0.127858	± 0.5	0.217
NonRESBGvnCC2pi	Non-resonant background normalization for vn ${\rm CC}2\pi$	2.11523	± 0.5	0.079
NonRESBGvbarpCC1pi	Non-resonant background normalization for $\bar\nu p~{\rm CC1}\pi$	0.127858	± 0.5	0.013
NonRESBGvbarpCC2pi	Non-resonant background normalization for $\bar\nu p~{\rm CC}2\pi$	2.11523	± 0.5	0.010
NonRESBGvbarnCC1pi	Non-resonant background normalization for $\bar{\nu}n~{\rm CC1}\pi$	0.007713	± 0.5	0.010
NonRESBGvbarnCC2pi	Non-resonant background normalization for $\bar\nu n~{\rm CC}2\pi$	0.787999	± 0.5	0.010
AhtBY	$A_{\rm HT}$ higher-twist parameter in the Bodek-Yang model scaling variable ξ_w	0.538	± 0.25	0.010
BhtBY	BHT higher-twist parameter in the Bodek-Yang model scaling variable $\xi_{\rm w}$	0.305	± 0.25	0.010
CV1uBY	CV1u valence GRV98 PDF correction parameter in the Bodek-Yang model	0.291	± 0.3	0.010
CV2uBY	CV2u valence GRV98 PDF correction parameter in the Bodek-Yang model	0.189	± 0.4	0.010

AGKYxF1pi	Hadronization parameter, applicable to true DIS interactions only	-0.385	± 0.2	0.108
AGKYpT1pi	adronization parameter, applicable to true DIS interactions only $1/6.625 \pm 0.03$		± 0.03	0.034
Final State Interaction Parameters	5			
MFP_{π}	π mean free path	0 ± 0.2		0.032
MFPN	Nucleon mean free path 0 ± 0.2		± 0.2	1.212
$FrCEx_{\pi}$	Fractional cross section for π charge exchange 0 ± 0.5		± 0.5	0.159
\texttt{FrInel}_{π}	Fractional cross section for π inelastic scattering	0	± 0.4	0.133
$FrAbs_{\pi}$	Fractional cross section for π absorption	0	± 0.3	0.906
FrCEx _N	Fractional cross section for nucleon charge exchange	0	± 0.2	0.953
\texttt{FrInel}_N	Fractional cross section for nucleon inelastic scattering	0	± 0.5	0.289
FrAbs_N	Fractional cross section for nucleon absorption	0	± 0.4	0.906
Delta Resonant Decay Parameters				
RDecBR1gamma	Normalization for $\Delta \to \gamma$ decays	Nominal BR	± 0.5	0.042
RDecBR1eta	Normalization for $\Delta \rightarrow \eta$ decays	Nominal BR	± 0.5	0.513
Coherent Parameters				
NormCCCOH	Scaling factor for CCCOH π production total cross section	Nominal	100% increase	0.027
NormNCCOH	Scaling factor for NCCOH π production total cross section	Nominal	100% increase	0.016

Hadronisation Parameters

Variation	Description	
Wire Mod x position	Wire modification of x position	
Wire Mod (y,z) position	Wire modification of (y,z) position	
Wire Mod θ_{XZ}	Wire modification of angle in XZ plane	
Wire Mod θ_{YZ}	Wire modification of angle in YZ plane	
Light Yield Attenuation	Attenuation of LY response in detector over time	
Light Yield Down	Turn down the light yield in the detector by 25%	
Light Yield Rayleigh	Increase Rayleigh scattering length from 60 cm to 90 cm	
Recombination	Reduce value of β' in the Modified Box Model	
SCE	Use an alternative Space Charge Map	





FIG. 1. The flux-integrated (a) single- and (b-c) double- (in $\delta \alpha_T$ bins) differential CC1p0 π cross sections as a function of the transverse missing momentum δp_T . Inner and outer error bars show the statistical and total (statistical and shape systematic) uncertainty at the 1 σ , or 68%, confidence level. The gray band shows the separate normalization systematic uncertainty. Colored lines show the results of theoretical cross section calculations with (solid line) and without (dashed line) FSI based on the GENIE (blue) and GiBUU (orange) event generators.



FIG. 2. The flux-integrated (a) single- and (b-c) double- (in δp_T bins) differential CC1p0 π cross sections as a function of the angle $\delta \alpha_T$. Inner and outer error bars show the statistical and total (statistical and shape systematic) uncertainty at the 1 σ , or 68%, confidence level. The gray band shows the separate normalization systematic uncertainty. Colored lines show the results of theoretical cross section calculations with a number of FSI-modeling choices based on the GENIE event generator.



FIG. 3. The flux-integrated (a) single- and (b-c) double- (in $\delta p_{T,y}$ bins) differential CC1p0 π cross sections as a function of the transverse three-momentum transfer component, $\delta p_{T,x}$. Inner and outer error bars show the statistical and total (statistical and shape systematic) uncertainty at the 1 σ , or 68%, confidence level. The gray band shows the separate normalization systematic uncertainty. Colored lines show the results of theoretical cross section calculations with a number of event generators. The standard deviation (σ_{Data}) of a Gaussian fit to the data is shown on each panel.

QE-dominated region

Phys. Rev. Lett. 131, 101802 (2023)

* Phys. Rev. D 101, 033003 (2020)





- Flat distribution indicative of absence of proton FSI
- Shape and normalization differences across FSI models

G21 = GENIE v3.0.6 G21_11b_00_000 SuSAv2 QE & MEC*, hA/hN/G4 = FSI modeling options ¹²⁹

Transverse Orientation $\delta \alpha_{T}$ Cross Section

All events





- First neutrino-argon differential cross section in $\delta \alpha_{T}$
- Sensitive to proton FSI modeling
- Data favors FSI addition
- Shape differences observed

G21 = GENIE v3.0.6 G21_11b_00_000 SuSAv2 QE & MEC*, hA/hN/G4 = FSI modeling options ¹³⁰

arXiv:2301.03706

* <u>Phys. Rev. D 101, 033003 (2020)</u>

- Extension to 2D for the first time on argon
- Probe regions with greater model discrimination power





- Primarily contributions from MEC/RES & QE events undergoing FSI
- More assymmetric behavior compared to 1D result
- No-FSI contribution lower than FSI ones
- High $\delta \alpha_{\pi}$ & high δp_{π} part of phase-space ideal to test FSI / multinucleon effect sensitivity

G21 = GENIE v3.0.6 G21 11b 00 000 SuSAv2 QE & MEC^{*}, hA/hN/G4 = FSI modeling options ¹³¹

* Phys. Rev. D 101, 033003 (2020)

VI. MODELING CONFIGURATIONS

The nominal MC neutrino interaction prediction (G18) uses the local Fermi gas (LFG) model [50], the Nieves CCQE scattering prescription [51] which includes Coulomb corrections for the outgoing muon [52] and random phase approximation (RPA) corrections [53]. Additionally, it uses the Nieves MEC model [54], the KLN-BS RES [55–58] and Berger-Sehgal coherent (COH) [59] scattering models, the hA2018 FSI model [60], and MicroBooNE-specific tuning of model parameters [38].

Our results are also compared to a number of alternative event generators. GiBUU 2021 (GiBUU) uses similar models, but they are implemented in a coherent way by solving the Boltzmann-Uehling-Uhlenbeck transport equation [61]. The modeling includes the LFG model [50], a standard CCQE expression [62], an empirical MEC model and a dedicated spin dependent resonance amplitude calculation following the MAID analysis [61]. The DIS model is from PYTHIA [63]. GIBUU's FSI treatment propagates the hadrons through the residual nucleus in a nuclear potential which is consistent with the initial state. NuWro v19.02.2 (NuWro) uses the LFG model [50], the Llewellyn Smith model for QE events [64], the Nieves model for MEC events [65], the Adler-Rarita-Schwinger formalism to calculate the Δ resonance explicitly [58], the BS COH [59] scattering model and an intranuclear cascade model for FSI [65]. NEUT v5.4.0 (NEUT) uses the LFG model [50], the Nieves CCQE scattering prescription [51], the Nieves MEC model [54], the BS RES [55–58] and BS COH [59] scattering models, and FSI with Oset medium corrections for pions [35, 36].

In addition to the alternative event generators, our results are compared to a number of different GENIE configurations. These include an older version, GENIE v2.12.10 (Gv2) [35, 36], which uses the Bodek-Ritchie Fermi Gas model, the Llewellyn Smith CCQE scattering prescription [64], the empirical MEC model [66], a Rein-Sehgal RES and COH scattering model [67], and a data driven FSI model denoted as "hA" [68]. Another model, "Untuned", uses the GENIE v3.0.6 G18_10a_02_11a configuration without additional MicroBooNE-specific Finally, the newly added theory-driven tuning. GENIE v3.2.0 G21_11b_00_000 configuration (G21) is shown. This includes the SuSAv2 prediction for the QE and MEC scattering parts [69] and the hN2018 FSI model [70]. The modeling options for RES, DIS, and COH interactions are the same as for G18.

To quantify the data-simulation agreement, the χ^2 /bins ratio data comparison for each generator is shown on all the figures and is calculated by taking into account the total covariance matrix. Ratios close to unity are indicative of a sufficiently accurate modeling performance. Theoretical uncertainties on the models themselves are not included.



FIG. 1. Fake data studies for δp_T using (left) NuWro, (center) GENIE without the MicroBooNE tune (NoTune), and (right) twice the weights for MEC events (TwiceMEC) as fake data samples.



FIG. 2. Fake data studies for $\delta \alpha_{\rm T}$ using (left) NuWro, (center) GENIE without the MicroBooNE tune (NoTune), and (right) twice the weights for MEC events (TwiceMEC) as fake data samples.



FIG. 3. Fake data studies for $\delta p_{T,x}$ using (left) NuWro, (center) GENIE without the MicroBooNE tune (NoTune), and (right) twice the weights for MEC events (TwiceMEC) as fake data samples.



FIG. 4. Fake data studies for δp_T in $\delta \alpha_T$ bins using (left) NuWro, (center) GENIE without the MicroBooNE tune (NoTune), and (right) twice the weights for MEC events (TwiceMEC) as fake data samples.



FIG. 5. Fake data studies for $\delta \alpha_{\rm T}$ in $\delta p_{\rm T}$ bins using (left) NuWro, (center) GENIE without the MicroBooNE tune (NoTune), and (right) twice the weights for MEC events (TwiceMEC) as fake data samples.



FIG. 6. Fake data studies for $\delta p_{T,x}$ in $\delta p_{T,y}$ bins using (left) NuWro, (center) GENIE without the MicroBooNE tune (NoTune), and (right) twice the weights for MEC events (TwiceMEC) as fake data samples.



FIG. 7. (Left) extracted cross section as a function of the vertex z distribution. (Right) vertex z efficiency function.



FIG. 8. Cross section interaction breakdown for (top) all the selected events, (middle) events with $\delta \alpha_{\rm T} < 45^{\circ}$, and (bottom) events with $135^{\circ} < \delta \alpha_{\rm T} < 180^{\circ}$. The breakdown is shown for (first column) the G18 configuration with FSI effects, (second column) the G18 configuration without FSI effects, (third column) GiB with FSI effects, and (forth column) GiB without FSI effects.



FIG. 9. Cross section interaction breakdown for (top) all the selected events, (middle) events with $\delta p_T < 0.2 \,\text{GeV/c}$, and (bottom) events with $\delta p_T > 0.4 \,\text{GeV/c}$. The breakdown is shown for (first column) the G21 hA configuration with the hA2018 FSI model, (second column) the G21 hN configuration with the hN FSI model, (third column) the G21 G4 configuration with the G4 FSI model, and (forth column) the G21 No FSI configuration without FSI effects.



FIG. 10. Cross section interaction breakdown for (top) all the selected events, (middle) events with $\delta p_{T,y} < -0.15 \text{ GeV/c}$, and (bottom) events with $-0.15 < \delta p_{T,y} < 0.15 \text{ GeV/c}$. The breakdown is shown for (first column) the G18 LFG configuration, (second column) the G18 RFG configuration, (third column) the G18 EffSF configuration, and (forth column) the G18 No RPA configuration.



FIG. 11. The flux-integrated (a) single- and (b-c) double- (in $\delta \alpha_{\rm T}$ bins) differential CC1p0 π cross sections as a function of the transverse missing momentum, $\delta p_{\rm T}$. Inner and outer error bars show the statistical and total (statistical and shape systematic) uncertainty at the 1 σ , or 68%, confidence level. The gray band shows the separate normalization systematic uncertainty. Colored lines show the results of theoretical cross section calculations with a number of G21 FSI modeling variations.



FIG. 12. The flux-integrated (a) single- and (b-c) double- (in δp_T bins) differential CC1p0 π cross sections as a function of the angle $\delta \alpha_T$. Inner and outer error bars show the statistical and total (statistical and shape systematic) uncertainty at the 1 σ , or 68%, confidence level. The gray band shows the separate normalization systematic uncertainty. Colored lines show the results of theoretical cross section calculations with a number of QE-modeling choices based on the GENIE event generator.


FIG. 13. The flux-integrated (a) single- and (b-c) double- (in $\delta p_{T,y}$ bins) differential CC1p0 π cross sections as a function of the transverse three-momentum transfer component, $\delta p_{T,x}$. Inner and outer error bars show the statistical and total (statistical and shape systematic) uncertainty at the 1 σ , or 68%, confidence level. The gray band shows the separate normalization systematic uncertainty. Colored lines show the results of theoretical cross section calculations with a number of G18 FSI modeling variations. The standard deviation (σ_{Data}) of a Gaussian fit to the data is shown on each panel.

Transverse Component $\delta p_{T,x}$





 $\delta p_{T,x} = \delta p_T \bullet sin \delta \alpha_T$

- Symmetric around 0 GeV/c
- **QE** dominance in central region
- MEC/RES events primarily in the tail

Transverse Component $\delta p_{T,x}$ Cross Section





 $\delta p_{T,x} = \delta p_T \bullet \sin \delta \alpha_T$

- G18 LFG = GENIE v3.0.6 G18_10a_02_11a (G18) + uB Tune with local Fermi gas
- G18 no-RPA = G18 w/o RPA effects
- G18 RFG = G18 with relativistic Fermi gas (RFG)
- G18 EffSF = G18 with spectral function (EffSF)





 $\delta p_{T,y} = \delta p_T \cdot \cos \delta \alpha_T$

- Asymmetric due to $\delta \alpha_{T}$ enhancement at ~180°
- QE dominance in central region
- Spread of tail sensitive to FSI strength & MEC/RES

High Statistics→Into the Multiverse!

- First neutrino-argon differential cross section in TKI variables
- Sensitive to initial nucleon motion & proton FSI modeling



δp_{T,x}

δp_{T,v}













MicroBooNE Data 6.79e+20 POT
MC uB Tune



Figure 56: Closure test for δp_T .



Generalized Kinematic Imbalance (GKI)

Straightforward extension to 3D by considering longitudinal component of missing moment However, an assumption on the incoming energy has to be made First attempt in <u>Phys. Rev. C 95, 065501 (2017)</u> using CCQE interactions off a bound stationary neutron

Leveraging calorimetric energy estimator definition

$$E_{\rm cal} = E_{\mu} + K_p + B$$

to obtain the longitudinal component of missing momentum

$$p_L = p_L^{\mu} + p_L^{p} - E_{\text{cal}}$$

and the energy transfer vector

$$\vec{q} = E_{\rm cal}\hat{z} - \vec{p}_{\mu}$$







-G18 --- Gv2 --- NEUT --- NuWro --- GiBUU



-G18 --- Gv2 --- NEUT --- NuWro --- GiBUU

0.15

 $\frac{d\sigma}{d\alpha_{3D}^{3D}} \left[10^{-38} \frac{cm^2}{deg Ar} \right]$

20









-G18 --- Gv2 --- NEUT --- NuWro --- GiBUU





-G18 --- Gv2 --- NEUT --- NuWro --- GiBUU



-G18 --- Gv2 -- NEUT --- NuWro --- GiBUU













0.8

0.6

















Into the GKI multiverse!



- Extended tail to higher values
- FSI-dominated region

- GiBUU shift to the right, yet lowest χ^2
- Gv2 yields worst agreement
- Other generators yield comparable ratios

Into the GKI multiverse!



- QE-dominated region
- Most generators result in comparable results

- Gv2 yields worst agreement
- G18T results in good agreement
Why electrons?

- Common vector current
- Identical nuclear effects
- Monoenergetic beams
- High statistics
 - Precision measurements

Any model must work for electrons, or it won't work for neutrinos !



Similar ν & e Distributions



Phys. Rev. D 103, 113003 (2021)

Jefferson Laboratory



- Electron beam accelerator facility
- Energies up to 12 GeV
- Using Hall B & CLAS detector



e4v Data-Mining With CLAS

- Charged particle threshold similar to ν tracking detectors
- ~50% of " 4π " coverage



e4v Data-Mining With CLAS

- Charged particle threshold similar to ν tracking detectors
- ~50% of " 4π " coverage
- Energies: 1, 2 & 4 GeV
- Targets: ⁴He, ¹²C, ⁵⁶Fe



Playing The QE-like Neutrino Game



1 proton (> 300 MeV/c)
No π[±] (> 70 MeV/c)

Phys. Rev. Lett. 125, 201803 (2020) arXiv:2301.03706 arXiv:2301.03700



- 1 proton (> 300 MeV/c)
 No π[±] (> 150 MeV/c)

Nature 599, 565-570 (2021)

Study energy reconstructionTest against GENIE event generator

QE Energy Reconstruction



Cherenkov detectors Assuming QE interaction Using lepton kinematics

$$E_{QE} = \frac{2M\epsilon + 2ME_l - m_l^2}{2(M - E_l + |k_l|\cos\theta_l)}$$

SK

QE Energy Reconstruction

C @ 1.16 GeV



- Relevant for T2K
- Overestimation of QE peak & RES tail

Calorimetric Energy Reconstruction



Tracking detectors Calorimetric sum Using all detected particles

$$\mathbf{E}_{cal} = \mathbf{E}_l + \mathbf{T}_p + \boldsymbol{\epsilon}_{\mathbf{B}}$$

Calorimetric Energy Reconstruction



 $\mathbf{E}_{cal} = \mathbf{E}_l + \mathbf{T}_p + \epsilon$

Nature 599, 565–570 (2021)

E_{cal} Nucleus & Energy Dependence



Nucleus & Energy Dependence



Nucleus & Energy Dependence



Nucleus & Energy Dependence



Transverse Momentum



Nature 599, 565-570 (2021)

$$\mathbf{P}_{\mathrm{T}} = \left| \mathbf{P}_{\mathrm{T}}^{e'} + \mathbf{P}_{\mathrm{T}}^{p} \right|$$



- P_T sensitivity to nuclear effects (fermi motion, final-state interactions, ...)
- Overestimation of QE peak & RES tail

Energy Reconstruction In P_T Slices













- Even more differential! Into the 3D multiverse Taking advantage of massive statistics
- More nuclear sensitivity variables Help decide what to tune
- $1p1\pi$ exclusive cross sections







- More with CLAS6
- Proton transparency studies
 - Observable that summarizes strength of intranuclear rescattering
 - Measurement study $T = N_{prot}^{detected} / N_{prot}^{PWIA}$
 - Test event generator FSI models in GENIE across multiple nuclear targets
 - Testing dependency of transparency calculations on PWIA normalization



- More with CLAS6
- Prepare suite of cross sections in same variables as neutrinos for first e-GENIE tune
 - Compare with neutrino $CC0\pi$ tune
- *Cal* data already being used to validate event generators





2206.11050 [hep-ph]



- New data with CLAS12 Targets: ⁴He, ¹²C, ⁴⁰Ar, ¹²⁰Sn
- 2 6 GeV beam energies
- More phase-space coverage ($\theta_e > 5^\circ$)
- Neutrons!







Why electrons?

- Common vector current
- Identical nuclear effects
- Monoenergetic beams
- High statistics
 - Precision measurements

Any model must work for electrons, or it won't work for neutrinos !



Similar ν & e Distributions



Phys. Rev. D 103, 113003 (2021)

Jefferson Laboratory



- Electron beam accelerator facility
- Energies up to 12 GeV
- Using Hall B & CLAS detector



e4v Data-Mining With CLAS

- Charged particle threshold similar to ν tracking detectors
- ~50% of " 4π " coverage



e4v Data-Mining With CLAS

- Charged particle threshold similar to ν tracking detectors
- ~50% of " 4π " coverage
- Energies: 1, 2 & 4 GeV
- Targets: ⁴He, ¹²C, ⁵⁶Fe



Playing The QE-like Neutrino Game



1 proton (> 300 MeV/c)
No π[±] (> 70 MeV/c)

Phys. Rev. Lett. 125, 201803 (2020) arXiv:2301.03706 arXiv:2301.03700



- 1 proton (> 300 MeV/c)
 No π[±] (> 150 MeV/c)

Nature 599, 565-570 (2021)

Study energy reconstructionTest against GENIE event generator

QE Energy Reconstruction



Cherenkov detectors Assuming QE interaction Using lepton kinematics

$$\mathbf{E}_{QE} = \frac{2\mathbf{M}\epsilon + 2\mathbf{M}\mathbf{E}_l - \mathbf{m}_l^2}{2\left(\mathbf{M} - \mathbf{E}_l + |\mathbf{k}_l|\cos\theta_l\right)}$$

SK

QE Energy Reconstruction

C @ 1.16 GeV



- Relevant for T2K
- Overestimation of QE peak & RES tail

Calorimetric Energy Reconstruction



Tracking detectors Calorimetric sum Using all detected particles

$$\mathbf{E}_{cal} = \mathbf{E}_l + \mathbf{T}_p + \boldsymbol{\epsilon}_{\mathbf{B}}$$

Calorimetric Energy Reconstruction



 $\mathbf{E}_{cal} = \mathbf{E}_l + \mathbf{T}_p + \epsilon$

Nature 599, 565–570 (2021)

E_{cal} Nucleus & Energy Dependence



216
Nucleus & Energy Dependence



Nucleus & Energy Dependence



Nucleus & Energy Dependence



Transverse Momentum



Nature 599, 565-570 (2021)

$$\mathbf{P}_{\mathrm{T}} = \left| \mathbf{P}_{\mathrm{T}}^{e^{*}} + \mathbf{P}_{\mathrm{T}}^{p} \right|$$



- P_T sensitivity to nuclear effects (fermi motion, final-state interactions, ...)
- Overestimation of QE peak & RES tail

Energy Reconstruction In P_T Slices



221











- Even more differential! Into the 3D multiverse Taking advantage of massive statistics
- More nuclear sensitivity variables Help decide what to tune
- $1p1\pi$ exclusive cross sections







227

- More with CLAS6
- Proton transparency studies
 - Observable that summarizes strength of intranuclear rescattering
 - Measurement study $T = N_{prot}^{detected} / N_{prot}^{PWIA}$
 - Test event generator FSI models in GENIE across multiple nuclear targets
 - Testing dependency of transparency calculations on PWIA normalization



- More with CLAS6
- Prepare suite of cross sections in same variables as neutrinos for first e-GENIE tune
 - Compare with neutrino $CC0\pi$ tune
- *Cal* data already being used to validate event generators





2206.11050 [hep-ph]



- New data with CLAS12 Targets: ⁴He, ¹²C, ⁴⁰Ar, ¹²⁰Sn
- 2 6 GeV beam energies
- More phase-space coverage ($\theta_e > 5^\circ$)
- Neutrons!







Mismodelling Impact On Mixing Parameters

Charged current cross sections obtained using GENIE for the DUNE near detector (left) and far detector (right) oscillated fluxes



Issues Identified & Fixed In G2018



SuSav2 Offers More Accurate Prediction



Probing The Neutrino Phase-Space With Electrons

QE Events



Consistent Treatment Of MEC Events With SuSav2

Unique chance to constraint one of least understood interaction channels



Electron results scaled by Q^4

Phys. Rev. D 103, 113003 (2021)

Inclusive C cross sections

Phys. Rev. D 103, 113003 (2021)



Inclusive C/Fe cross sections Phys. Rev. D 103, 113003 (2021)



Inclusive H cross sections

Phys. Rev. D 103, 113003 (2021)



Sanity Check With Inclusive Cross Sections



M.Khachatryan, A.Papadopoulou, et al. Nature 599, 565–570 (2021)

Detected Hadron Multiplicities



Q⁴ Scaling Effect



241

Available Nuclear Models

Phys. Rev. D 103, 113003 (2021)



242

SuSav2 Configuration / GEM21_11b_00_000

	Electrons	Neutrinos
QE	SuSav2	SuSav2
MEC	SuSav2	SuSav2
RES	Berger-Sehgal	Berger-Sehgal
DIS	AGKY	AGKY
FSI	hN2018	hN2018
Nuclear Model	Relativistic Mean Field	Relativistic Mean Field

G2018 Model Configuration

	Electrons	Neutrinos
QE	Rosenbluth	Nieves
MEC	Empirical	Nieves
RES	Berger-Sehgal	Berger-Sehgal
DIS	AGKY	AGKY
FSI	hA2018	hA2018
Nuclear Model	Local Fermi Gas	Local Fermi Gas

Closure Test

- Use GENIE files
- Filter specific topologies (e.g. $1p0\pi p + 1p1\pi$)
- Subtracted & True $1p0\pi$ are in good agreement



Well defined signal definition: Min θ_{e} Cut

(a) 1.1 GeV: $\theta = 17 + 7 / P$

@ 2.2 GeV: $\theta = 16 + 10.5 / P$

(a) 4.4 GeV: $\theta = 13.5 + 15 / P$

See backup for p / $\pi^{+/-}$ definitions

• We do not acceptance correct below min θ



Well defined signal definition: Min θ_{e} Cut

(a) 1.1 GeV: $\theta = 17 + 7 / P$

@ 2.2 GeV: $\theta = 16 + 10.5 / P$

@ 4.4 GeV: $\theta = 13.5 + 15 / P$

See backup for p / $\pi^{+/-}$ definitions



• We do not acceptance correct below min θ

Background Subtraction

Non-(e,e'p) interactions lead to multi-hadron final states Gaps can make them look like (e,e'p) events



Data Driven Correction

Non-(e,e'p) interactions lead to multi-hadron final states Gaps make them look like (e,e'p) events

- Use measured (e,e'p π) events
- Rotate p, π around q to determine π detection efficiency
- Subtract undetected (e,e'pπ)
- Repeat for higher hadron multiplicities



Data Driven Correction

Non-(e,e'p) interactions lead to multi-hadron final states Gaps can make them look like (e,e'p) events

- Use measured (e,e'p π) events
- Rotate p, π around q to determine π detection efficiency
- Subtract for undetected (e,e'p π)
- Repeat for higher hadron multiplicities (2p, 3p, 2p+1π, ...)



Subtraction Effect



Systematics: Sector Dependence



252
Systematics: Sector Dependence



Quantifying uncertainty by using unweighted variance & by subtracting variance from statistical uncertainty



- Playing this game across all nuclei & energies
- Division by $\sqrt{N}_{sectors}$
- Flat uncertainty of 6%

Calorimetric energy reconstruction using the $1p0\pi$ channel



- Area normalized results
- No information with respect to absolute scale
- G2018 offset potentially due to binding energy issue
 - +Data
 - -SuSav2 (Total) -QE -MEC -RES-DIS

Data

- Divide # events by integrated charge & target thickness to get xsec in µb
- \bullet Divide by bin width to get $\mu b/GeV$

Simulation

- Get GENIE total cross section for E_e^{\prime} / target A & Q2 > Q2_{min}
- xsec = (Selected detected events / all generated events) * total xsec / bin width

No corrections for CLAS acceptance or for bremsstrahlung radiation

Step #2: Normalized Yield



- Start from reco / true ratio w/o radiation to obtain acceptance correction
- Average on a bin-by-bin basis x = |SuSav2 + G2018| / 2
- Due to offset, G2018 Ecal predictions have been shifted by 10/25/36 MeV for 4He/12C/56Fe respectively

Step #3a: Example 12C @ 1.1 GeV



Use reco / true ratio to obtain acceptance correction

Step #3b: Radiation Correction



Averaged Acceptance Correction Uncertainty Over True Beam Energy

On a bin-by-bin basis

x = |SuSav2 - G2018| / Sqrt(12)

Bin Entry = x / Average * 100 %

Same recipe as for acceptance correction but, to avoid infinities, will use average (1 bin) around the peak and average(reco) / average(true) for correction factor

Excluding Radiation

⁵⁶Fe @ 2.2 GeV



262

Correction Factors



After both acceptance & radiation corrections, without systematics yet



Systematics

Source	Uncertainty (%)		
Detector acceptance Identification cuts φ _{qπ} cross section dependence Number of rotations	2,2.1,4.7 (@ 1.1,2.2,4.4 GeV)		
Sector dependence	6		
Acceptance correction	2-15		
Overall normalization	3		
Electron inefficiency	2		

Energy Reconstruction Accuracy

		$1.159 {\rm GeV}$		$2.257 { m ~GeV}$		$4.453~{ m GeV}$	
		Peak	Peak	Peak	Peak	Peak	Peak
		Fraction	Sum $[\mu b]$	Fraction	Sum $[\mu b]$	Fraction	Sum $[\mu b]$
⁴ He	Data	-	-	41	0.48	38	0.15
	SuSAv2		-	45	1.31	22	0.14
	G2018	-	-	39	0.93	24	0.16
¹² C	Data	39	4.13	31	1.26	32	0.34
	SuSAv2	44	5.33	27	1.76	12	0.20
	G2018	51	6.53	37	2.44	23	0.43
⁵⁶ Fe	Data	-	-	20	3.73	23	1.01
	SuSAv2	-	-	21	5.28	10	0.58
	G2018	-	-	30	8.22	19	1.48









Into The 3D e4v Multiverse!



A.Papadopoulou, et al, In preparation



Missing Momentum Approximation



$${
m p}_{
m n,proxy}=\sqrt{\delta {
m p}_{
m L}^2+\delta {
m p}_{
m T}^2}$$

Under QE assumption

Phys. Rev. Lett. 121, 022504 (2018)

Fails To Reproduce True Missing Momentum

2.261 GeV



A.Papadopoulou, et al, In preparation

 ${
m p}_{
m n,proxy}=\sqrt{\delta {
m p}_{
m L}^2+\delta {
m p}_{
m T}^2}$

Under QE assumption

Phys. Rev. Lett. 121, 022504 (2018)

True missing momentum

 $P_{miss} = |p - q|$

p = proton 3-vector q = momentum transfer

The e4v Result Factory Continued!



Inclusive Results



The e4v Result Factory Continued!

Scan over multiple anglesResults on Argon soon

```
e4v Collaboration
In preparation
```

Nuclear Sensitivity Variables



Double Differential Results

